A Multi-wavelength Study on the X-ray Emissions from Young Stellar Objects in Orion Molecular Cloud 2 and 3

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Abstract

We made a multi-wavelength study to reveal the variety and to understand the mechanisms of X-ray emissions from young stellar objects (YSOs), selecting Orion molecular cloud (OMC) 2 and 3 as our study field. OMC-2 and OMC-3 are intermediate mass star-forming regions, which contain YSOs at all evolutional classes from class 0 protostars to T Tauri stars in a wide range of mass from early-type stars to brown dwarfs.

We conducted deep observations on OMC-2 and OMC-3 in the X-ray and near-infrared (NIR) band respectively using the *Chandra X-ray Observatory* and the University of Hawaii 88 inch (2.2 m) telescope. In the X-ray band, we detected 385 sources in the Advanced CCD Imaging Spectrometer (ACIS)-I image of 17×17 arcmin^2, which is complete down to $F_{\rm X} \sim 10^{-14.5}$ ergs s^{-1} cm^{-2} with the faintest detected source of $F_{\rm X} \sim 10^{-15.5}$ ergs s^{-1} cm^{-2} in the 0.5–8.0 keV energy range. In the NIR band, we obtained the J- (1.2 μm), H- (1.6 μm), and K-band (2.2 μm) Quick Infrared Camera (QUIRC) images to extract 1448 NIR sources in a 512 arcmin^2 region. The survey is complete down to $J \sim 17.5$, $H \sim 16.5$, and $K \sim 16.0$ mag, matching well with the *Chandra* limit.

Combining the 2MASS (Two Micron All Sky Survey) and our QUIRC data, we identified the NIR counterpart for 278 (∼72%) X-ray sources (NIR-identified [NIR-IDed] X-ray sources). Most of these sources are YSOs that belong to OMC-2 and OMC-3 considering their magnitude and luminosity function in the K band. The rests of the X-ray sources are unidentified with NIR sources (NIR-unidentified [NIR-unIDed] X-ray sources).

For NIR-IDed X-ray sources, we estimated their mass and evolutional class using their J-, H-, K-band flux. We also derived their X-ray flux variability through the X-ray temporal analysis, and their plasma temperature and X-ray luminosity by the X-ray spectral analysis. By comparing the averaged X-ray properties among different mass ranges, we found that YSOs in the intermediate ($2.0\ M_\odot \leq M < 10.0\ M_\odot$), low ($0.2\ M_\odot \leq M < 2.0\ M_\odot$), and very low ($M < 0.2\ M_\odot$) mass ranges have the same X-ray emission properties in contrast to the high mass ($M \geq 10.0\ M_\odot$) sources. We further revealed that the X-ray emissions from intermediate to very low mass YSOs consist of two thin-thermal plasma components of different temperatures ($k_{\rm B}T \sim 1$ keV and 2–3 keV). Based on the time-sliced X-ray spectroscopy of some bright variable sources and on comparison with the sun and other main and pre-main-sequence sources, we proposed that the soft X-ray component is from coronae while the hard component is due to flares.

Most of the NIR-unIDed X-ray sources are background extragalactic sources from their hard X-ray spectra. However, the spatial distribution of these sources has an excess along the ridge of star-forming cloud cores, which indicates that some of the NIR-unIDed X-ray sources are related to star formations. We made follow-up imaging observations using the University of Hawaii 88 inch (2.2 m) telescope, the Subaru telescope, and the Infrared Telescope Facility in the J, H, K, L' (3.8 μm), and H_2 $v = 1 - 0$ S(1) (2.12 μm) bands in addition to the centimeter interferometer imaging observation with Very Large Array (VLA). Four NIR-unIDed X-ray sources are associated with jet and outflow systems and share many multi-wavelength characteristics in common. We proposed that these X-ray emissions can be explained by the high temperature plasma induced by protostellar jets. Other NIR-unIDed X-ray sources along the cloud core do not have such association. However, their heavily absorbed X-ray spectra and the association of some with sub-millimeter cores infer that they are heavily embedded X-ray-emitting YSOs, such as class 0 objects.

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Chapter 1

Introduction

The night sky, without stars, should be boring. The stars we see are the mixture of young and old, which indicates that our universe keeps forming stars and not all stars were born at the beginning. The sites of star formation, however, had been long inaccessible to human beings. They are born deep in dense molecular clouds, so the visual light from these baby stars is almost extinct before reaching our eyes. These star-forming clouds appear dark in visual images, hence are called dark clouds.

One of the representatives is the Orion molecular cloud (Fig. 1.1). Between two spectacular nebulae, there exists a dark lane. It appears to contain nothing at a glance, but this is where stars are born. Babies are surrounded by dense blankets of molecular clouds.

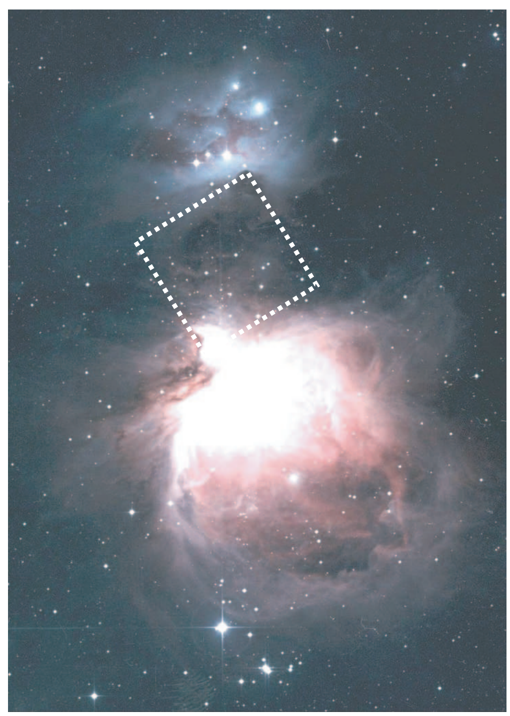

Figure 1.1: Visual image ($60' \times 94'$) of the Orion nebula (M42), the brightest nebula seen from the Earth (Guinness World Records 2003[58]). Our study field (*the dashed square*) is located $\sim 20'$ north of the nebula. Courtesy; the Kiso Observatory, the University of Tokyo.

The veil was lifted at the advent of new astronomy in the 20th century, when we became able to see the universe at any electromagnetic wavelengths. Figure 1.2 shows the absorption or scattering cross section of electromagnetic waves by the interstellar medium (ISM). Typically, cloud cores have the visual extinction of $A_V > 10$ mag. In other words, the visual light ($\lambda \approx 550$ nm) from forming stars is attenuated to less than 10^{-4} times. However, when we observe at the wavelengths of lower cross sections like X-ray, infrared, and radio lights, we can delve into the dense clouds.

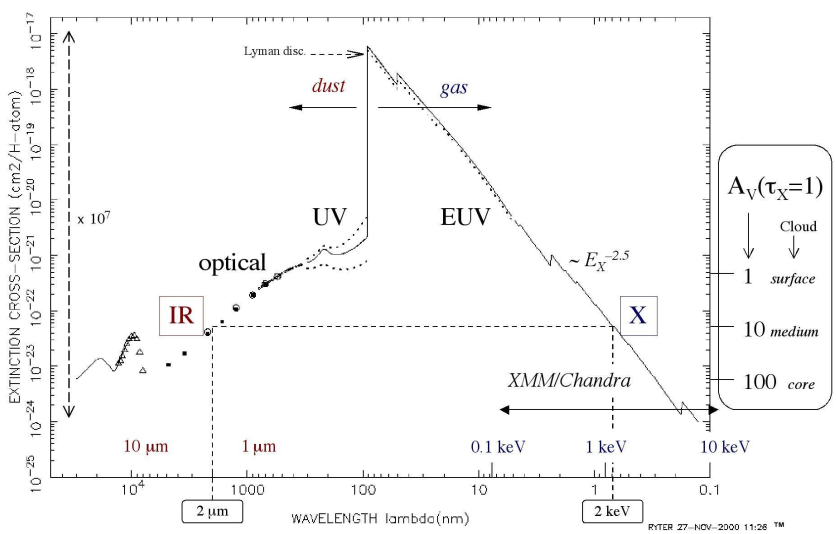

Figure 1.2: Absorption or scattering cross section by ISM of the electromagnetic wave at a given wavelength in the infrared (IR), optical, ultraviolet (UV), extreme ultraviolet (EUV), and X-ray band (Ryter 1996[168]). The *Chandra* and *XMM-Newton* satellites cover the range of 0.1–10.0 keV. The wavelengths of lower cross section have larger penetrating power into dense matter.

In 1967, Becklin & Neugebauer[19] discovered an infrared star in the Orion nebula. This object, unlike seven other objects found in their study field, had no optical counterpart. Successive near-infrared (NIR) observations found a dozen of sources with similar characteristics in dark clouds, and revealed that they are the stars at their birth. This was the first observational evidence of forming stars.

These stars are yet to begin the nuclear fusion at their center, hence are called pre-main-sequence sources or young stellar objects (YSOs). They are cool (T = 100–1000 K at the surface) and embedded sources that are just generated from gravitationally contracting cold ($T \sim 10$ K) gas. It was astonishing, therefore, that strong X-ray emissions corresponding to the temperature of $T \sim 10$ MK were found from these sources. Babies were found to be crying loud! In 1980's, the *Einstein Observatory* detected soft ($E < 2$ keV) X-ray emissions from T

Tauri stars (TTSs), which comprise a class of YSOs with the mass of 0.5–2.0 $M_{\odot}$[1] and the age of ∼1 Myr (Feigelson, & DeCampli 1981[50]; Montmerle et al. 1983[133]). Using the *Advanced Satellite for Cosmology and Astrophysics* (*ASCA*) satellite, Koyama et al. (1996)[107] further discovered hard ($E > 2$ keV) X-ray emissions from a cluster of class I protostars, which are low-mass YSOs younger than TTSs (∼0.1 Myr). Protostars are severely extinct by dense cores and are hard to access even with NIR observations.

The X-ray studies on YSOs are important in a wide area of astronomy and astrophysics. First, the X-ray emissions from these stars are so intense that they can affect heavily on the circumstellar environment. Glassgold, Feigelson, & Montmerle (1999)[70] discussed that the photoionization of circumstellar matter by X-rays can surpass the collision ionization by cosmic rays inside the distance from the core of

$$r \approx 0.02 \times \left(\frac{L_{\rm X}}{10^{29}\ {\rm ergs\ s^{-1}}} \right)\ [{\rm pc}], \tag{1.1}$$

where $L_{\rm X}$ is the X-ray luminosity of YSOs. At active phases of YSOs, $L_{\rm X}$ can exceed 10^{30}–10^{31} ergs s^{-1}, which indicates that the X-ray ionization is the dominant mechanism inside the whole cloud core of typically ∼0.1 pc. As circumstellar matter should be ionized before being coupled with the magnetic field, the ionization rate and efficiency are key parameters in star formations. Without establishing the average features and revealing the mechanism of X-ray emissions from YSOs, therefore, we would not be able to understand the star formation process, a long-standing subject in astronomy and astrophysics.

Once the ubiquity of the X-ray emission is established for protostars as well as for TTSs, X-ray observations will be one of the most efficient tools to conduct a census on star-forming clouds. With a larger penetrating power than NIR observations (particularly of hard X-rays; Fig. 1.2), with a better spatial resolution than FIR–millimeter observations, and with a very wide field of view, we will soon encounter a large number of "X-ray-selected" YSOs. The investigation on these sources, we believe, will proceed our understanding on star births.

X-ray observations on protostars and TTSs with the mass of ∼1 $M_{\odot}$ also give us a picture on the past activities of our sun. Its enhanced activity by 3–5 orders of magnitude in the past should have given a large impact on the evolution of the solar system.

Finally, YSOs are one of the enigmatic sources in the universe that emit radiation by releasing their gravitational energy, such as neutron stars, galactic black holes, and active

[1]Masses in the unit of solar mass (2.0×10^{33} g) are given with $M_{\odot}$.

galactic nuclei (AGNs). These sources have many processes in common; formation of jets, gravitational energy release through accretion disks, and acceleration of charged particles. YSOs are much closer to us than the other types of sources, which enables us to investigate fine structures of these high energy phenomena.

Previous studies with the *Einstein*, *ROSAT*, and *ASCA* satellites on the X-ray emissions from YSOs are mainly focused on low-mass sources. Higher mass YSOs have lower population densities and evolve more quickly, making these samples fewer and more distant. Lower mass YSOs are also difficult to observe because of their intrinsic faintness. The X-ray spectra of low-mass YSOs revealed that the X-ray emissions are of thin-thermal plasma origin. In addition, their X-ray light curves often show flares with a rapid rise and a slow decay, which are similar in profile but higher in flux and temperature compared to solar flares. These lead to the general consensus that the X-ray emissions from low-mass YSOs are solar-like; YSOs have thin-thermal plasma to emit X-rays with occasional flares triggered by magnetic reconnections. However, not all observational results are accountable by a simple extension of the solar flares. Moreover, YSOs of different mass ranges ($M < 0.2\ M_\odot$ or $M > 2.0\ M_\odot$) are not well studied. What are different between YSOs and our sun? Do higher-mass or lower-mass YSOs possess high temperature plasma like low-mass YSOs? If so, how is the plasma generated? Are there any other X-ray emission mechanisms besides flares? The aim of this thesis is to give an answer to these questions through a multi-wavelength study on Orion molecular cloud 2 and 3 (OMC-2 and OMC-3).

Our study field is located at the center of the dark lane in Figure 1.1 and contains YSOs at all evolutional phases from class 0 protostars to TTSs in a wide range of mass from early-type stars to brown dwarfs. OMC-2 and OMC-3 are full of phenomena related to star formation; radio jets, molecular outflows, Herbig-Haro (HH) objects, sources with accretion disks, dust emissions, etc. In particular, the condensation of protostars in OMC-2 and OMC-3 is the highest among all near-by star forming regions (Reipurth, Rodríguez, & Chini 1999[162]). They are proximate to us at a distance of $\sim$450 pc (Genzel & Stutzki 1989[66]) and have an appropriate size of $10' \times 20'$ that can be covered with one field of view of the current X-ray cameras. All these features make these clouds ideal sites for our study.

In order to study the X-ray emissions from YSOs, we employ a multi-wavelength approach. OMC-2 and OMC-3, although hard to see in the optical light, can be accessed with X-ray, NIR, and radio observations (see our X-ray and NIR images of the field in Figs. 5.2

and 5.9). Information obtained at different wavelengths complements with each other. In addition to the X-ray observations that trace high temperature (T =1–100 MK) plasma of YSOs, NIR–millimeter observations detect blackbody emissions from stars, accretion disks and dusts, which indicate the mass and evolutional phase. Outflows and jets from YSOs can be seen through the optical, NIR, millimeter, and centimeter observations. In this thesis, we conduct a deep X-ray and NIR observations on OMC-2 and OMC-3 respectively using the *Chandra X-ray Observatory* and the University of Hawaii 88 inch (2.2 m) telescope. Follow-up observations are also performed on interesting sources using the Subaru telescope, the Infrared Telescope Facility (IRTF), and the University of Hawaii 88 inch (2.2 m) telescope at the NIR band and the Very Large Array (VLA) at the centimeter band.

The plan of this thesis is as follows. We review on the recent progress of X-ray observations on YSOs in Chap. 2 and previous multi-wavelength observations of our study field in Chap. 3. The basic features of the telescopes and instruments used in this work are summarized in Chap. 4. In Chap. 5, we explain our X-ray and NIR observations on OMC-2 and OMC-3. In total, 385 X-ray sources are found in our observation. By correlating them with our 1448 NIR sources as well as NIR all-sky survey catalog sources, we discriminate the NIR-identified (NIR-IDed) and NIR-unidentified (NIR-unIDed) X-ray sources. The following two chapters (Chap. 6 and Chap. 7) are devoted for the NIR and X-ray properties of the NIR-IDed X-ray sources, where we estimate the mass and the evolutional phase of our samples and discuss the X-ray spectral and temporal features of these sources. In Chap. 8, we deal with the NIR-unIDed X-ray sources, where we focus on sources related to star formation. In Chap. 9, we discuss the X-ray emission mechanisms of both the NIR-IDed and NIR-unIDed sources. Finally, we conclude in Chap. 10. Figure 10.1 will help readers to follow the flow of this thesis.

Chapter 2

X-ray Observations on YSOs

Contents

In the first section (Sect. 2.1) of this chapter, we briefly review the basic feature of YSOs. In the following section (Sect. 2.2), we review the X-ray observations on YSOs. As the progress made on low-mass YSOs with the *Einstein*, *ROSAT*, and *ASCA* satellites are concisely summarized in a review paper by Feigelson & Montmerle (1999)[53], we focus mainly on recent advances made with the following *Chandra* and *XMM-Newton* observations. The problems to be solved on the X-ray emissions from YSOs are summarized in the last subsection, for which this thesis tries to give an answer.

2.1 Basic Features of YSOs

2.1.1 Evolution and Classification

A star formation begins with the onset of gravitational collapse of molecular cloud cores with the density of 10^{-17}–10^{-16} g cm^{-3}. Gas and dust keep contracting freely on the central core until the density reaches $\sim 10^{-13}$ g cm^{-3}. At this density, the core becomes optically thick. The thermal pressure adiabatically increases and finally balances with the gravitational force to stop dynamical accretion. YSOs at this stage are called protostars. Protostars commonly accompany jets and outflows, which function to release angular momentum. Without the mechanism of angular momentum release, they can not keep accumulating mass because matters inside the Keplerian radius can not accrete onto their surface.

After most of the circumstellar matter accretes onto the star or is blown away by jets and outflows, YSOs quasi-statically contract with increasing temperature and density. YSOs at this stage are called T Tauri stars (TTSs). The contraction continues with a constant surface temperature, which makes TTSs move down along the Hayashi track on the Hertzsprung-Russell (H-R) diagram (Hayashi 1966[81]). The circumstellar disk gradually disappears as TTSs evolve. TTSs are divided into two classes; classical TTSs (cTTSs) that are in the earlier phase of TTSs with an optically thick accretion disk and weak-lined TTSs (wTTSs) that are in the later phase with an optically thin or no accretion disk. Finally, when nuclear burning starts at the center, YSOs turn into the main sequence stars. In general, "pre-main-sequence source" indicates TTSs, while "YSO" indicates the collection of protostars and TTSs.

The evolution of low-mass YSOs, which have been the main targets of previous observations due to their proximity and ample samples, is traced by several measurements. One is to use the equivalent width (EW) of the H_α emission in the optical band, which is a direct indicator of mass accretion. CTTSs are defined as YSOs with $EW > 10$ Å while wTTSs are with $EW < 10$ Å. As cTTSs are younger ($\sim 10^6$ yr in case of low-mass YSOs) than wTTSs ($\sim 10^7$ yr) and have mass accretion from their circumstellar disk, they show larger EW values than wTTSs.

YSOs in much younger phases than TTSs are invisible with the optical light. Lada & Wilking (1984)[112], Adams et al. (1987)[1], and Lada (1991)[113] established a classification scheme using the NIR to mid-infrared (MIR) excess, in which younger YSOs show stronger NIR excess emission in addition to their photospheric blackbody emission. This excess emission comes from circumstellar disks and envelopes that have lower temperature ($T =$ 100–1000 K) than the photosphere ($T \sim 3000$ K). The spectral index between 2.2 μm (the K band) and 10 μm (the M band) in the spectral energy distributions (SEDs) of YSOs;

$$\alpha = -\left.\frac{d\log \nu F_\nu}{d\log \nu}\right|_{2.2-10\ \mu\mathrm{m}} \tag{2.1}$$

is used to classify sources into class I, class II, and class III (Fig. 2.1). As class I sources are younger ($\sim 10^5$ yr), they show larger NIR excess hence larger α values than class II and class III sources. When the M-band flux is not available, which is often the case because of the difficulty of sensitive observations in this band, the J (1.2 μm)-, H (1.6 μm)-, and K-band fluxes are alternatively used to discriminate among classes using the color-color diagram (Sect. 6.1). CTTSs usually have the class II SED, while wTTSs have the class III SED. In this thesis, we use cTTSs and wTTSs for the same meaning as class II and class III sources.

In addition to the class I, II, and III objects, André, Ward-Thompson, & Barsony (1993)[5] introduced class 0 objects, which are younger than class I. The lower temperature and higher obscuration make this class sources invisible even with NIR and MIR lights. Barsony (1994)[17] defined the class 0 sources with the following features: (1) undetectable at the wavelengths of $\lambda < 10$ μm, (2) a high ratio of $L_{\mathrm{smm}}/L_{\mathrm{bol}}$, where L_{smm} and L_{bol} are the sub-millimeter and bolometric luminosity, (3) a relatively narrow SED resembling that of a single blackbody at $T \leq 30$ K, (4) presence of a molecular outflow, (5) existence of centimeter continuum emission, and (6) presence of HH objects. Class 0 objects are considered to be at the age of 10^4–10^5 yr, which are the youngest stellar objects that are accessible with the current observational technology.

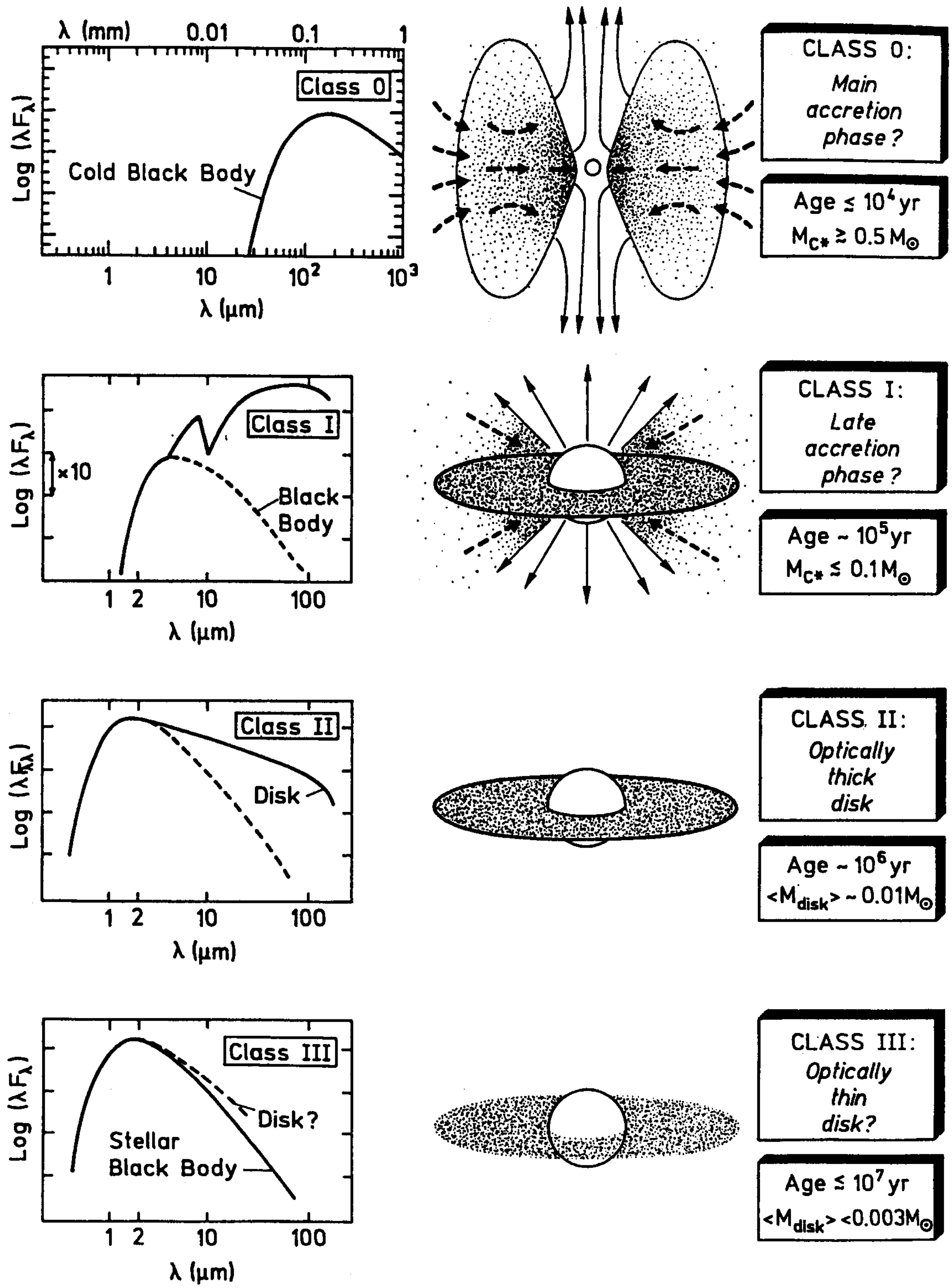

Figure 2.1: Classification of YSOs and their SEDs (Bachiller 1996[13])

Table 2.1: Multi-wavelength properties of low-mass YSOs

	YSOs				main
	class 0	class I	class II	class III	sequence
properties	infalling protostar	evolved protostar	classical T Tauri	weak-lined T Tauri	
age (yr)	10^4	10^5	10^6	10^7	$>10^7$
accretion disk	thick	thick	thick	thin or nonexistent	planetary system (?)
X-ray emission	?	strong	strong	strong	weak
jet and outflow	yes	yes	yes	no	no
non-thermal radio	?	yes	no (?)	yes	yes

2.1.2 Jets and Outflows

Jets and outflows from YSOs are traced using several observational tools. From the origin, radio jets at the distance of ∼100 AU, optical jets and fast neutral winds at the distance of ∼0.01 pc, and molecular outflows at the distance of 0.1–1 pc are ubiquitously observed from protostars. More evolved YSOs accompany weak or no jet and outflow phenomena.

Radio Jets

Radio jets from YSOs were first studied by Cohen, Bieging, & Schwartz (1982)[37]. They appear to be thermal emissions at the centimeter band and are detected from almost all the class 0 and some of the class I protostars. The observational characteristics of these emissions in near-by star-forming regions are summarized as (1) the elongated morphology at sub-arcsecond scale, (2) association with both high and low luminosity objects, (3) alignment within a few degrees with the large scale outflow, (4) weak flux with the flux density typically in the range of 0.5–10 mJy, (5) positive or flat spectral indices, and (6) dynamical time scales of only a few years (Anglada 1996[8]).

Because these emissions are associated with low as well as high luminosity YSOs, the free-free emission from the H_{II} region produced by stellar photo-ionization is unlikely. The alignment with and the elongation in the direction of global outflows strongly infer that these emissions are related to jet and outflow phenomena. Curiel, Cantó, & Rodríguez (1987)[42] and Curiel et al. (1989)[43] proposed an idea of shock-induced ionization, where protostellar jets collide into high density ambient matter at the speed of a few times of 100 km s^{-1}, and the UV photons are irradiated from the shock front to produce H_{II} regions. An empirical relation is seen between the outflow momentum rate and the centimeter flux, which supports this scenario (Anglada 1992[7]). Martí, Rodríguez, & Reipurth (1998)[122] monitored one of the radio jets from YSOs and found that a pair of centimeter condensations are traveling away from the central core at the speed of ~500 km s^{-1}.

Optical Jets

Because Herbig (1951)[82] and Haro (1952)[78] were the first to identify knotty structures in the Orion region, they are called HH objects. Together with optical jets that were detected later by Mundt & Fried (1983)[138] and Mundt et al. (1984)[139], these are the typical phenomena that characterize star-forming regions. These emissions are considered to originate from cooling regions of high velocity shocks that are produced as a result of highly collimated jets from YSOs colliding into dense ambient matter.

HH objects and optical jets altogether have the following characteristics: (1) knotty structures sometimes spaced periodically, (2) bow-shock-shaped HH objects located at the end of a series of knots, (3) half of the HH objects and optical jets with a bipolar structure, (4) collimated structure of the opening angles in the range of 5°–20°, and (5) the typical length, velocity and number density of optical jets in the range of 0.01–1 pc, 200–600 km s^{-1}, and 10–200 cm^{-3} (Edwards, Ray, & Mundt 1993[46]).

Fast Neutral Winds

CTTSs are known to posses high velocity (100–300 km s^{-1}) wind based on the P Cygni-type profile in various lines in the optical wavelengths (Kuhi 1964[110]; Hartmann 1982[80]). Similar high velocity winds were also detected from much younger phase of YSOs (protostars) using the neutral H_I line (Giovanardi et al. 1992[69]). Although cTTSs and protostars are in different phases of the YSO evolution, the stellar wind from these sources is considered to have the same origin because they have similar velocities, mechanical energies, and momentum flux. The origin of the phenomena is not well known, but it is somewhat related to accretion disks. A detailed discussion on this topic can be found in Edwards et al. (1993)[46].

Molecular Outflows

Circumstellar matter is dragged by high velocity and high density jet from protostars, which forms molecular outflows moving at the velocity of 10–200 km s^{-1}. They are well traced by rotational transition lines of molecules in the millimeter band and vibrational-rotational transition lines in the NIR band. Snell, Loren, & Plambeck (1980)[174] was the first to detect a bipolar CO outflow from L1551 region. More than 200 outflows are cataloged to date.

The vibrational-rotational transition lines of H_2 and the rotational transition lines of CO are commonly used. The H_2 and CO outflows together with the optical jets are powerful tracers of jet and outflow activities from protostars that work complementarily, where the lowest velocity ($\sim$10 km s^{-1}) component is traced by CO, the moderate velocity ($\sim$100 km s^{-1}) component by H_2, and the highest velocity ($\sim$500 km s^{-1}) component in the optical wavelengths.

About a half of molecular outflows are observed to have a bipolar structure with a pair of red-shifted and blue-shifted robes and the powering source is always at the center of the symmetry. They are moderately collimated with the opening angle of 10°–50°. Recent progress of observations on molecular outflows can be found in Fukui et al. (1993)[59], Eislöffel et al. (2000)[47], and Richer et al. (2000)[163].

2.1.3 Magnetic Activities

As YSOs are fully convective in their interior and rotate much faster ($\sim$1 day) than main sequence stars ($\sim$30 days in case of the sun), they are expected to possess enhanced magnetic activities. This is supported by three major observational indications: (1) X-ray flares that trace the magnetic reconnections, (2) centimeter gyro-synchrotron emissions that indicate accelerated charged particles to the energy of $\sim$1 MeV, and (3) various optical photometric and spectroscopic studies. X-ray emissions are separately reviewed in Sect. 2.2.

Gyro-synchrotron Emissions

In addition to the thermal free-free emissions, YSOs emit non-thermal emissions in the centimeter band. These emissions are considered to originate from the gyro-synchrotron emissions for the following reasons (André 1996[6]). First, they show flux variability of the time scale of hours to days, including intense "radio flares" (Feigelson & Montmerle 1985[51]; Becker & White 1985[18]; Bieging & Cohen 1989[23]). Second, they have a moderate degree of circular polarization of $|V/I|$ <5% at outbursts and $|V/I|$ <20% at quiescence. Third, a brightness temperature of $>10^7$ K is measured by Very Long Baseline Interferometer (VLBI) observations, directly confirming the size of the emitting regions. Electrons with the energy of ∼1 MeV are gyrating around large-scale magnetic fields with the size of ∼5 stellar radii and the strength of 1–10 G, which emit non-thermal emissions detectable at the centimeter band.

These emissions are only seen in evolved YSOs except for a class I protostar (Feigelson, Carkner, & Wilking 1998[52]). The centimeter flux of this emission is somewhat related to the X-ray luminosity, which can be extended to magnetically-active main sequence stars such as RS CVns and dMe stars. A recent simultaneous observation on an RS CVn-type binary in the X-ray and the centimeter band using *XMM-Newton* and VLA reported the detection of the Neupert effect, in which the time derivative of the X-ray light curve is proportional to the centimeter light curve (Güdel et al. 2002[75]). This confirms to understand the phenomena in solar analog; i.e., accelerated electrons via magnetic reconnections reach the chromosphere to heat the plasma that emits X-rays.

Optical Studies

The most successful approach to measure the magnetic field strength of TTSs is the Zeeman splitting measurements, although this is applicable only to ideal cases. Some TTS samples were studied and revealed to have a strong magnetic field with the strength of 0.5–4 kG (Johnstone & Penston 1986[97]; Johnstone & Penston 1987[98]; Johns-Krull, Valenti & Koresko 1999[95]).

There are many indirect methods to investigate magnetic activities. Doppler imaging technique showed that large cool spots are moving on the surface of some TTSs. H_α and Mg_{II} emission lines, which are the tracers of chromospheric activities, are other indicators of strong magnetic activities of these sources.

2.2 X-ray Observations on YSOs

2.2.1 Low-mass YSOs: (1) Recent Results with *Chandra*

In X-ray observations as well as observations in other wavelengths, low mass (0.2–2.0 $M_{\odot}$) YSOs have been the main targets. Previous studies using *Einstein*, *ROSAT*, and *ASCA* are concisely summarized in a review paper by Feigelson & Montmerle (1999)[53]. Here, we briefly review recent results made with *Chandra* (Table 2.2).

Two capabilities are required for the X-ray studies on YSOs: (1) the spatial resolution of $\sim 1''$ to resolve each YSO member in a usually crowded star-forming region and (2) the sensitivity in the hard X-ray band ($E > 2$ keV) to penetrate through dark clouds. The preceding X-ray observatories did not meet either of these two requirements; *ASCA* the former and *Einstein* and *ROSAT* the latter. *Chandra* has both capabilities, which makes this satellite the best tool for YSO studies. The charge coupled device (CCD) detector onboard *Chandra* further enables us to obtain temporal ($\Delta t \sim 3$ s) and spectral ($\Delta E \sim$ 100 eV) information on a photon basis. With a wide field of view (FOV) of $17' \times 17'$, we can simultaneously derive the light curve and spectrum of ~100 X-ray sources in a star-forming region (Fig. 2.3). Table 2.3 summarizes the survey studies made with *Chandra*.

The X-ray properties of TTSs were found to have the following characteristics: (1) thin-thermal plasma spectrum with the X-ray luminosity of $L_{\rm X} < 10^{30}$–10^{31} ergs s^{-1} and the plasma temperature of $k_{\rm B}T = 0.5$–5 keV, (2) flare-like flux variability with some hints of spectral hardening, (3) an empirical relation of $L_{\rm X}/L_{\rm bol} = 10^{-3}$–$10^{-5}$, and (4) no clear difference in X-ray properties between cTTSs and wTTSs. All these results are already obtained with previous *Einstein*, *ROSAT*, and *ASCA* observations (Feigelson & Montmerle 1999[53]). *Chandra* studies on TTSs confirmed these results with larger number of samples.

Major advances are seen in the studies of class I protostars. In the era of *ASCA* and *ROSAT*, only a limited number of class I samples were identified (Fig. 2.3, Koyama et al. 1996[107]; Kamata et al. 1997[103]; Ozawa et al. 1999[145]; Tsuboi et al. 2000[187] using *ASCA* and Grosso et al. 1997[72]; Neuhäuser & Preibisch 1997[141]; Grosso 2001[73] using *ROSAT*). With poor photon statistics, it was difficult to address their X-ray properties. A textbook case of the X-ray studies on class I protostars was presented by Imanishi, Koyama, & Tsuboi (2001a)[91] using the *Chandra* deep exposure observation on the ρ Ophiuchi dark cloud (Fig. 2.4). They detected X-rays from ~70% of class I protostars and their candidates.

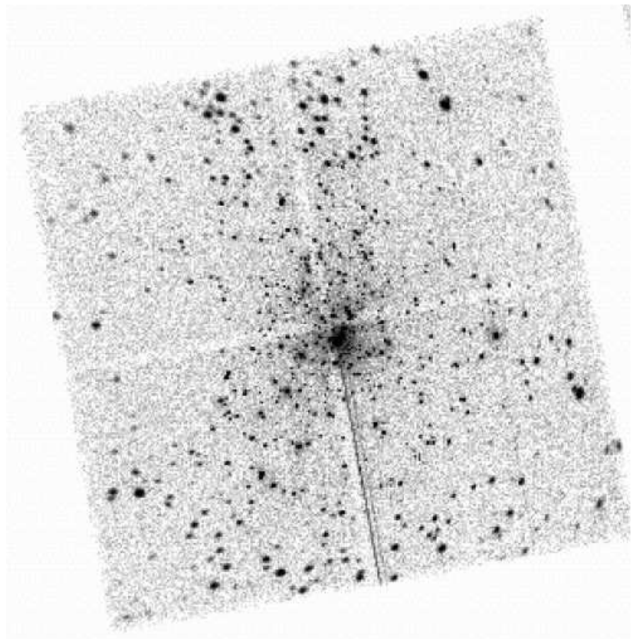

Figure 2.2: *Chandra*/ACIS image of ONC (Feigelson et al. 2002b[55]). Inside a $17' \times 17'$ region of ACIS-I, $\sim$1000 X-ray sources were detected.

Moreover, they investigated their X-ray spectra and light curves and found that class I protostars have similar X-ray features with TTSs, with higher plasma temperature, larger column density and X-ray luminosity, and more occasional flare-like events.

The X-ray sources in star-forming regions that are not identified with NIR and optical sources can be an interesting topic. The number of these sources is 30–100 in some (Table 2.3). Although this number depends on the significance level of the source detection algorithm and the completeness limit of the NIR and optical catalog, it generally exceeds the expected number of background sources, particularly in cases of Orion Nebula Cluster (ONC; Feigelson et al. 2002b[55]) and OMC-2/3 (Tsujimoto et al. 2002a[189]; 2003[192] [this thesis]) observations. The nature of these X-ray sources is not studied at all yet. However, considering the high sensitivity of *Chandra* observations that are comparable or deeper than optical and NIR observations, some of these optically and NIR unidentified sources can be a new class of YSOs.

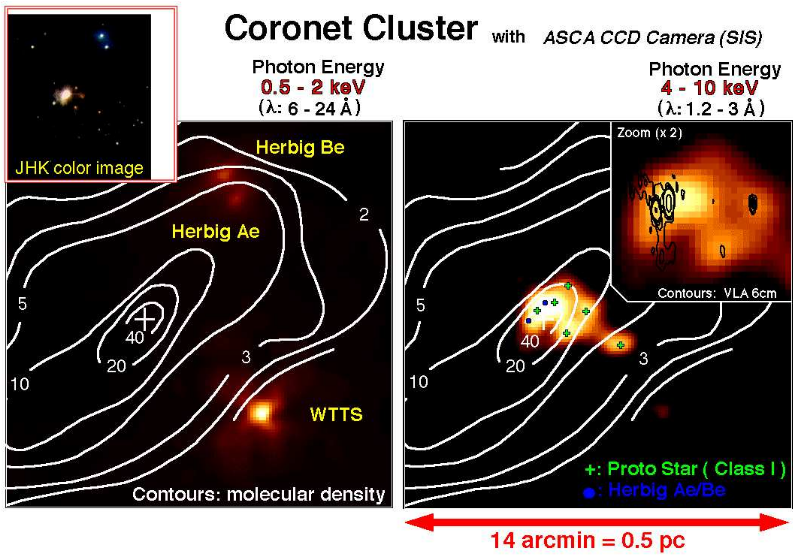

Figure 2.3: Hard X-ray detections from class I protostars in R Coronae Australis with *ASCA* (http://www-cr.scphys.kyoto-u.ac.jp/IAU/gallery/gallery.html). Unlike soft X-rays (*left*), hard X-rays (*right*) can penetrate into the densest part of the cloud core, detecting hard X-ray emissions for the first time from class I protostars (Koyama et al. 1996[107]).

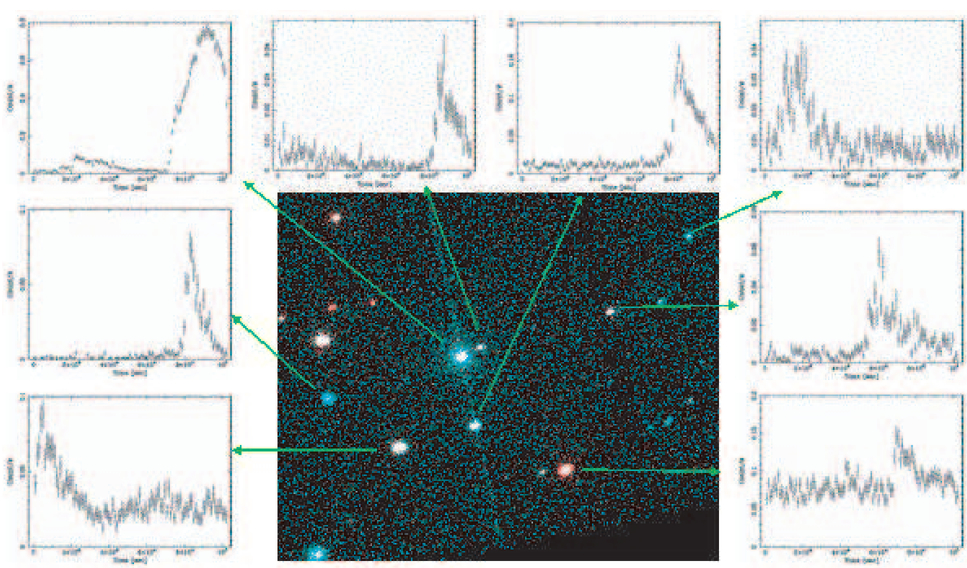

Figure 2.4: X-ray detections and their light curves from class I protostars in the ρ Ophiuchi cloud with *Chandra*. Courtesy; Kensuke Imanishi (Kyoto University).

Table 2.2: *Chandra* and *XMM-Newton* studies of star-forming regions

categories	regions (references)
survey studies . .	ρ Oph (Imanishi et al. 2001a[91]; 2001b[92]; 2002[93])
	IC 348 (Preibisch & Zinnecker 2001[155]; 2002a[156])
	NGC 1333 (Getman et al. 2002[67])
	ONC (Garmire et al. 2000[62]; Schulz et al. 2001[169], Feigelson et al. 2002a[54]; 2002b[55]; 2003[56])
	OMC-2/3 (Tsuboi et al. 2001[188]; Tsujimoto et al. 2002a[189])
intermediate- &	Mon R2 (Kohno et al. 2002[106]; Preibisch et al. 2002b[157]), Sgr B2 (Takagi et al. 2002[178])
high-mass YSOs	W 3 (Hofner et al. 2002[90]), IRAS 19410+2336 (Beuther et al. 2002[22]), M 8 (Rauw et al. 2002[159])[a]
TTSs (grating) .	TW Hya (Kastner et al. 2002[104])
HH objects	HH-2 (Pravdo et al. 2001[153]), L1551 IRS5 (Favata et al. 2002[49])[a]
diffuse emissions	NGC 3603 (Moffat et al. 2002[132]), RCW 38 (Wolk et al. 2002[195])
	Arches Cluster (Yusef-Zadeh et al. 2002[198]), M 8 (Rauw et al. 2002[159])[a]

[a] Studies by *XMM-Newton*. Other studies were made with *Chandra*.

Table 2.3: Detected numbers of X-ray sources in *Chandra* survey studies

region	distance	exposure	number of sources				references
	(pc)	(ks)	all sources (SL[a])	class I[b]	BD[c]	no ID[d](CL[e])	
ρ Oph	165	101	87 (1×10^{-7})	18	7	28 (K < 15)	Imanishi et al. (2001a)[91]; 2001b[92]
IC 348	310	53	215 (3×10^{-6})	13	2	39 (K < 19)	Preibisch & Zinnecker (2001)[155]
NGC 1333	318	38	127 (1×10^{-6})	8	N/A	32 (I < 23)	Getman et al. (2002)[67]
ONC	450	83	1075 (1×10^{-5})	N/A	30	101 (V < 20)	Feigelson et al. (2002a)[54]
OMC-2/3	450	89	385 (1×10^{-5})	13	12	107 (K < 16)	Tsujimoto et al. (2002a[189]; 2003[192]) [this thesis]

[a] The significance level in the X-ray source detection algorithm. The wavdetect program was used in all studies.

[b] The number of X-ray-emitting class I protostars.

[c] The number of X-ray-emitting brown dwarfs, which includes brown dwarf candidates identified with NIR photometry observations alone.

[d] The number of X-ray sources not identified with optical and NIR sources.

[e] The band and the limiting magnitude of the optical and NIR catalogs that were used to correlate with X-ray sources.

2.2.2 Low-mass YSOs: (2) X-ray Emission Mechanisms

The X-ray emission mechanism of YSOs should account for many observational results obtained so far: (1) the soft X-ray emissions ($E < 2$ keV) as well as hard X-ray emissions ($E > 2$ keV), (2) flare-like variability as well as quiescent emissions, and (3) some empirical relations, such as $L_{\rm X}/L_{\rm bol} \sim 10^{-3}$–$10^{-5}$. The proposed scenarios do not successfully account for all of these pieces of evidence.

The most favored idea is the magnetic reconnection model, in which rapid energy release by reconnections of magnetic field lines (flares) generates and maintains the high temperature plasma. This model well explains the flare-like flux variability with rapid rise and slow decay and hard X-ray spectra exceeding 2 keV.

Four geometries of magnetic field lines are considered (Feigelson & Montmerle 1999[53]; Fig. 2.5): (1) both feet on the stellar surface, (2) connecting the star and the circumstellar disk, (3) above the corotation radius, and (4) both feet on the circumstellar disk. The geometry (1), which is the same with the sun, is most plausible. The reconnections are caused as a result of differential rotation and convection of the star. In contrast, Montmerle et al. (2000)[134] proposed a magnetic field geometry bridging between the star and the disk (geometry 2). This is based on the *ASCA* result by Tsuboi et al. (2000)[187], who detected quasi-periodic flares from a class I protostar in the ρ Ophiuchi cloud (YLW 15; Fig. 2.6). They assumed the radiative cooling for the flux decay in these flares, and derived that the magnetic loop length is ∼14 times of the stellar radius. Favata, Micela, Reale (2001)[48] noted that the loop length is overestimated unless we take the recurrent heating during flares into account. At present, no clear difference is seen between cTTSs and wTTSs in X-ray luminosity functions, which casts some doubt on the geometry where the disk plays a key role (geometry 2 and 4).

Whatever the geometry is, the magnetic reconnection model does not account for the steady X-ray emissions at the quiescent phase. In the case of our sun, attempts have not been successful to explain the steady emissions (the solar coronae) by the integration of small flares (Shimizu 1995[170]; Aschwanden et al. 2000[9]; Aschwanden & Charbonneau 2002[10]). Shimizu (1995)[170] discussed, using the *Yohkoh* Soft X-ray Telescope data, that the occurrence rate of flares is expressed with a single power-law function of the flare energy (E) as

$$\frac{dN}{dE} \propto E^{-\alpha}. \tag{2.2}$$

He derived $\alpha = 1.5$–1.6 in the energy range down to $\sim 10^{27}$ ergs. The α value of less than

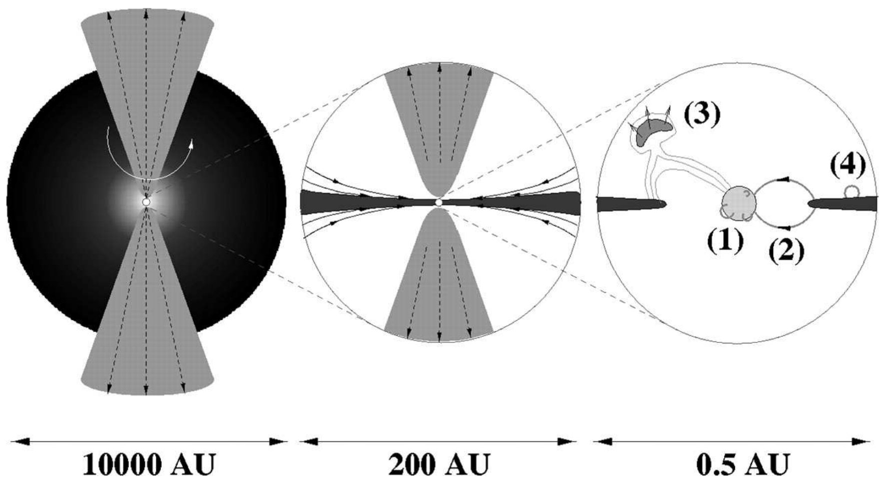

Figure 2.5: Magnetic field geometries of YSOs (Feigelson & Montmerle 1999[53]).

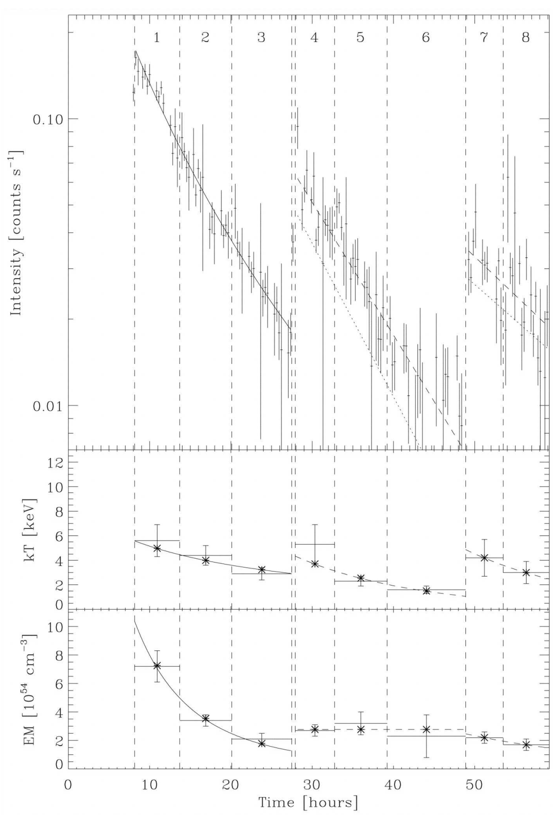

Figure 2.6: Quasi-periodic flare from a protostar in the ρ Ophiuchi cloud detected with *ASCA* (Tsuboi et al. 2000 [187]). The time evolution of the X-ray luminosity, plasma temperature, and emission measure are given in the upper, middle, and lower panels, respectively.

2 indicates that smaller flares are not the dominant source of the total energy released by flares of all scales, if we assume that the power-law can be extrapolated to lower energies. In fact, the integration of all observed flares account for at most only one fifth of the heating rate required for the active-region coronae. With this in mind, we have to deal with flares and coronae as different phenomena, although the heating mechanism of coronae is still quite puzzling.

Another difficulty that the magnetic activity scenario is confronted with is that we see no clear relation between the X-ray luminosity and the rotation period (e.g.; Feigelson et al. 2003[56]). This is in sharp contrast with magnetically active main sequence stars, where we see a strong anti-correlation between the soft X-ray luminosity and rotation period (Pallavicini et al. 1981[147]). This should be a natural deduction when the magnetic activity is the origin of the X-ray emissions.

Skinner & Walter (1998)[173] presented an interesting result using the *ASCA* observation on the brightest cTTSs in the Taurus-Auriga complex (SU Aur). They argued that there exist two different X-ray emissions by constructing the differential emission measure (DEM) distribution of this source. The bimodal plasma temperature distribution is quite similar to those seen in active late-type main sequence stars with X-ray emissions of coronal origin. If the X-ray emissions from YSOs are commonly the combination of some different mechanisms, all previous discussions correlating L_{X} or $k_{\mathrm{B}}T$ with stellar parameters should be re-examined.

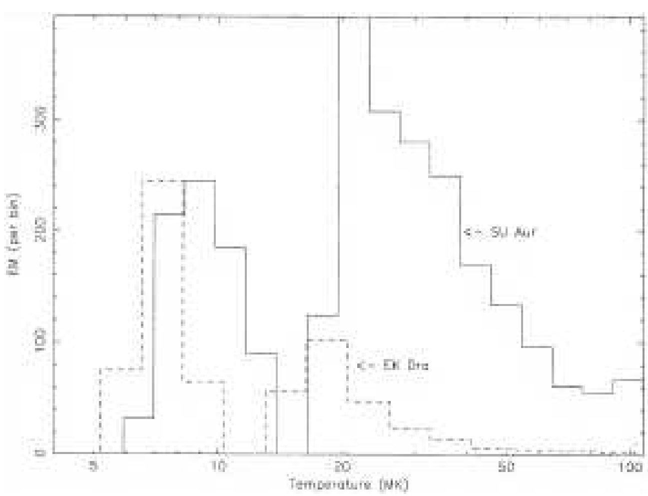

Figure 2.7: Differential emission measure distributions of SU Aur (a pre-main-sequence source; *solid*) and EK Dra (a main sequence star; *dashed*) derived from *ASCA* observations (Skinner & Walter 1998[173])

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Kastner et al. (2002)[104] proposed yet another scenario based on their high-resolution grating spectroscopy observation on a cTTS (TW Hya) using *Chandra* (Fig. 2.8). They argued that the X-ray emission arises from the hot spot or spots on the stellar surface located at the root of accretion funnels. Three lines of evidence support this idea: (1) The DEM distribution is sharply peaked at $T \sim 3$ MK, which is in sharp contrast with broad DEMs seen in coronally active late-type stars. (2) The temperature of $T \sim 3$ MK is consistent with the value expected from the adiabatic shock caused by gas accretion at the speed of 150–300 km s^{-1}, which is measured from the broadening of H_α line emissions. (3) The high-resolution spectroscopy enabled them to measure the plasma density using the density-sensitive triplet lines of O_{VII} and Ne_{IX}. In combination with the emission measure, the emitting volume was derived to be only 10^{-6} of the stellar volume, which is smaller than typical coronal X-ray sources. We have no other grating spectroscopy results on YSOs besides this source. Further studies with a larger number of samples are quite prospective.

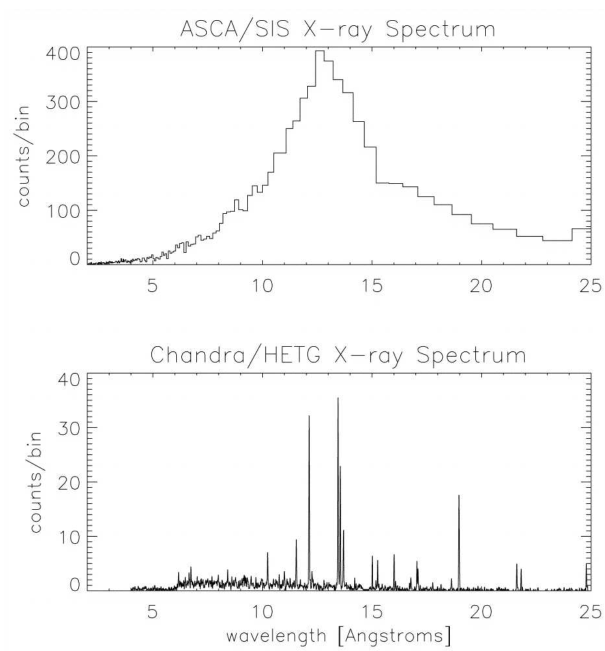


Figure 2.8: X-ray spectra of TW Hya taken with *ASCA*/SIS (a CCD spectrometer; *upper pannel*) and *Chandra*/HETG (a grating spectrometer; *lower panel*). The density sensitive triplet lines of O_{VII} and Ne_{IX} are clearly resolved at ∼22 Å and ∼11 Å in the grating spectrum, which serve as a powerful tool to determine the electron density of the plasma (Kastner et al. 2002[104]).

Zhekov, Palla, & Myasnikov (1994)[201] presented a model calculation to account for the soft X-ray emission with colliding winds from YSO binaries at the speed of 300–500 km s^{-1}

and the mass loss rate of 10^{-8}–10^{-6} $M_\odot$ yr^{-1}. This model, although it well explains the soft X-ray temperature and the X-ray luminosity, has difficulties in understanding the hard X-ray emissions and flaring temporal behaviors of YSO X-rays.

2.2.3 Very Low Mass YSOs (=Young Brown Dwarfs)

Recent X-ray observations are also capable to study brown dwarfs in star-forming regions. Brown dwarfs are tiny objects with mass less than 0.08 $M_\odot$ that never possess nuclear burning during their life. A series of studies using *ROSAT* were the first to report the detection of X-ray emissions from young brown dwarfs (Neuhäuser & Comerón 1998[142]; Neuhäuser et al. 1999[143]) in the Chamaeleon, Taurus, and ρ Ophiuchi regions. The number of X-ray-emitting brown dwarfs drastically increased with the following *Chandra* studies on the Orion Nebula Cluster (Garmire et al. 2000[62]; Feigelson et al. 2002a[54]), ρ Ophiuchi (Imanishi et al. 2001b[92]), and IC 348 (Preibisch & Zinnecker 2001[155]; 2002a[156]). At present, X-ray observations in near-by star forming regions have reached the brown dwarf limit.

When observed with a deep exposure, X-ray observations on brown dwarfs are no longer a mere detection experiment. Imanishi et al. (2001b)[92] studied the X-ray properties of four X-ray-brightest brown dwarfs and their candidates in the ρ Ophiuchi cloud with a ∼100 ks exposure with *Chandra*. They found that brown dwarfs have similar X-ray properties with low-mass YSOs with: (1) the temperature of ∼1–2.5 keV, (2) flux variability from some sources, and (3) $L_{\rm X}/L_{\rm bol} = 10^{-3}$–$10^{-5}$. Based on this similarity in X-ray properties, they discussed that young brown dwarfs have the same X-ray emission mechanism with low mass YSOs.

Interestingly enough, X-ray emissions are detected from only a few field brown dwarfs (Fleming, Giampapa, & Schmitt 2000[57]; Rutledge et al. 2000[167]) in spite of their much closer distance to us. It appears that brown dwarfs decrease their X-ray activity as they evolve. This behavior is commonly seen in low mass stars, which are X-ray active in the pre-main-sequence stage and become inactive as they evolve into the main sequence stage. However, as brown dwarfs have no clear definitions of pre-main-sequence and main sequence with no nuclear fusion in their life, it is not quite clear why their X-ray activity changes in the same manner as low mass stars.

2.2.4 Intermediate-mass YSOs

Intermediate-mass YSOs in the mass range of 2.0–10.0 $M_\odot$ are still not clear to have any X-ray emissions. Zinnecker & Preibisch (1994)[202] studied 21 Herbig Ae/Be sources (intermediate-mass counterpart of TTSs) with *ROSAT* and found that 11 of them emit X-rays. Hamaguchi (2000)[77] obtained a similar X-ray detection rate with the *ASCA* data. He also discussed, based on the hard X-ray spectra and flare-like variability of some Herbig Ae/Be stars, that these sources also have magnetic activity similar to low-mass YSOs. Zinnecker & Preibisch (1994)[202], on the contrary, raised a possibility that the X-ray emissions from these sources can be from their low-mass companion and not intrinsically from the intermediate-mass YSOs. Feigelson et al. (2002a)[54] studied the ONC sources with A or B spectral types and discussed that the X-ray emissions can be fully attributable to their low-mass companions. Their samples, however, are not confined to pre-main-sequence sources and may include main sequence intermediate-mass sources. This can mislead the conclusion because intermediate-mass main sequence sources do not emit X-rays at all due to their lack of convection zones.

2.2.5 High-mass YSOs

Some *Chandra* observations detected X-ray emissions from ultra-compact H_{II} regions in massive star-forming regions at a few kpc away. The ultra-compact H_{II} regions are due to the UV photoionization from the central young massive objects of >10 $M_\odot$ and are traced using the free-free emissions in the centimeter wavelengths. The X-ray emissions were reported from the infrared source "n" in the BN/KL region in OMC-1 at 450 pc (Garmire et al. 2000[62]), IRS 1, IRS 2, IRS 3, and a_s in the Monoceros R2 cloud at 830 pc (Kohno et al. 2002[106]; Preibisch et al. 2002b[157]), and IRS 2, IRS 2a, and IRS 3a in the W 3 complex at 2.3 kpc (Hofner et al. 2002[90]).

The most notable characteristics of these X-ray emissions are their hard spectra of $k_{\rm B}T >$ 2 keV and X-ray variability. These features are similar to those of low-mass YSOs, indicating that these X-ray emissions are also of magnetic origin. On the contrary, the $L_{\rm X}/L_{\rm bol}$ values of some high mass YSOs are in the range of 10^{-6}–10^{-8} (e.g., Kohno et al. 2002[106]), which are too small for an empirical value of 10^{-3}–10^{-5} for low-mass YSOs and are comparable to high-mass main sequence sources that emit X-rays of stellar wind origin. The X-ray emission mechanism of these sources is not still clear.

Whatever the X-ray emission mechanism of high-mass YSOs, their hard X-ray emissions provide us with a unique tool to search for young massive stars in dense molecular clouds. The penetrating power of hard X-rays is comparable to radio bands. Moreover, unlike millimeter and centimeter emissions, we can trace the central star with their X-ray emissions. Takagi et al. (2002)[178] illustrated this unique capability by detecting highly absorbed ($N_{\rm H}$ = 1–4×10^{23} cm^{-2}) hard ($k_{\rm B}T$ = 5–10 keV) X-ray emissions from ultra-compact $\rm H_{II}$ regions in the Sagittarius B2 cloud at the distance of 8.5 kpc away.

2.2.6 Diffuse X-ray Emissions

High sensitivity observations using *Chandra* and *XMM-Newton* extended the X-ray studies to weaker but new X-ray emissions. One example is HH objects. In strong adiabatic shocks, the plasma temperature of T is expected at the shock velocity of v_s (Raga, Noriega-Crespo, Velázquez 2002[158]), where

$$T = 1.5 \times 10^5 \left(\frac{v_{\rm s}}{100 \text{ km s}^{-1}}\right)^2 \text{ [K]}. \tag{2.3}$$

With the typical velocity of optical jets (200–300 km s^{-1}), we can expect soft X-ray emissions from HH objects.

Pravdo et al. (2001)[153] reported the detection of X-ray emission from HH-2 using *Chandra*. The emission has a very soft spectrum with an extended spatial distribution. They proposed that this emission originates from the plasma with the temperature of $\sim$10^6 K and the plasma is generated by collision of the protostellar jet with ISM at the speed of $\sim$250 km s^{-1}. A similar soft X-ray emission was later reported by Favata et al. (2002)[49] from L1551 IRS5 using *XMM-Newton*.

Another example of new findings is the diffuse X-ray emissions from massive star-forming regions. Moffat et al. (2002)[132] reported the detection of the diffuse emission with the extent of $\sim 2'$ (4 pc) from the center of NGC 3603. The emission is of thermal origin with $k_{\rm B}T \sim 3.1$ keV and the total luminosity of $L_{\rm X} \sim 2 \times 10^{34}$ ergs s^{-1}. They probably arise from merging or colliding hot stellar winds of massive stars. Similar diffuse thermal emissions were also found by Yusef-Zadeh et al. (2002)[198] in the Arches Cluster and by Rauw et al. (2002)[159] in the Lagoon Nebula (M8). Wolk et al. (2002)[195] reported, on the other hand, that the diffuse X-ray emission found from RCW 38 is of non-thermal origin. They insisted that this emission might be attributable to synchrotron emissions from hidden supernova

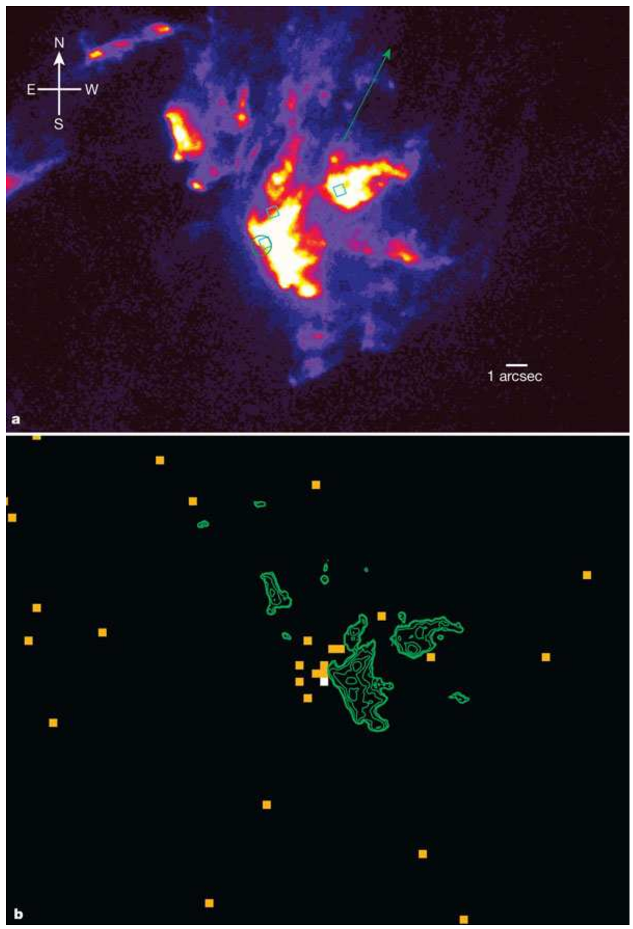


Figure 2.9: Soft X-ray detection from a HH object (HH-2) with *Chandra* (Pravdo et al. 2001[153]). The upper panel gives the *Hubble Space Telescope* image with the position of the *Chandra* and VLA sources with a circle and squares, respectively. The lower panel gives the X-ray event map of *Chandra* with the contours giving the H_α intensity.

remnants. These diffuse X-ray emissions from star-forming regions are restricted to massive star-forming regions at the distance of a few kpc, which makes it difficult to assess the contamination by a complex of point sources. Further studies are necessary to establish the diffuse X-ray emissions from star-forming regions.

2.2.7 Problems to be Addressed

Among a number of problems raised on X-ray emissions from YSOs, we try to give an answer for the following issues in this thesis.

(1) Are there any class I protostars or brown dwarfs in our field? The X-ray emitting samples of these sources are still limited, so it is important to enrich them for further investigations. The combination of our sensitive X-ray and NIR observations is quite useful for this purpose.

(2) Do intermediate-mass YSOs intrinsically emit X-rays? The possible contamination by low-mass binary companions or main sequence intermediate-mass sources prevented previous studies from drawing a definite conclusion. We confine the sample to ostensively single, pre-main-sequence intermediate-mass sources, and compare the binary rate and the X-ray detection rate of these well-defined intermediate-mass YSO samples to address whether they are intrinsic X-ray emitters.

(3) Are there any differences in averaged X-ray features among sources in different mass ranges? Using large number of samples, we can first compare the averaged features of YSOs spanning from intermediate-mass to brown dwarfs in a single star-forming region.

(4) What are the X-ray emission mechanisms of all these samples? Skinner & Walter (1998)[173] raised a possibility that the X-ray emission mechanism of YSOs can be the combination of different mechanisms that is represented by a bimodal structure in plasma temperatures. Their result alone is not conclusive, because their source (SU Aur) is a spectroscopic binary (Bouvier et al. 1986[25]) that can mimic two-temperature plasma. Two things are required. First, we need to try multi-temperature plasma models for the X-ray spectral fittings. All previous *Chandra* papers present only the result of one-temperature fittings. Second, we need to deal with a large number of samples collectively to statistically discuss whether the bimodal temperature structure is attributable to the binarity or not.

(5) What is the nature of X-ray sources that have no NIR counterpart? We conduct

multi-wavelength follow-up observations to reveal their nature.

Chapter 3

Orion Molecular Cloud 2 and 3

Contents

In this chapter, we review the past observations of OMC-2 and OMC-3. In Sect. 3.1, we summarize the global properties of OMC-2 and OMC-3 and illustrate that these regions are one of the best targets for our study. In the following sections, past observations (Table 3.3) in the radio (Sect. 3.2), NIR (Sect. 3.3), optical (Sect. 3.4), and X-ray (Sect. 3.5) band are reviewed. In particular, we stress how to utilize these results in this thesis and how to complement them with our new observations.

3.1 Global Properties

The Orion molecular cloud (OMC) is the cradle of star formation studies. Its proximity to us and the brightness have been attracting many astronomers and astrophysicist who intend to understand the formation of stars in a wide mass range. OMC is located 15–20° south to the Galactic plane (Fig. 3.1) and is comprised of a large complex of molecular clouds (Fig. 3.2). Among them, the Orion A, Orion B, Northern Filament, and λ Ori clouds are physically associated with each other, judging from the fact that they are located at similar distances and have similar velocities. The association of these clouds with Monoceros R2 and Southern Filament is not certain yet (Maddalena et al. 1986[120]).

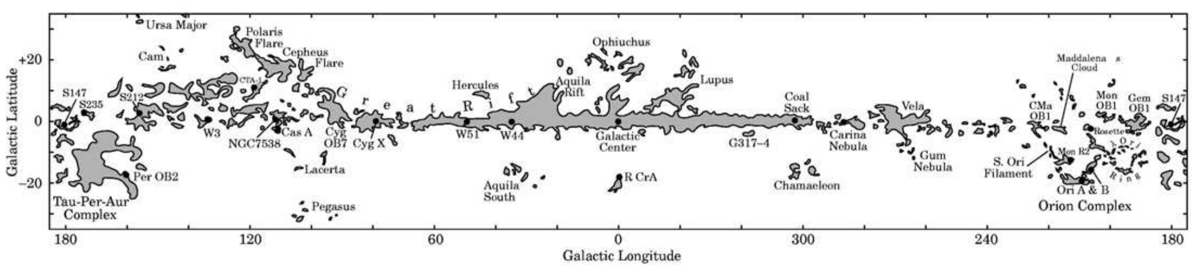


Figure 3.1: CO intensity map along the Galactic plane (Dame, Hartmann, & Thaddeus 2001[44]). The Orion A (Ori. A) molecular cloud, at the right side of the panel, is located at the Galactic longitude and latitude of $l \sim 210°$ and $b \sim -20°$.

OMC-1 is the most prominent and densest cloud core in the Orion A. It contains a young stellar cluster called the Orion Nebula Cluster (ONC) that embraces an OB association known as the Trapezium. The Trapezium is the powering source of the spectacular H_{II} region of the Orion (M42; Fig. 1.1). OMC-2 and OMC-3 are located $\sim$12′ and $\sim$23′ north to ONC. OMC-2 was first discovered by chance as a complex of NIR and molecular emissions

(Gatley et al. 1974[63]). OMC-3 is another intensity peak of radio emissions at the north of OMC-2.

Successive studies in OMC-2 revealed that it is a high density as well as high column density core with the kinematic temperature (24 K) much lower than that of OMC-1 (70 K). OMC-3 has even lower kinematic temperature of 19 K (Castets & Langer 1995[30]). These indicate that both OMC-2 and OMC-3 are the sites of on-going star formation. The star formation activity first occurred in OMC-1 and propagated into the north to OMC-2 and OMC-3, which are currently active star-forming clouds. To illustrate this, these clouds are full of phenomena related to star formations; jets, molecular outflows, HH objects, accretion disks, dust emissions, etc.

Throughout this thesis, we adopt the age of OMC-2 and OMC-3 to be ∼1 Myr, which is estimated based either on the H-R diagram of some bright NIR sources in OMC-2 (Johnson et al. 1990[96]), the K-band luminosity function of cluster members (Ali & DePoy 1995[2]), direct spectroscopic evidence (Hodapp & Deane 1993[87]; Hillenbrand 1997[85]), the presence of rich outflows and NIR excesses commonly seen in YSOs (Strom, Strom, & Merrill 1993[175]; Chen & Tokunaga 1994[34]; Jones et al. 1994[101]), and the I- and V-band magnitudes of member sources (Rebull et al. 2000[160]). The distance to OMC-2 and OMC-3 is assumed to be ∼450 pc (Genzel & Stutzki 1989[66]).

3.2 Radio Observations

3.2.1 Identifying Protostellar Cores with Dust Emissions

Radio continuum observations in OMC-2 and OMC-3 have been conducted at 1.3 mm (Mezger et al. 1990[129]; Chini et al. 1997[35]), 850 μm, 450 μm (Johnstone & Bally 1999[99]), and 350 μm (Lis et al. 1998[117]). All these observations aim to detect dust continuum emissions of protostellar cores and to determine the position and distribution of contracting cores, which pilots the star formation studies in these regions.

Mezger, Zylka, & Wink (1990)[129] obtained a 1.3 mm map of OMC-1 and OMC-2 with the angular resolution of $\sim 90''$ and $\sim 11''$, respectively. In OMC-2, they identified six millimeter clumps (FIR 1–FIR 6). Chini et al. (1997)[35] took a higher-resolution image of OMC-2 and OMC-3 at 1.3 mm and detected a chain of 21 cores (MMS 1–MMS 10 in OMC-3,

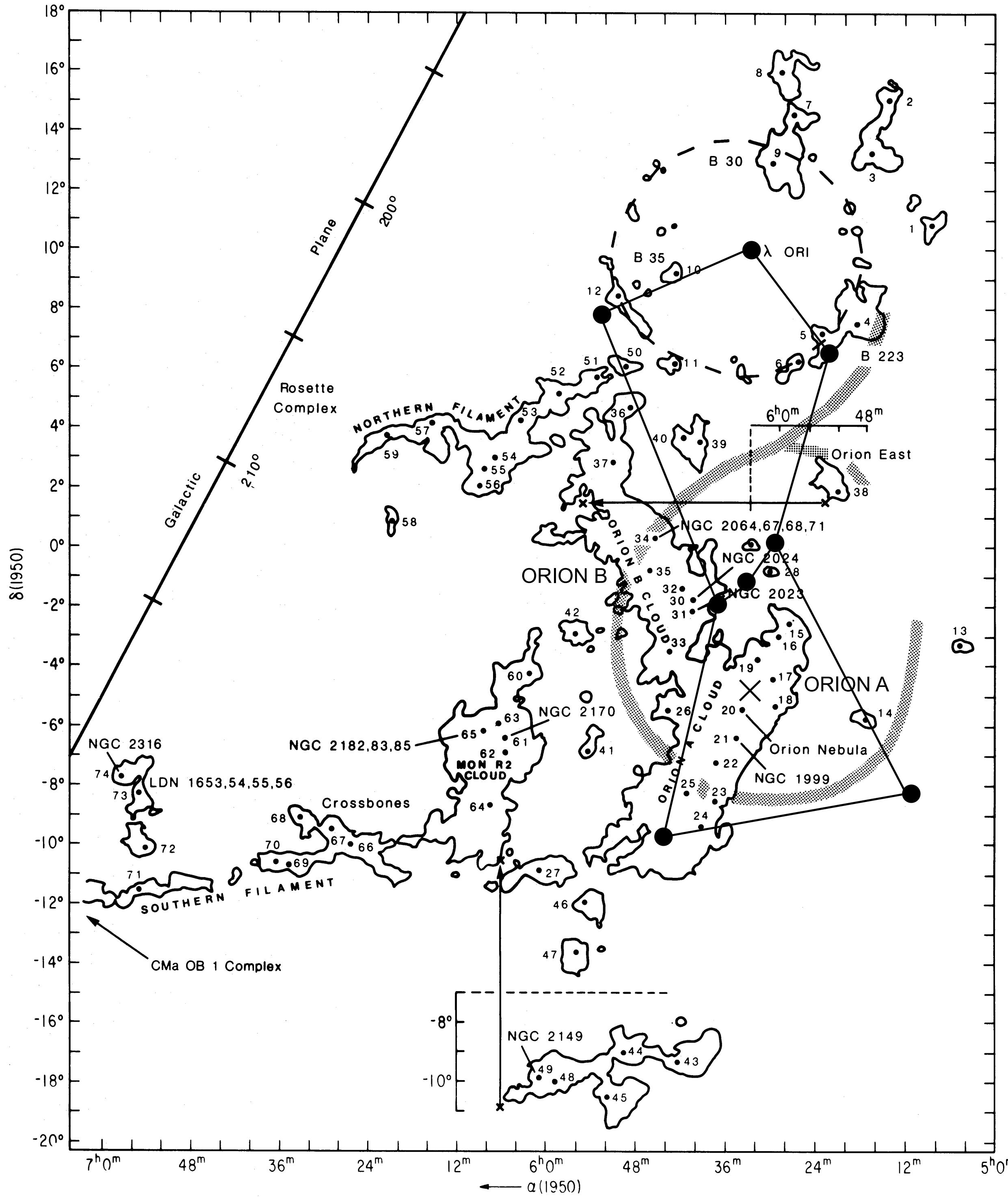


Figure 3.2: Schematic view of the Orion A, Orion B, and Monoceros R2 cloud complex (Maddalena et al. 1986[120]), and major stars of the Orion. OMC-2 and OMC-3 belong to the Orion A, and are located at the position marked with the cross.

FIR 1a–FIR 6d in OMC-2; Table 3.1) using the Institut de Radio Astronomie Millimétrique (IRAM) telescope. Combining the James Clerk Maxwell Telescope (JCMT) photometry from 350 μm to 2 mm and the *Infrared Astronomical Satellite* (*IRAS*) photometry from 12 μm to 100 μm, they showed that most cores have a high ratio of sub-millimeter to bolometric luminosity ($L_{\rm smm}/L_{\rm bol}$), insisting that these cores are class 0 protostar candidates. They also derived the temperature and mass of each condensation (Tables 3.1 and 3.2). Figure 3.3 shows the 1.3 mm intensity map of OMC-2 and OMC-3. Throughout this thesis, we use this map as the landmark of these regions.

Lis et al. (1998)[117] followed with a 350 μm imaging observation with a larger FOV including OMC-2 and OMC-3. They detected 350 μm emissions from almost all the 1.3 mm cores and 10 additional ones (Fig. 3.4).

Johnstone & Bally (1999)[99] observed this region in 450 μm and 850 μm and found similar "integral-shaped" structure identified in 1.3 mm. They constructed the spectral index map, which constrains the dust emissivity and temperature at each position.

3.2.2 Centimeter Observations of Free-free Emissions

Reipurth et al. (1999)[162] observed a $6' \times 15'$ region in OMC-2 and OMC-3 with VLA. The D configuration in 3.6 cm was used, which yields an angular resolution of $\sim 8''$ (Sect. 4.5). In total, 14 sources were detected above $\sim$0.1 mJy, ten of which are associated with the protostellar cores seen in the 1.3 mm continuum. Some 3.6 cm sources are also associated with H_2 (Yu, Bally, & Devine 1997[199]) and CO and HCO^+ (Aso et al. 2000[11]) outflows, HH objects (Reipurth, Bally, & Devine 1997[161]), H_2O masers (Morris & Knapp 1976[135]; Genzel & Downes 1979[65]), and sub-millimeter emissions (Lis et al. 1998[117]). Due to these associations, these 3.6 cm emissions are considered to be free-free emissions from the $H_{\rm II}$ regions generated by shocks from protostar jets.

The detection of thermal free-free centimeter emissions gives strong evidence that there is a class 0 or class I protostar at its origin, with active mass accretion. Seven and four 3.6 cm emissions were respectively found in OMC-2 and OMC-3. In one of the protostellar cores in OMC-3, we found that D configuration observations are not fine enough to correlate with our X-ray and NIR data, so we obtained the A configuration image in this region. The details of analysis and result of this observation are discussed in Sect. 8.4.

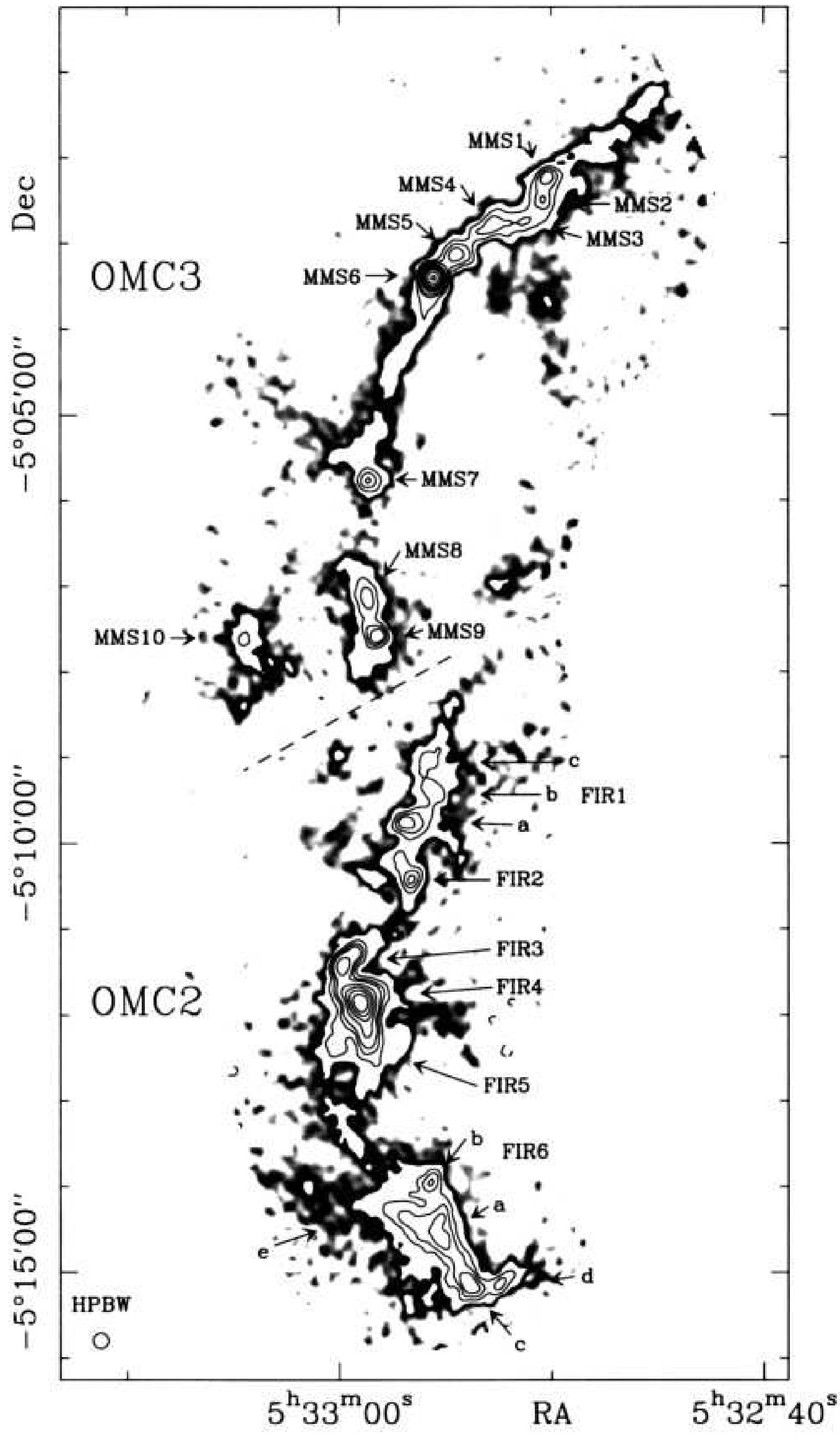


Figure 3.3: Distribution of protostellar cores seen in 1.3 mm (Chini et al. 1997[35]). OMC-2 and OMC-3 are separated by the dashed line into south and north. The equinox is in B1950.0.

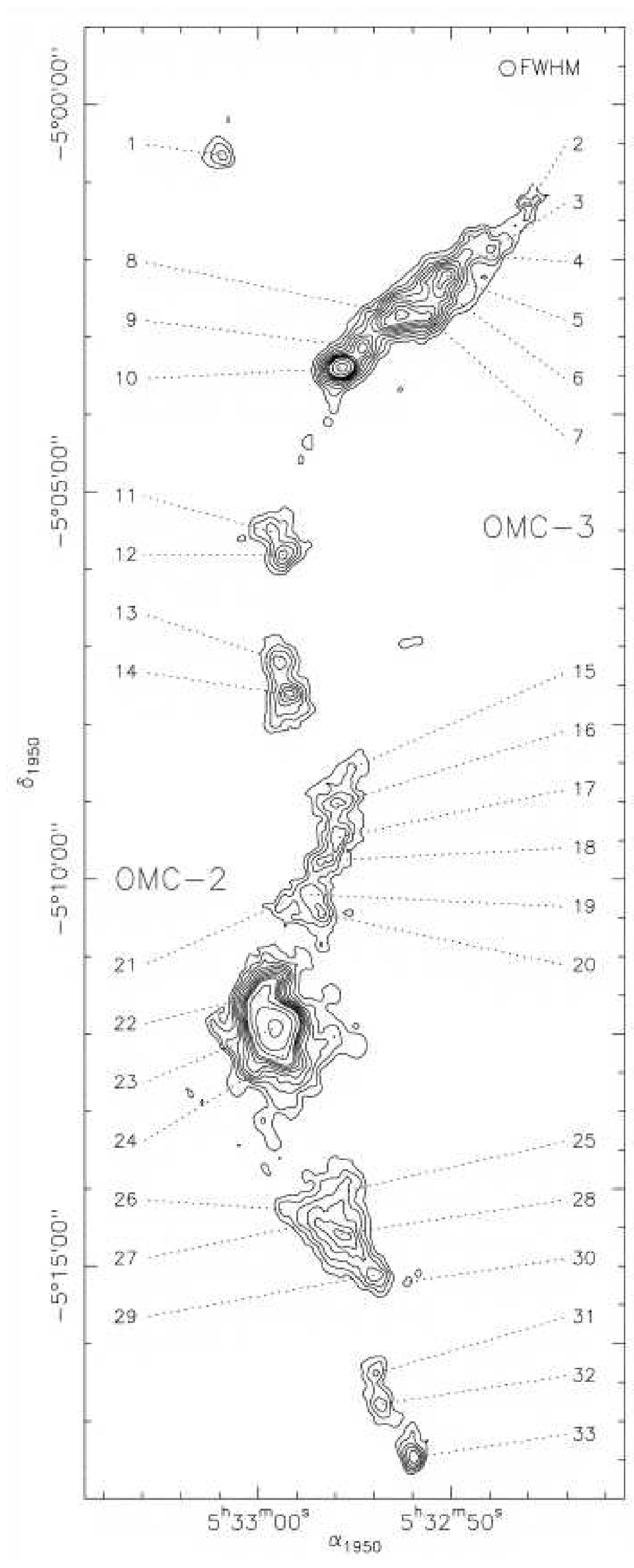


Figure 3.4: Distribution of protostellar cores seen in 350 μm (Lis et al. 1998[117]). The equinox is in B1950.0.

Table 3.1: Properties of protostellar cores in OMC-2 and OMC-3 (1)

object[a]	R.A. (J2000.0)	decl. (J2000.0)	counterpart 350 μm[b]	$H^{13}CO^{+}$[c]	H_2[d]	3.6 cm[e]
			OMC−3			
MMS 1	05:35:18	−05:00:20	CSO 5	....	......	....
MMS 2	05:35:18	−05:00:35	CSO 6	AC 3	flow B	VLA 1
MMS 3[f]	05:35:19	−05:00:51	CSO 7	....	......	....
MMS 4	05:35:20	−05:00:53	CSO 8	....	......	....
MMS 5	05:35:22	−05:01:14	CSO 9	....	flow C	....
MMS 6	05:35:23	−05:01:32	CSO 10	AC 4	flow A	VLA 3
MMS 7	05:35:26	−05:03:53	CSO 12	AC 8	flow F	VLA 4
MMS 8	05:35:27	−05:05:17	CSO 13	....	......	....
MMS 9	05:35:26	−05:05:42	CSO 14	AC 10	flow H	VLA 5
MMS 10	05:35:32	−05:05:42	....	AC 12	......	....
			OMC−2			
FIR 1c	05:35:24	−05:07:10	CSO 16	....	......	VLA 7
FIR 1b	05:35:23	−05:07:32	CSO 17	AC 14	......	....
FIR 1a	05:35:25	−05:07:53	CSO 18	....	......	VLA 8
FIR 2	05:35:24	−05:08:33	CSO 20	AC 15	flow O	....
FIR 3	05:35:28	−05:09:33	CSO 22	....	flow J	VLA 11
FIR 4	05:35:27	−05:10:00	CSO 23	AC 17	......	VLA 12
FIR 5	05:35:26	−05:10:23	CSO 24	....	......	....
FIR 6a	05:35:23	−05:12:36	CSO 28	....	......	VLA 14
FIR 6b	05:35:23	−05:12:03	CSO 25	....	flow L	....
FIR 6c	05:35:21	−05:13:15	CSO 29	....	......	....
FIR 6d	05:35:20	−05:13:15	CSO 30	....	......	....

[a] Nomenclatures follow Chini et al. (1997)[35] for OMC-3 sources and Mezger, Zylka, & Wink (1990)[129] for OMC-2 sources.
[b] Lis et al. (1998)[117].
[c] Aso et al. (2000)[11].
[d] Yu, Bally, & Devine (1997)[199].
[e] Reipurth, Rodrígues, & Chini (1999)[162].
[f] The position of MMS 3 in Chini et al. (1997)[35] is incorrect. The corrected coordinate is given in Tsuboi et al. (2001)[188].

Table 3.2: Properties of protostellar cores in OMC-2 and OMC-3 (2)

object[a]	$L_{\rm bol}$[a]	$L_{\rm bol}/L_{\rm smm}$[a]	T_d[a]	M_{gas}[a]	$\dot{M}$[b]	$\dot{P}$[b]
	($L_\odot$)		(K)	($M_\odot$)	($M_\odot$ yr^{-1})	($M_\odot$ km s^{-1} yr^{-1})
			OMC−3			
MMS 1	<55	<61	20–25	18	....	....
MMS 2	.....	....	.....	.....	7.4×10^{-6}	3.0×10^{-5}
MMS 3	.....	....	.....	.....	ditto[e]	ditto[e]
MMS 4	<56	<72	20–25	11	ditto[e]	ditto[e]
MMS 5	.....	....	.....	.....	6.7×10^{-6}	1.9×10^{-5}
MMS 6	<60	<50	15–25	36	....	....
MMS 7	76	129	26	8	1.4×10^{-5}	8.8×10^{-5}
MMS 8	<89	<178	20	9	....	....
MMS 9	<94	<145	20	10	3.6×10^{-5}	15.1×10^{-5}
MMS 10	.....	....	.....	.....	3.3×10^{-5}	11.3×10^{-5}
			OMC−2			
FIR 1c	128	356	33	5	3.1×10^{-5}	9.8×10^{-5}
FIR 1b	.....	....	.....	.....	ditto[f]	ditto[f]
FIR 1a	<138	....	.....	.....	ditto[f]	ditto[f]
FIR 2	<157	<320	20	8	....	....

[a] Chini et al. (1997)[35].
[b] Aso et al. (2000)[11].
[c] Mass loss rate of the outflows.
[d] Momentum loss rate of the outflows.
[e] MMS 2, MMS 3, and MMS 4 were not resolved in the observations by Aso et al. (2000)[11].
[f] FIR 1a, FIR 1b, and FIR 1c were not resolved in the observations by Aso et al. (2000)[11].

3.2.3 Millimeter Observations of Molecular Cores and Outflows

Mapping observations with various density tracers have been conducted on OMC-2 and OMC-3; with ^{13}CO by Bally et al. (1987)[14], C^{18}O by Dutrey et al. (1993)[45], ^{13}CO, C^{18}O, C^{32}S, and C^{34}S by Castes & Langer (1995)[30], and $C_I(^3P_1–^3P_0)$ and CO ($J = 3 - 2$) by Ikeda et al. (1999)[94]. These observations identified a filamentary structure extending from OMC-1 to OMC-2 and OMC-3. Hundreds of dust condensations along this filament were extracted by two high resolution mapping observations in CS ($J = 1-0$) by Tatematsu et al. (1993)[179] and in NH_3 by Cesaroni & Wilson (1994)[31].

With the Nobeyama 45 m telescope, Aso et al. (2000)[11] observed OMC-2 and OMC-3 in $H^{13}CO^+$ ($J = 1 - 0$), HCO^+ ($J = 1 - 0$), and CO ($J = 1 - 0$) lines, detecting eight molecular outflows in addition to 18 molecular cores. Blue and red lobes of these outflows are well aligned with the outflows seen in the H_2 observation (Yu et al. 1997[199]). The spatial resolution of their map (the beam width of $\sim 15''$) is comparable to those taken in the continuum observations by Chini et al. (1997)[35] and Lis et al. (1998)[117], which enabled them to determine the origin, velocity, mass loss rate, and outflow momentum rate of these outflows. No interferometer millimeter observation has been conducted on these regions.

3.2.4 Polarization Measurements

Matthews & Wilson (2000)[124] and Matthews, Wilson, & Fiege (2001)[125] measured the polarization of dust emissions in OMC-3 at 850 μm. They found that the polarization vectors are highly aligned. The field direction in the plane of the sky ($B_\perp$) was found to be perpendicular to the filamentary structure of OMC-3.

3.3 NIR Observations

3.3.1 Bright Discrete Sources in OMC-2

A series of NIR studies in OMC-2 and OMC-3 has been made concentrating on bright discrete sources. Gatley et al. (1974)[63] was the first to observe OMC-2 with 10 broad-bands from 1.6 μm to 20 μm and found five separate components (IRS 1–IRS 5). Based on their lack of the optical counterparts, their SEDs, and the association with CO emissions, they suggested that these NIR sources are protostars. The nature of these emissions was the main issue in the following three papers published once in four years (Thronson & Thompson 1982[182]; Pendleton et al. 1986[149]; Johnson et al. 1990[96]).

Thronson & Thompson (1982)[182] measured the K-band spectra of IRS 3 and IRS 4, and revealed that IRS 4 shows many emission lines of molecular hydrogen. They pointed out that the series of H_2 $v = 2 - 1$ lines are relatively strong compared to $v = 1 - 0$ lines instead of their low absolute flux, hence they are more likely produced by UV pumping. This was the first candidate of UV-pumped H_2 sources, although the shock pumping origin of these H_2 emission lines could not be ruled out.

Pendleton et al. (1986)[149] conducted NIR–FIR imaging and NIR polarimetry observations of OMC-2. They showed that IRS 1 and IRS 4 are reflection nebulae based on their high degree of polarization. Their FIR map revealed compact sources at IRS 1 and IRS 4, which are considered to be responsible for the reflection.

The final conclusion on the nature of these NIR sources was derived by Johnson et al. (1990)[96], who compiled the SEDs of 11 discrete sources including IRS 1–IRS 5. Most of these SEDs are accountable with the combination of NIR emissions from reddened stellar photosphere and thermal emissions from circumstellar matter, indicating the pre-main-sequence nature of these sources. By de-reddening these SEDs, they also estimated the temperature, extinction, mass and age of these sources. The picture of OMC-2 was thus established to be a cluster of low-luminosity, low- to intermediate-mass YSOs with ages of $\sim$1 Myr embedded in an extended dust cloud.

3.3.2 Surveys with Broad-band Imaging Observations

In 1990's, the advent of large format NIR arrays prompted survey studies of star-forming clouds. Jones et al. (1994)[101] measured the J-, H-, and K-band magnitude of 219 sources in a $15' \times 5'$ region containing OMC-2. The 90% completeness limit of their survey was $K \sim 14$ mag. They found that almost all the sources brighter than 13 mag are cloud members and some sources show flux variability.

Ali & DePoy (1995)[2] conducted a similar survey study in a larger field (1472 arcmin^2) with a deeper exposure ($K \sim 14.5$ mag at the 90% completeness) but in the K band alone. They detected 3548 sources and constructed the surface density map of these sources, identifying two density peaks; i.e., the Trapezium and OMC-2. They also studied the K-band luminosity function (KLF) of these peaks and found that KLF of the Trapezium is consistent with a Miller & Scalo (1979)[131] initial mass function (IMF) with the age of $\sim10^6$ yr. No age estimate of OMC-2 was available due to the paucity of sources ($\sim$33).

Carpenter (2000)[26] and Carpenter, Hillenbrand, & Skrutskie (2001b)[28] extended the survey study of the Orion A cloud with 2MASS data, which provides the J-, H-, and K_s-band magnitude down to 16.0, 15.4, and 14.8 mag at the 93.8% completeness, respectively. They showed that there is an enhanced stellar surface density over a $0.4^\circ\times2.4^\circ$ region containing the Trapezium, OMC-2 and OMC-3.

One of the unique outcomes of 2MASS project was presented by Carpenter, Hillenbrand, & Skrutskie (2001b)[28], who systematically studied the NIR variability properties of pre-main-sequence sources in a $0.84^\circ\times6^\circ$ region across OMC. A total of 1235 variables are classified into several patterns; periodic stars, eclipsing systems, stars steadily increase their brightness, stars that change colors redder or bluer, long-term variables, and so on. Models such as hot or cool spots on the surface, changes in the mass accretion rate, inner disk radius, or extinction are employed to account for their variability. This is the most comprehensive study of NIR variability seen in pre-main-sequence sources.

We found that all above NIR surveys are not deep enough to match our X-ray observation of OMC-2 and OMC-3. Therefore, we conducted the deeper ($K \sim 16$ mag) and more comprehensive NIR observations in these regions (this thesis; Tsujimoto et al. 2003[192]). The details of analysis and result of this observations are discussed in Sect. 5.2.

3.3.3 Narrow-band Studies of Molecular Outflows

The vibrational-rotational transition of $v = 1-0$ S(1) works as an effective coolant of the excited hydrogen molecules, so this emission line is commonly used as a tracer of outflows from protostars (Bally et al. 1993[15]; Hodapp & Ladd 1995[88]). Yu et al. (1997)[199] conducted a narrow-band imaging observation in H_2 $v = 1-0$ S(1) (2.12 μm) and found $\sim$80 sources in OMC-2 and OMC-3. These sources consist a dozen of collimated outflow aligned from east to west and with millimeter cores at their center. Based on the high condensation and level of outflow activity seen in H_2, they described that OMC-2 and OMC-3 are undergoing a “microburst” of star formation.

We further obtained the H_2 image of OMC-3 with a better spatial resolution using the University of Hawaii 88 inch (2.2 m) telescope and with a much higher sensitivity using the Subaru telescope in order to make it easier to identify the powering source of the outflows in crowded regions. The details of analysis and result of these observations are discussed in Sects. 8.2 and 8.3.

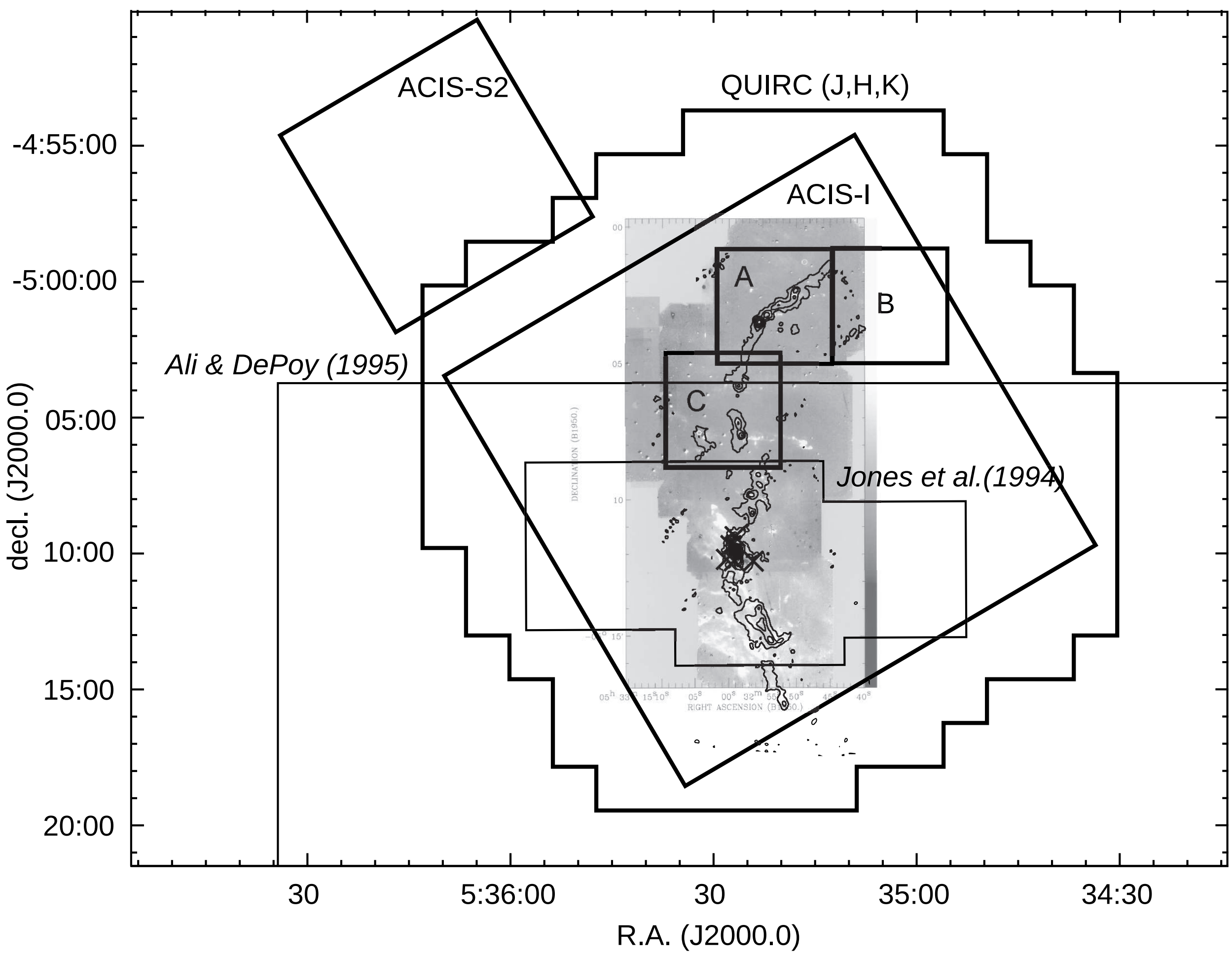


Figure 3.5: Past NIR observations of OMC-2 and OMC-3. The gray scale represents the continuum-subtracted H_2 emission (Yu et al. 1997[199]), while the contours give the 1.3 mm intensity (Chini et al. 1997[35]). The pluses are the position of IRS 1–IRS 5 (Johnson et al. 1990[96]). The survey fields of Jones et al. (1994)[101] and Ali & DePoy (1995)[2] are shown with thin lines. The FOVs of observations conducted in this thesis; *Chandra*, QUIRC *J*, *H*, and *K* band, and QUIRC H_2 band (A, B, and C) are shown with thick lines (see Sect. 8.2).

3.4 Optical Observations

3.4.1 Membership Studies with Proper Motion Measurements

Membership studies of OMC have been conducted mainly in the optical wavelength through the measurement of proper motions (Parenago 1954[148]; Jones & Walker 1988[100]; van Altena et al. 1988[193]; McNamara et al. 1989[126]). Tian et al. (1996)[183] examined the relative proper motions and the membership probabilities for 333 sources with the plates taken over a period of 83 years in the Shanghai Observatory. 184 sources have higher membership probability than 70%, which they consider are the Orion members.

3.4.2 Detection of Circumstellar Disks

Two optical techniques were employed to identify sources with circumstellar disks in these regions; one is by H_α emission and the other by UV excess (Fig. 3.6).

Herbig & Bell (1988)[83] compiled a catalog of 735 pre-main-sequence sources in OMC, which were observed with the slit spectrograph or at equivalent resolution. The equivalent width of H_α emission ($EW[H_\alpha]$) was measured for these sources, with which they were classified into cTTSs ($EW[H_\alpha] > 10$ Å) and wTTSs ($EW[H_\alpha] < 10$ Å).

Rebull et al. (2000)[160] surveyed the ONC flanking fields covering OMC-2 and OMC-3 with the U, V, and I bands. The purpose of this work is to determine the $(U{-}V)$ and $(V{-}I)$ colors of the Orion sources and to pick up sources with circumstellar disks based on their UV excess emissions. About 5000 sources were examined for their colors and ∼230 of them were found to have a UV excess of more magnitudes than 0.5 mag, a signature of circumstellar disk. They also illustrated the consistency of this method to search for disks with other better-established methods such as NIR excess or H_α emission.

In this thesis, we combine the NIR excess (this thesis; Tsujimoto et al. 2003[192]), UV excess (Rebull et al. 2000[160]), and H_α emission line (Herbig & Bell 1988[83]) data to discriminate sources with and without circumstellar disks. We recognize sources to have circumstellar disks if they have at least one positive detection among these three methods.

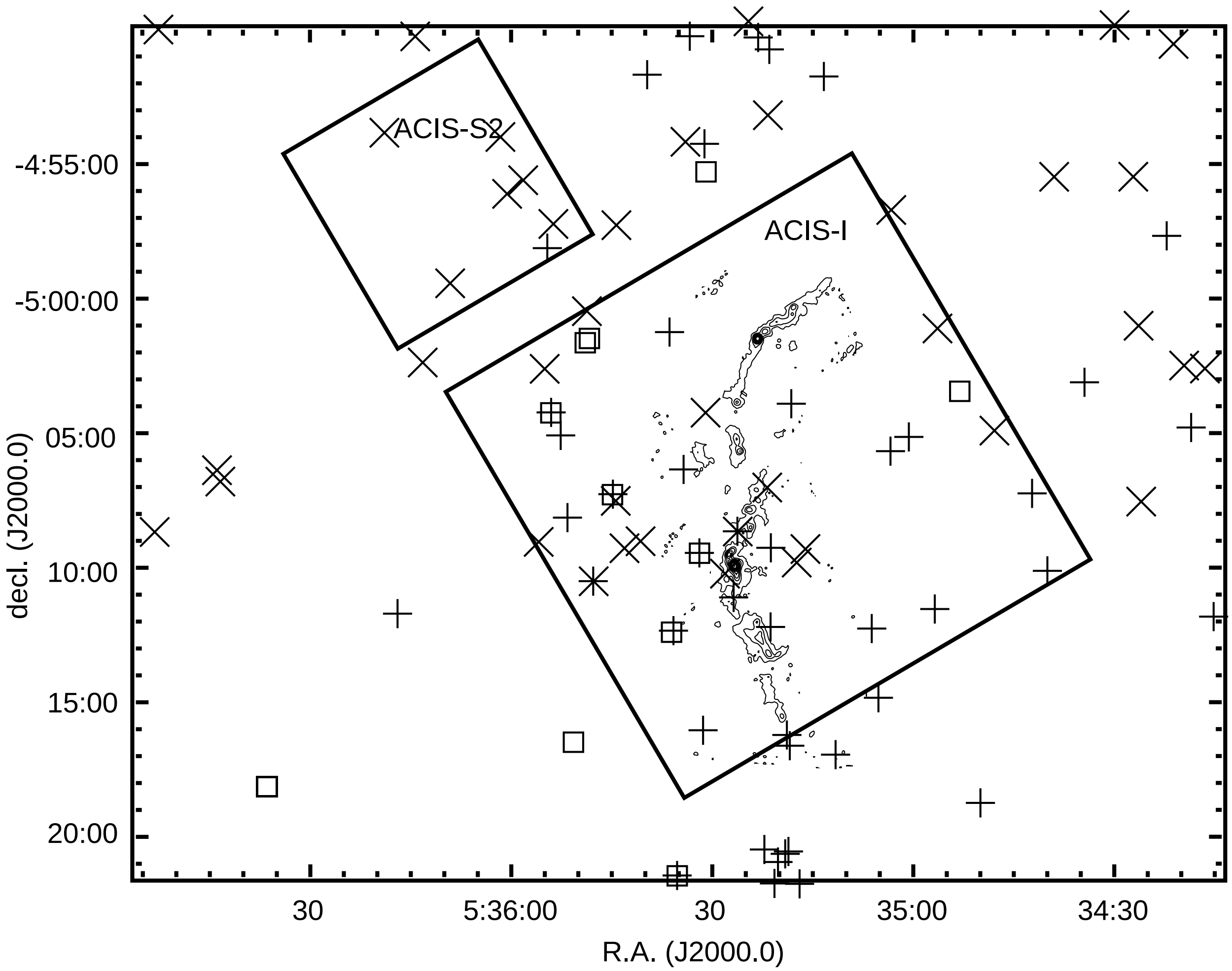


Figure 3.6: Past optical observations of OMC-2 and OMC-3. *squares*: sources with the equivalent width of H_α larger than 10 Å (Herbig & Bell 1988[83]). *Crosses*: sources with *UV* excess (Rebull et al. 2000[160]). *Pluses*: sources with the membership probability of more than 70% (Tian et al. 1996[183]). The contours represent the 1.3 mm intensity (Chini et al. 1997[35]). The *Chandra* FOVs are given in two oblique squares.

3.5 X-ray Observations

3.5.1 Surveys in the Soft X-ray Band

Two systematic X-ray studies of the Orion A were conducted by Gagné & Caillault (1994)[60] and Gagné, Caillault, & Stauffer (1995)[61] with Imaging Proportional Counter (IPC) onboard the *Einstein* satellite and with High-Resolution Imager (HRI) onboard *ROSAT* (Fig. 3.7). These observations provide the position and the count rate of 245 and 389 X-ray sources in a 4.5 square degree and a 0.8 square degree region, respectively. No spectroscopy has been made in these observations. In Gagné, Caillault, & Stauffer (1995)[61], X-ray emission was found from sources of all spectral types, ranging from massive O- and B-type stars to late-type pre-main-sequence sources. About 75 X-ray sources with a measured spectral type were investigated for any relations between the X-ray luminosity ($L_{\rm X}$ with an assumed X-ray spectrum) and the bolometric luminosity (L_{bol}), $v \sin i$, rotation period, and the effective temperature. They found that (1) $L_{\rm X}$ is related to L_{bol} with $L_{\rm X}/L_{\rm bol} < 10^{-3}$ and (2) $L_{\rm X}$ and $L_{\rm X}/L_{bol}$ do not appear to be related with the stellar rotation.

Geier, Wendker, & Wisotzki (1995)[64] observed the Orion region with Position-Sensitive Proportional Counter (PSPC) onboard *ROSAT* and detected 171 X-ray sources. The purpose of this observation was to confirm the diffuse X-ray emission reported by Ku & Chanan (1979)[109], but the trial was in vain because X-ray sources were too crowded. Unlike previous two detectors, PSPC has low-resolution spectral capability in the soft X-ray band (0.1–2.4 keV), so they also conducted spectral analyses for bright 95 sources to determine the plasma temperature and the interstellar absorption. Sources were separated into two groups based on these parameters: (1) those with the absorption of $N_{\rm H} < 2 \times 10^{21}$ cm^{-2} and the temperature of $T \sim 4 \times 10^6$ K (0.35 keV) and (2) those with the absorption of $N_{\rm H} \sim 10^{22}$ cm^{-2} and the temperature of $T \sim 10^7$ K (0.86 keV). The sources in the former group are randomly distributed, while the latter are concentrated on the Trapezium cluster. These results lead them to suspect that the former group consists of main sequence sources, while the latter consists of YSOs.

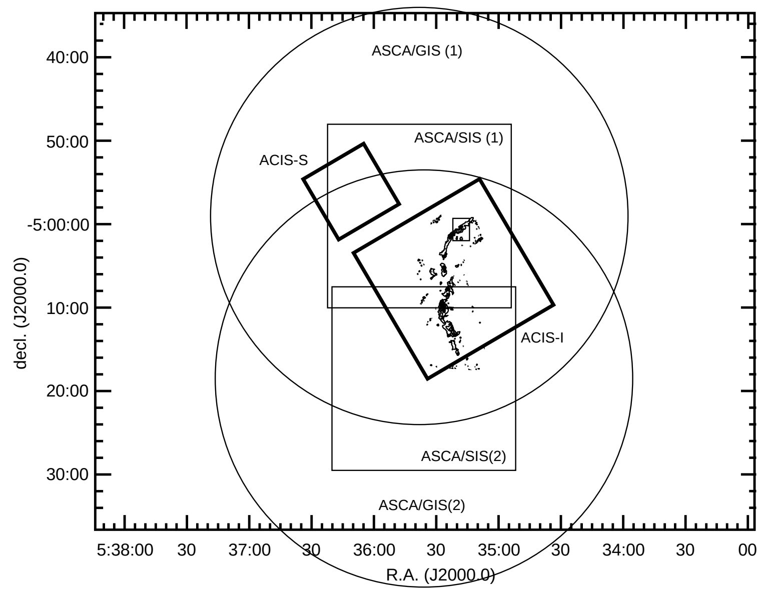
3.5.2 Surveys in the Hard X-ray Band

Yamauchi et al. (1996)[196] studied the whole Orion A cloud with two pointing observations with the *ASCA* satellite (Fig. 3.7). They detected 52 X-ray sources using Gas Imaging

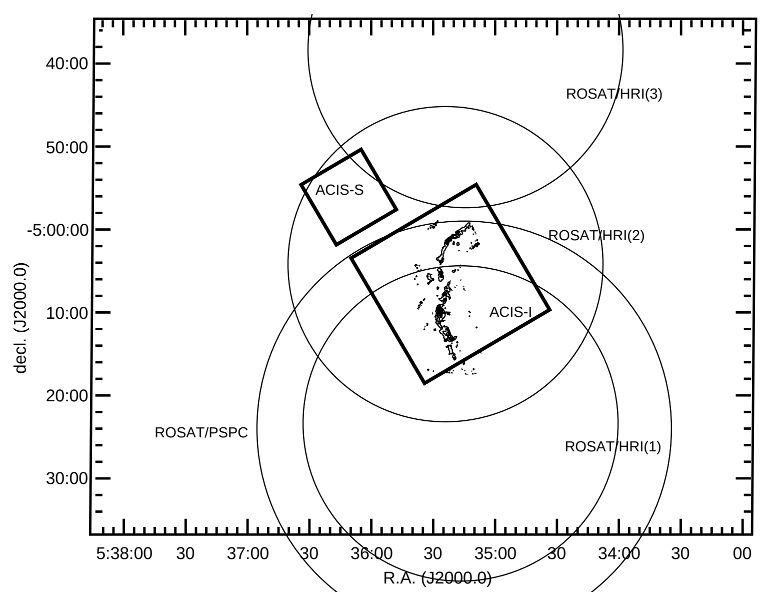
Spectrometer (GIS) and Solid-state Imaging Spectrometer (SIS). They also derived spectral parameters of five bright X-ray sources. An interesting result was obtained that these X-ray sources have a thin-thermal plasma of two temperatures; one is 0.7–1 keV and the other is 3–5 keV. This was also confirmed in a composite spectra of many discrete sources accumulated in large areas. This was the first hard X-ray imaging observation of the Orion region. In addition to lower temperature plasma of 0.7–1 keV found by *Einstein* and *ROSAT*, higher temperature plasma of 3–5 keV was first confirmed from discrete sources in the Orion.

Two important lessons can be drawn from these previous observations. First, we need to have the spatial resolution of $\sim 1''$ to resolve each X-ray source in star-forming regions like OMC. Second, the hard X-ray capability is inevitable to overcome high extinction and to detect high temperature plasma. Preceding X-ray observatories such as *Einstein*, *ROSAT*, and *ASCA* did not meet either of these requirements. We therefore made the X-ray observation on this region using the *Chandra X-ray Observatory*, which has $\sim 1''$–$5''$ spatial resolution and hard X-ray imaging and spectroscopy capability (Tsujimoto et al. 2002a[189]; this thesis).

The initial result of our *Chandra* observation was published by Tsuboi et al. (2001)[188], who concentrated on the hard X-ray detections from 1.3 mm cores in the northern part of OMC-3. Two hard X-ray sources were detected at the protostellar cores with a thin-thermal spectrum of $N_{\rm H} = 1$–3×10^{23} cm^{-2} and $L_{\rm X} \sim 10^{30}$ ergs s^{-1}. These protostellar cores are at the class 0 stage based on their SEDs. They proposed that these X-ray emissions are the first candidates of X-ray-emitting class 0 sources.



(a) *ASCA*



(b) *ROSAT*

Figure 3.7: FOVs of the X-ray observations on OMC-2 and OMC-3. The past observations are shown in thin circles and squares, while the *Chandra* observation conducted in this thesis is in the thick squares. The contours are the 1.3 mm intensity (Chini et al. 1997[35]). (*a*) The FOVs of two *ASCA* observations (Yamauchi et al. 1996[196]). (*b*) The FOVs of three *ROSAT*/HRI observations (Geier, Wendker, & Wisotzki 1995[64]) and one *ROSAT*/PSPC observations (Gagné et al. 1995[61]).

Table 3.3: Past observations of OMC-2 and OMC-3

	bands/lines	mode	observatory	reference
centimeter ...	3.6 cm	imaging	VLA	Reipurth et al. (1999)[162]
	3.6 cm	imaging	VLA	Tsujimoto et al. (2002c)[191] (this thesis)
millimeter	^{13}CO	imaging	AT&T Bell Lab.	Bally et al. (1987)[14]
	1.3 mm	imaging	IRTF, IRAM	Metzger et al. (1990)[129]
	CS	imaging	NRO	Tatematsu et al. (1993)[179]
	$C^{18}O$	imaging	AT&T Bell Lab.	Dutrey et al. (1993)[45]
	NH_3	imaging	Effelsberg	Cesaroni & Wilson (1994)[31]
	^{13}CO, $C^{18}O$, $C^{32}S$, $C^{34}S$	imaging	SEST	Castets & Langer (1995)[30]
	1.3mm	imaging	IRAM	Chini et al. (1997)[35]
	$H^{13}CO^+$, HCO^+, CO	imaging	NMA	Aso et al. (2000)[11]
	^{12}CO	imaging	NRAO	Yu et al. (2000)[200]
	CO	imaging	BIMA	Williams et al. (2003)[194]
sub-millimeter	40–400μm	imaging	LJO, KAO, MLO	Thronson et al. (1978)[181]
	C I, C II	imaging	KAO	Herrmann et al. (1997)[84]
	350μm	imaging	CSO	Lis et al. (1998)[117]
	450μm, 850μm	imaging	JCMT	Johnstone & Bally (1999)[99]
	C I, CO	imaging	Mt. Fuji	Ikeda et al. (1999)[94]
	850μm	polarimetry	JCMT	Matthews & Wilson (2000)[124]
NIR–MIR	1.6–20μm	imaging	Wilson, Hale	Gatley et al. (1974)[63]
	K	spectroscopy	Steward	Thronson & Thompson (1982)[182]
	1.25–100μm	imaging	IRTF, KAO	Pendleton et al. (1986)[149]
	J, H, K, L	imaging, polarimetry	WIRO	Johnson et al. (1990)[96]
	J, H, K	imaging	KPNO	Jones et al. (1994)[101]
	K	imaging	Perkins	Ali & DePoy (1995)[2]
	H_2	imaging	KPNO, CTIO	Yu et al. (1997)[199]
	J, H, K_s	imaging	2MASS	Carpenter (2000)[26]
	J, H, K_s	imaging	2MASS	Carpenter et al. (2001b)[28]
	J, H, K, L', H_2	imaging	Subaru, IRTF	Tsujimoto et al. (2002b)[190] (this thesis)
	J, H, K, H_2	imaging	UH88	Tsujimoto et al. (2003)[192] (this thesis)
optical	V	imaging	Shanghai	Tian et al. (1996)[183]
	U, V, I	imaging, spectroscopy	KPNO,WIYN	Rebull et al. (2000)[160]
X-ray	0.1–4.0 keV	imaging	*Einstein*	Gagné, & Caillault (1994)[60]
	0.1–2.4 keV	imaging, spectroscopy	*ROSAT*	Geier et al. (1995)[64]
	0.2–2.0 keV	imaging	*ROSAT*	Gagné et al. (1995)[61]

(*cont.*)

bands/lines	mode	observatory	reference
0.5–8.0 keV	imaging, spectroscopy	*ASCA*	Yamauchi et al. (1996)[196]
0.5–8.0 keV	imaging, spectroscopy	*Chandra*	Tsuboi et al. (2001)[188]
			Tsujimoto et al. (2002a)[189] (this thesis)

Table 3.4: Past NIR survey studies of OMC-2 and OMC-3

reference	area ($arcmin^2$)	region	band	completeness limit[a](mag)	num. of sources
Jones et al. (1994)[101]	90	OMC-2	J, H, K	∼14	219
Ali & DePoy (1995)[2]	1472	Trapezium, OMC-2	K	∼14.5	3548
Carpenter (2000)[26]	....	Orion A, B, etc.	J, H, K_s	∼14.8	....
Tsujimoto et al. (2003)[192] (this thesis)	512	OMC-2, OMC-3	J, H, K	∼16.0	1448

[a] The 90% completeness limit in the K band, except for Carpenter (2000)[26] with the 94% limit.

Table 3.5: Past X-ray survey studies of OMC-2 and OMC-3

reference	area ($arcmin^2$)	instrument	band (keV)	num. of sources
Gagné & Caillault (1994)[60]	16000	*Einstein*/IPC	0.1–4.0	245
Gagné et al. (1995)[61]	2900	*ROSAT*/HRI	0.2–2.0	389
Geier et al. (1995)[64]	1900	*ROSAT*/HRI	0.1–2.4	171
Yamauchi et al. (1996)[196]	3900	*ASCA*/GIS, SIS	0.5–8.0	52
Tsujimoto et al. (2002)[189] (this thesis)	360	*Chandra*/ACIS	0.5–8.0	398

Chapter 4

Observing Facilities and Instruments

Contents

In this chapter, we review on all telescopes and instruments that we used for this thesis. The first section is devoted for the *Chandra X-ray observatory* (Sect. 4.1), where we mention the basic features of the spacecraft, optics, and detectors. For detectors, ACIS onboard *Chandra*, which we actually used, is particularly focused. The following three sections are for NIR telescopes and instruments; QUIRC on the University of Hawaii 88 inch (2.2 m) telescope (Sect. 4.2), IRCS on the Subaru telescope (Sect. 4.3), and NSFCam on IRTF (Sect. 4.4). The last section deals with VLA (Sect. 4.5). Along with the properties of the array, the basic idea of radio interferometry is briefly reviewed.

4.1 *Chandra X-ray Observatory*

4.1.1 Spacecraft

The *Chandra X-ray Observatory* was successfully launched in July 1999 by National Aeronautics and Space Administration (NASA). The spacecraft circulates around the earth once in 64 hours on an elliptical orbit with the perigee and the apogee distance of 10000 km and 140000 km, respectively. It is comprised of several modules, including the solar panels, the mirror assembly, the telescope, and the integrated science instrument module (Fig. 4.1). Details of *Chandra* can be found in the *Chandra* Proposers' Observatory Guide (2001)[32].

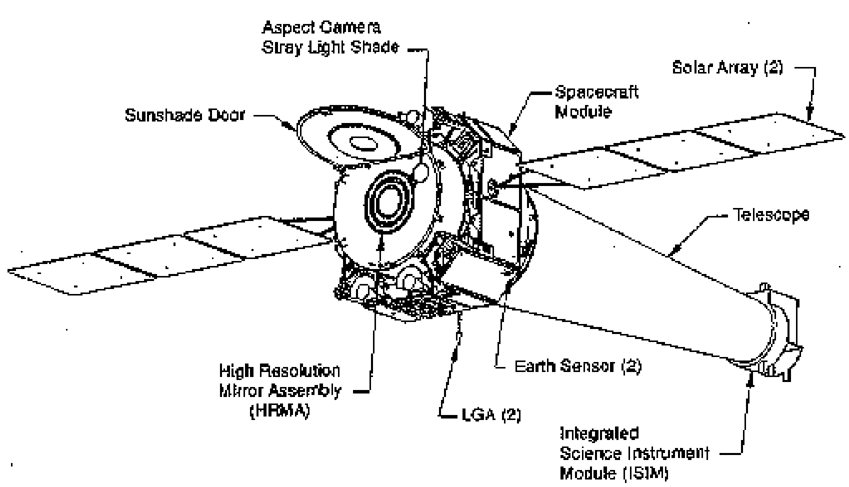


Figure 4.1: Schematic view of the *Chandra* spacecraft (*Chandra* Proposers' Observatory Guide 2001[32]).

The pointing accuracy of the spacecraft, which is measured and stabilized by the Pointing Control and Aspect Determination (PCAD) system, is $\sim 0.1'' \text{ s}^{-1}$. The *Chandra* line-of-site is kept dithered during observations drawing a Lissajous pattern. The dithering distributes photons over many detector pixels for several purposes; to prevent a bad pixel to ruin the entire observation, to reduce uncertainty due to pixel-to-pixel variations in the quantum efficiency, to pick up sources in the gaps between CCD chips, and to allow sub-sampling of the image.

4.1.2 Optics

Mirror Assembly

High Resolution Mirror Assembly (HRMA) is the mirror assembly carried on *Chandra*. It consists of a nested set of four paraboloid-hyperboloid (Wolter-1) grazing-incidence X-ray mirror pairs, with the focal length of ∼10 m and the largest mirror having a diameter of ∼1.2 m (Fig. 4.2). The most precisely shaped and aligned, and the smoothest X-ray mirrors ever constructed allow *Chandra* to obtain unprecedentedly sharp X-ray images.

Effective Area

The on-axis effective area of HRMA is shown in Figure 4.3, together with the expected effective areas when the detector quantum efficiency is convolved. The effective area decreases as the increasing off-axis angle (the vignetting effect; Fig. 4.4) depending on the energy.

Point Spread Function

The point spread function (PSF) is a spatial distribution function over the detector surface of incident X-ray photons at a given energy. The PSF is approximately shaped Gaussian. The sharpness of images is evaluated either by the full width half maximum (FWHM) or the encircled energy radius of PSFs. A circle of 50% encircled energy radius, which equals to the half of FWHM if the PSF is exactly Gaussian, accumulates the 50% of incident photons. Figure 4.5 gives the on-axis PSF, while Figure 4.6 shows the dependence of PSFs on the off-axis angle.

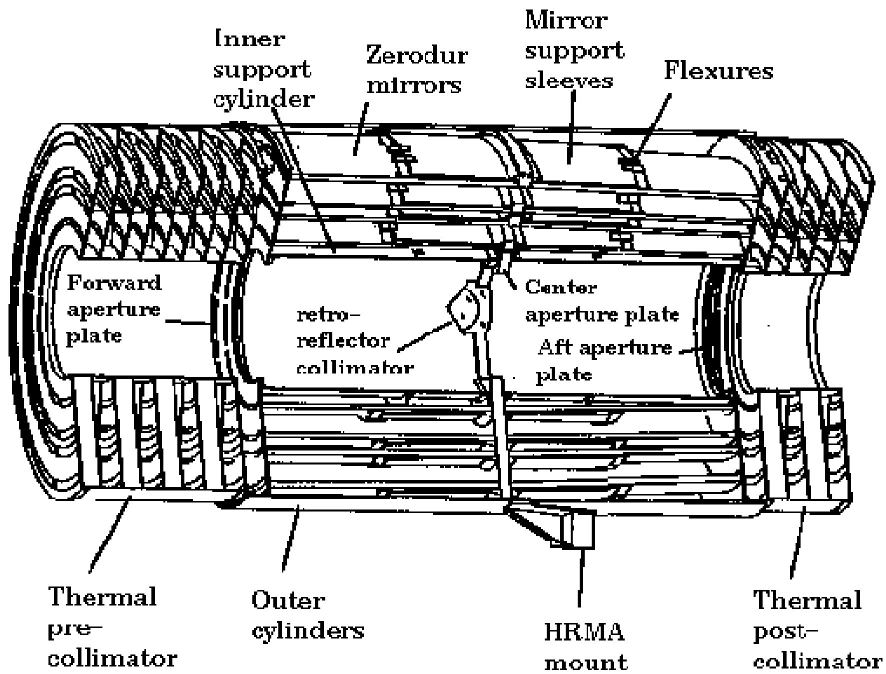


Figure 4.2: Configuration of four nested mirror pairs (*Chandra* Proposers' Observatory Guide 2001[32]).

4.1.3 Instrument (ACIS)

Configuration

Ten X-ray CCDs comprise the Advanced CCD Imaging Spectrometer (ACIS) array (Fig. 4.7). The array consists of two parts; ACIS-I with 2×2 CCDs (ACIS-I0, -I1, -I2, and -I3) and ACIS-S with 1×6 CCDs (ACIS-S0, -S1, -S2, -S3, -S4, and -S5), which are mainly used for the imaging-spectroscopic and the grating-spectroscopic purposes, respectively. The format of CCD chips is 1024×1024 pixels with the pixel scale of 0.492$''$ pixel^{-1}. All CCDs utilize front-illuminated CCDs except for ACIS-S1 and ACIS-S3 that are back-illuminated CCDs.

Energy Resolution

An X-ray CCD can work as a medium-resolution ($E/\Delta E$ = 10–50) spectrometer as well as an imager. When an incident X-ray photon with the energy of E keV is photoelectrically absorbed by silicon atoms in the depletion layer of the device, it emits a photoelectron with the corresponding kinematic energy. The photoelectron keeps ionizing other atoms until

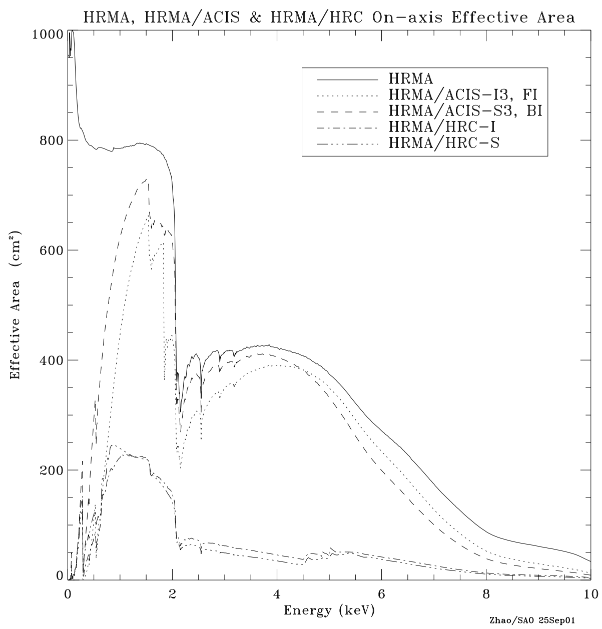


Figure 4.3: Energy dependence of the on-axis effective area of HRMA (*solid curve*). The expected effective areas convolved with the detector quantum efficiency are separately shown with dotted (front-illuminated ACIS), dashed (back-illuminated ACIS), dashed-and-dotted (HRC-I), and dashed-and-triplicated-dotted (HRC-S) curves (*Chandra* Proposers' Observatory Guide 2001[32]).

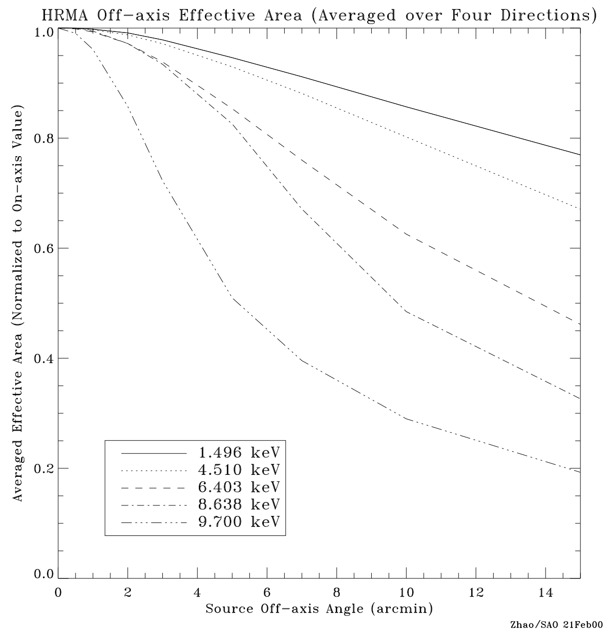


Figure 4.4: Spatial dependence of the effective area as a function of off-axis angle for several representative incident X-ray energies (*Chandra* Proposers' Observatory Guide 2001[32]).

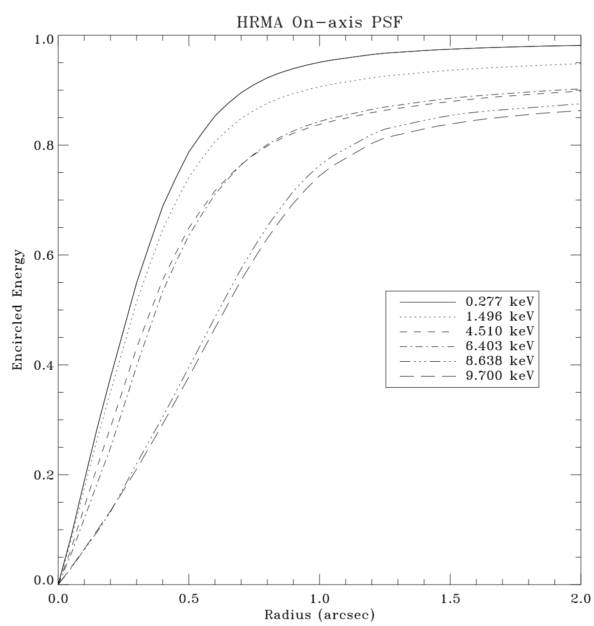


Figure 4.5: On-axis PSFs of several representative incident X-ray energies as a fraction of the encircled photons (energy) inside the circle with a given radius (*Chandra* Proposers' Observatory Guide 2001[32]).

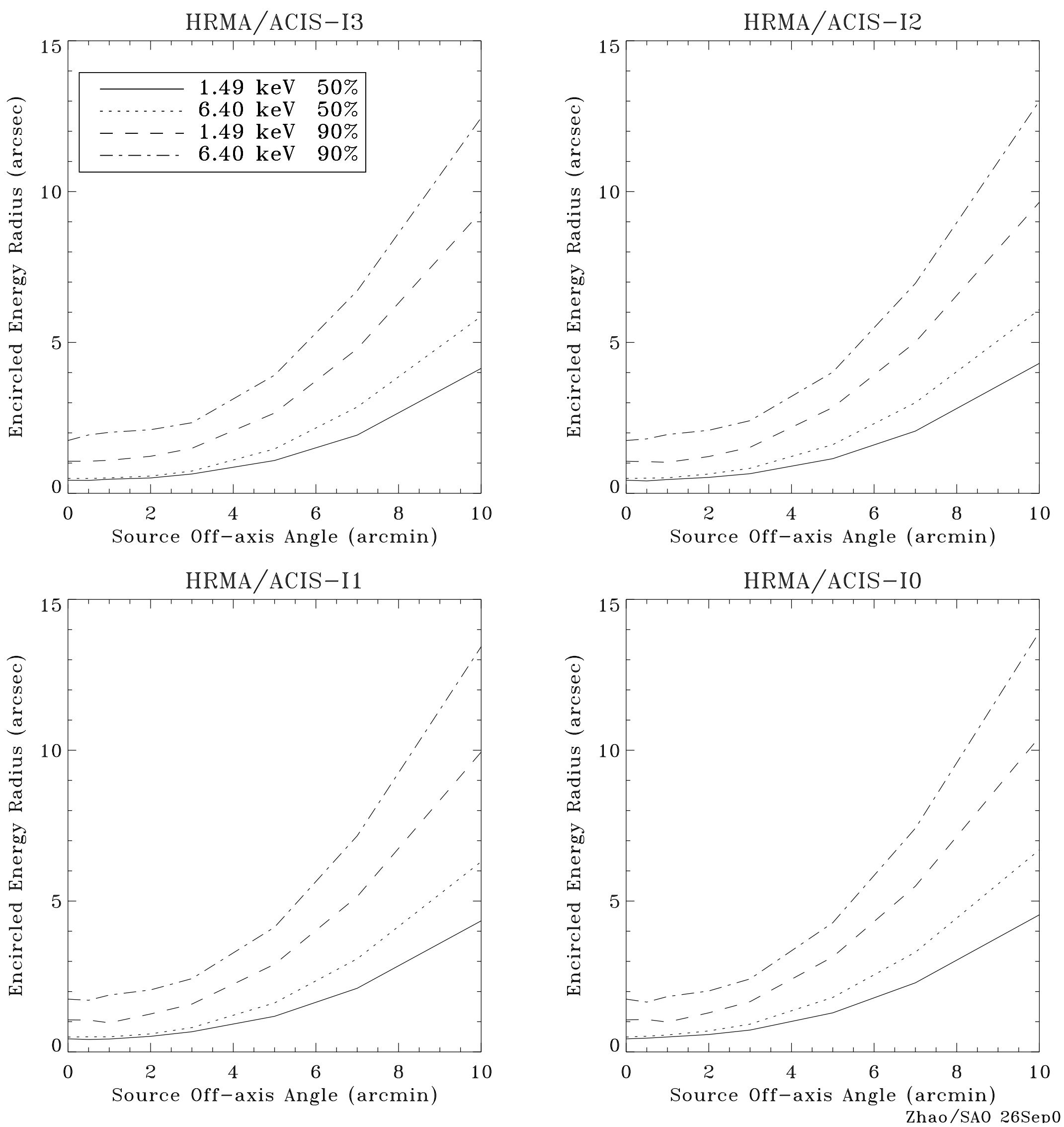


Figure 4.6: The 50% and 90% encircled energy radii at a given off-axis angle. The four panels give the radii of representative incident X-ray energies separately for ACIS-I0, -I1, -I2, and -I3 (*Chandra* Proposers' Observatory Guide 2001[32]).

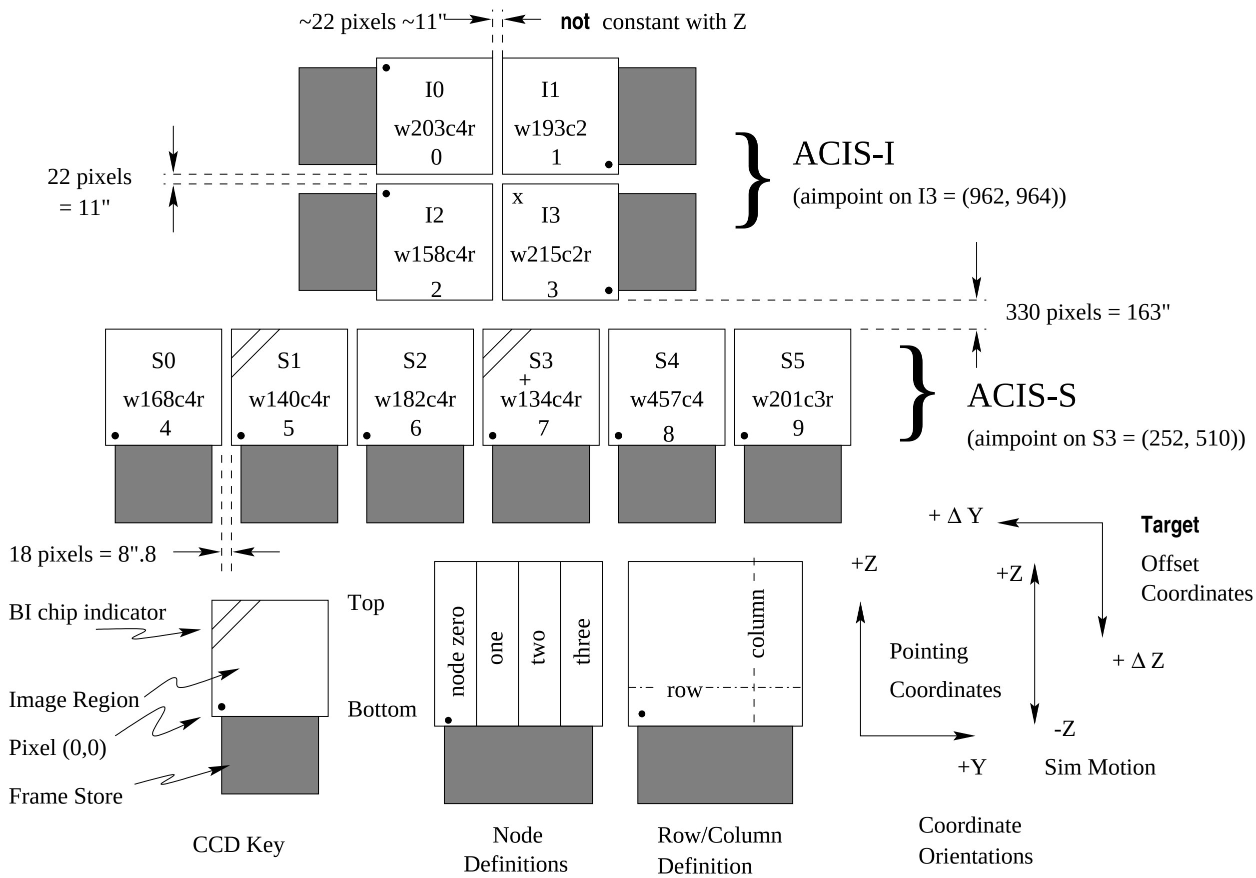


Figure 4.7: Configuration of the ACIS array. The top two panels show the ACIS-I and ACIS-S arrays, while the schematic view of each CCD chip is given at the bottom panel. In our observation, ACIS-I0, I1, I2, I3 and S2 were used with the aim point on the cross at the top left corner of ACIS-I3 (*Chandra* Proposers' Observatory Guide 2001[32]).

all energy is consumed, which finally produces $n = E/W$ electrons where W keV is the average energy required to ionize an atom. By measuring n with the device, the incident X-ray energy is determined on a photon basis. The energy resolution is determined by the convolution of the Poisson fluctuation of n and the readout noise (Fig. 4.8). In case of silicon ($W = 3.65$ keV), the lower limit of the energy resolution is ∼120 eV at 6 keV.

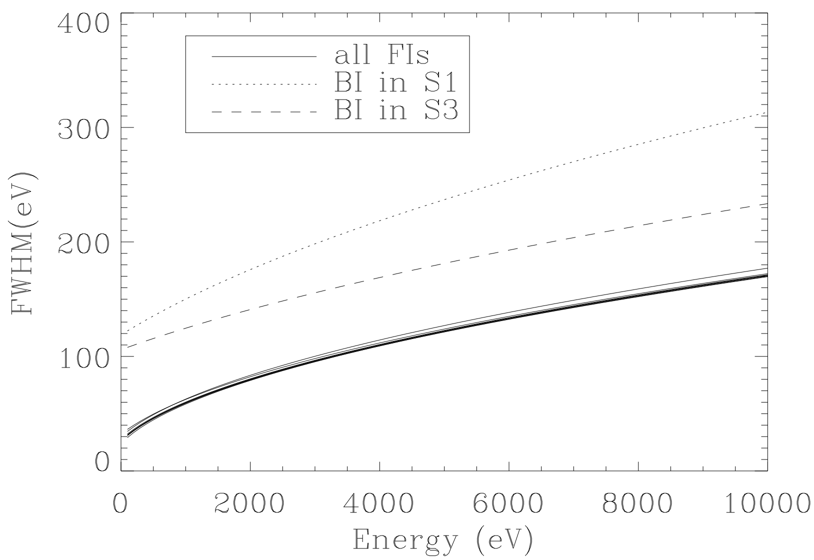


Figure 4.8: Pre-launch energy resolution of front-illuminated (*solid*), back-illuminated S1 (*dotted*), and S3 (*dashed*) chips (*Chandra* Proposers' Observatory Guide 2001[32]).

Unfortunately, the CCD chips onboard *Chandra* were damaged by charged particles on the orbit, degrading the energy resolution of ACIS. This is currently handled by lowering the temperature of the device to −110°C (as of our observation date) and by introducing a new set of detector response functions.

Effective Energy Range

Three factors are convolved to determine the energy range that an ACIS observation is sensitive to. The first factor is the quantum efficiency of the CCD chip itself. At higher energies, the efficiency is constrained by the thickness of the depletion layer, where X-ray photons with higher energy are more difficult to be photoelectrically absorbed. At lower energies, on the other hand, the efficiency is reduced by the absorption of X-ray photons by electrodes and insulators on the surface of the device. This makes back-illuminated chips to have higher quantum efficiency at lower energies than front-illuminated ones. The second factor is the transmission of the optical blocking filters (OBFs). ACIS has three OBFs with the thickness of 2000 Å (thick), 1200 Å (medium), and 400 Å (thin) to block photons in the optical wavelengths. Filters are composed of polyimide (a poly-carbonate plastic) sandwiched between two thin layers of aluminum. The third is the effective area of HRMA, which depends on the incident X-ray energy. Figure 4.9 shows the effective area of all these factors combined. In our observation, we used front-illuminated CCDs with the medium thickness filter, making our observation sensitive to X-ray photons in the 0.5–8.0 keV energy range.

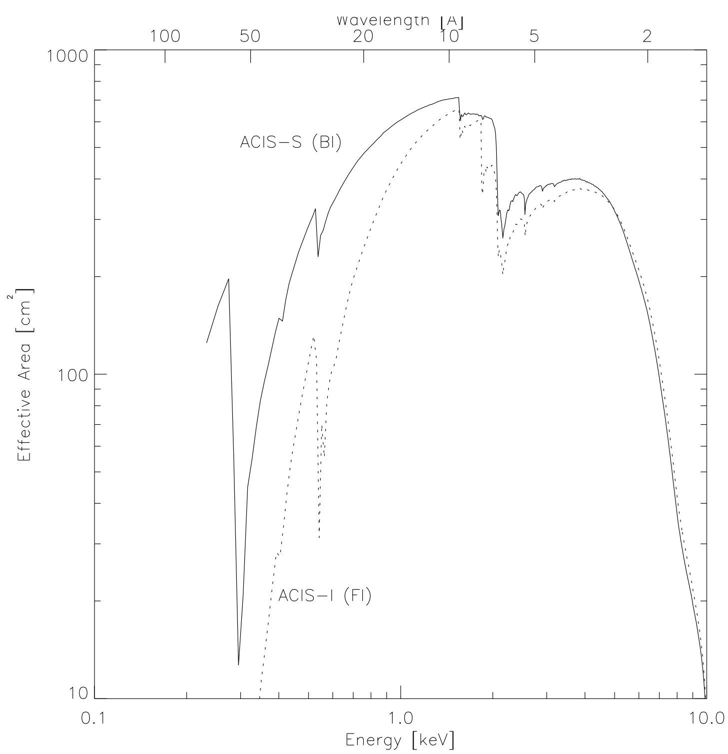


Figure 4.9: Expected effective area of ACIS with the dotted (front-illuminated chips) and the solid (back-illuminated chips) curves. The quantum efficiency, the filter transmission, and the effective area of the optical system are all combined (*Chandra* Proposers' Observatory Guide 2001[32]).

Background Events

The background events of ACIS are dominated by two components; cosmic ray (CR) events and the cosmic X-ray backgrounds (CXB). The CR events are reduced by various filtering procedures, including the grade filtering and the elimination of flaring events caused by CRs. A typical background spectrum after these procedures is given in Figure 4.10. The CR events are dominant in the lower energy band than 5 keV, while CXB in the higher energy band. The average count rate of this background spectrum is shown for some representative energy ranges (Table 4.1). The background count rate is roughly constant within a period of the same CCD temperature (Fig. 4.11). Unlike CXB, the rate of CR events also depend on the position on the chip (Fig. 4.12) by a factor of ∼20%.

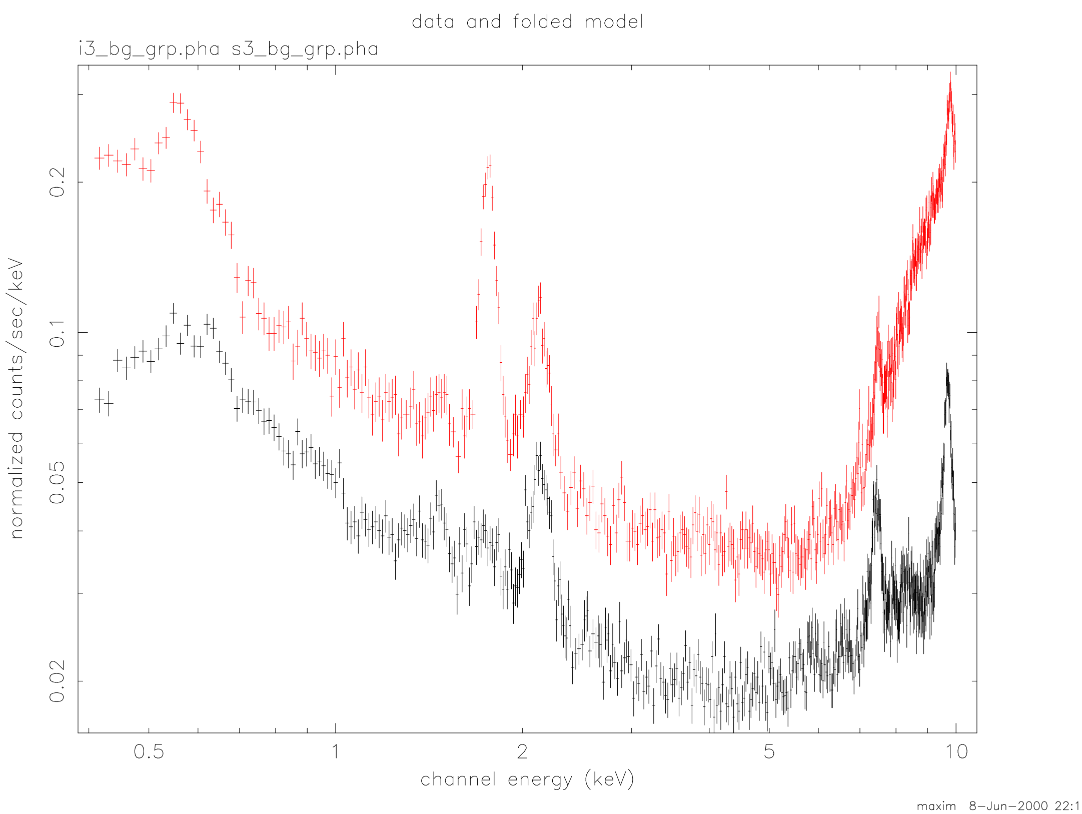


Figure 4.10: Background spectra of the back-illuminated ACIS-S3 chip (the upper spectrum) and the front-illuminated ACIS-I3 (the lower spectrum) chip. The spectra, containing both the CR and CXB components, are extracted from the whole relevant chips from the September 1999 – January 2000 data. The events with bad grades, bad pixels and bright celestial sources are eliminated (Markevitch 2001[121]).

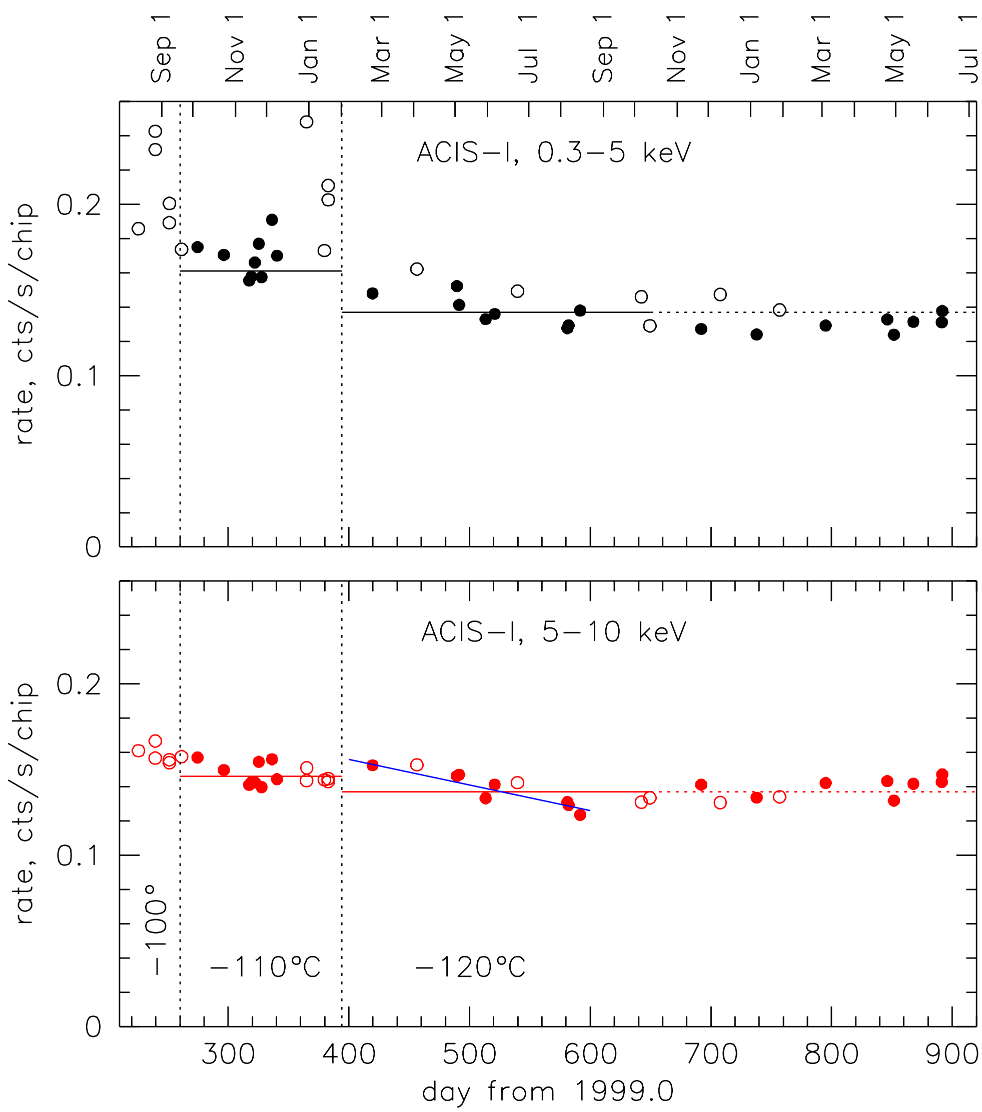


Figure 4.11: Long-term variation of background count rates of ACIS. The top panel shows the count rate of 0.3–5 keV that represents the CR component, while the bottom shows the rate of 5–10 keV that represents CXB. The CCD temperature is given at each period (Markevitch 2001[121]).

Table 4.1: Background count rate (s^{-1}) of *Chandra* ACIS chips for five representative energy ranges (Markevitch 2001[121]). The rates in 0.5–8 keV were estimated using these values and the standard background spectrum (Figure 4.10).

chip	I2[a]	I3[a]	S1[b]	S2[a]	S3[b]
0.3–10 keV	0.321	0.310	1.483	0.336	0.857
0.5–2 keV	0.083	0.076	0.168	0.086	0.160
0.5–7 keV	0.196	0.188	0.454	0.207	0.375
5–10 keV	0.153	0.154	1.038	0.162	0.506
10–12 keV	0.087	0.087	0.674	0.081	0.603
0.5–8 keV	0.226	0.220	.....	0.239	.....

[a] Front-illuminated chips.
[b] Back-illuminated chips.

Other Detectors

High-Resolution Imager (HRI) is another focal plane instrument. HRI is comprised of two imaging detectors using micro-channel plates; High Resolution Camera (HRC)-I designed for wide-field imaging and HRC-S designed to serve as a readout for the Low Energy Transmission Grating (LETG). HRI has no spectral resolution, but provides better temporal resolution than ACIS.

Each of the two instruments (ACIS and HRI) can be combined with one of the two gratings to conduct high-resolution ($E/\Delta E \sim 100$–1000) spectroscopy. High Energy Transmission Grating (HETG) and LETG are respectively optimized for grating spectroscopy of the high and low energy X-rays.

4.2 University of Hawaii 88 inch (2.2 m) Telescope

4.2.1 Telescope

The University of Hawaii 88 inch (2.2 m) telescope (UH88) is one of the telescopes at the summit of Mauna Kea, Hawaii, U.S.A. It was constructed in 1970 and has been under

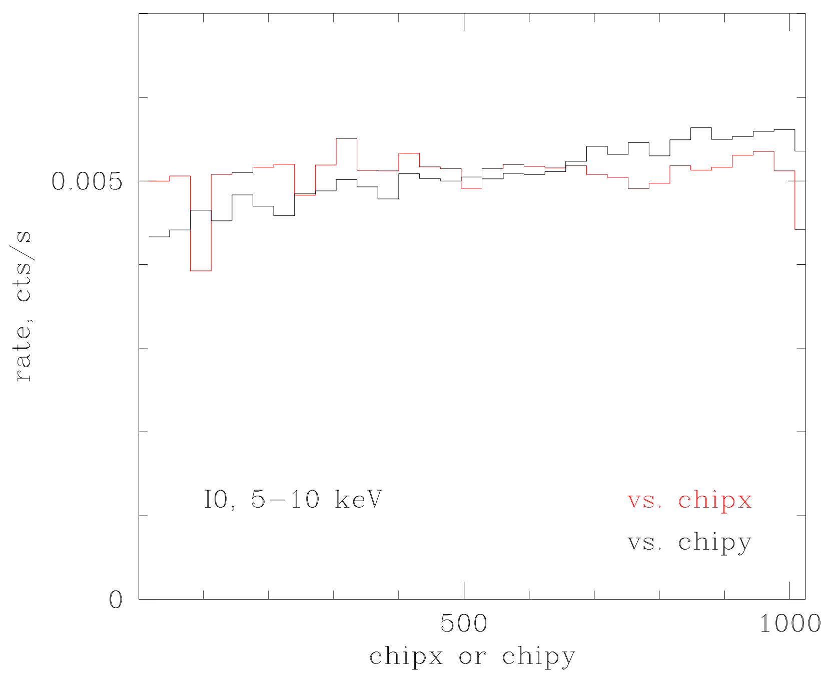


Figure 4.12: Spatial dependence of the background count rate of ACIS-I0. The profiles on the x-axis (*dotted*) and y-axis (*solid*) are shown. The count rate in the 5–10 keV range is used, where the CR component is dominant. The CXB component has no spatial dependence (Markevitch 2001[121]).

operation of the University of Hawaii. The telescope utilizes the Ritchey-Crétien optics system and has the Cassegrain focus. The primary mirror has a diameter of 88.13 inch (2.24 m) with the focal length of 22.50 m (Fig. 4.13).

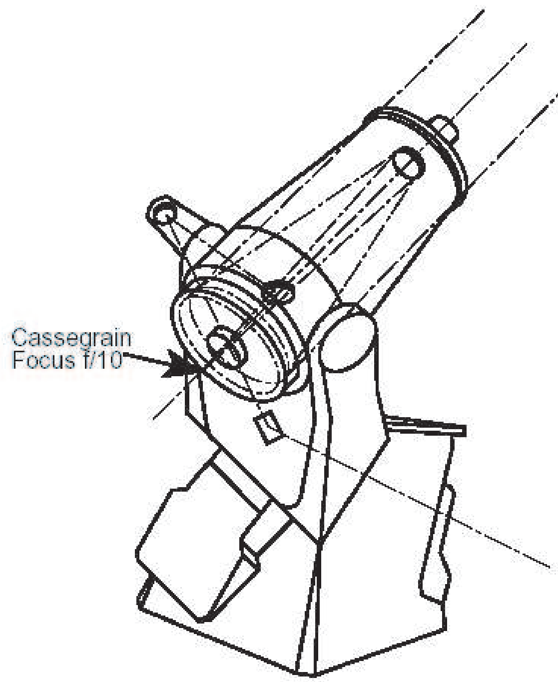


Figure 4.13: Schematic view of the University of Hawaii 88 inch (2.2 m) telescope. QUIRC is mounted on the Cassegrain focus of the telescope.

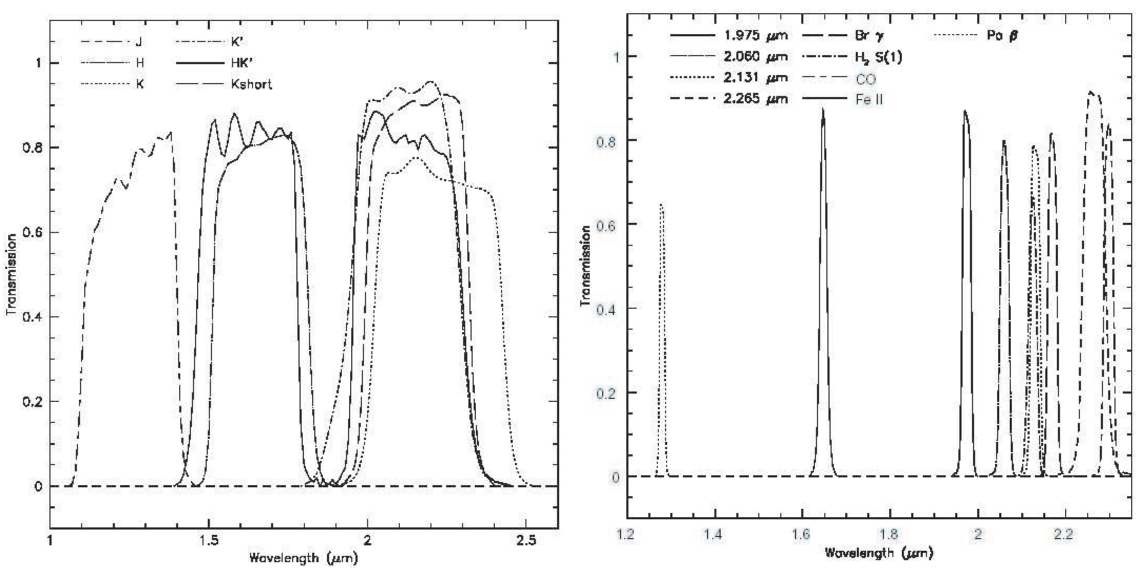
4.2.2 Instrument (QUIRC)

Quick Infrared Camera (QUIRC; Hodapp, Hora, & Metzger 1997[89]) is the only NIR camera mounted on the Cassegrain focus of the telescope. QUIRC has a HAWAII (HgCdTe Astronomical Wide Area Infrared Imaging) array produced by Rockwell Science Center. It consists of four quadrants, each of which has the format of 512×512 pixels with the pixel size of 18.5 μm×18.5 μm. It has two optics with different focal lengths ($f/10$ and $f/31$), which yields the pixel scale of 0.1886″ pixel^{-1} and 0.0608″ pixel^{-1}, and the FOV of 193″×193″ and 62″×62″, respectively. In our observation, we used the $f/10$ optics to maximize FOVs.

QUIRC is sensitive to the radiation from 1 μm to 2.5 μm. Two filter wheels with eight positions are combined to observe either in the broad-band (J, H, K, K_s, K', and

HK′) or in the narrow-band (H_2 $v = 1 - 0$ S(1) , Fe II, CO band head, Br γ, and their neighboring continuum). We used three broad-band (*J*, *H*, and *K*) and two narrow-band (H_2 $v = 1 - 0$ S(1) at 2.12 μm and *K-continuum* at 2.26 μm) filters. The transmissions of these filters are given in Figure 4.14.

The effective gain of the detector is 1.85 electrons ADU^{-1} and the linearity is kept better than 1% for values up to ∼44000 ADUs. The average detector dark current is $\leq$ 0.8 electrons s^{-1} and the readout noise is $\leq$ 15 electrons (r.m.s.). The camera sensitivity is ∼18.6 mag (*J* band), ∼17.8 mag (*H* band), and ∼16.2 mag (*K* band), assuming one minute on-source integration time, a PSF of 0.5″ FWHM, and 5 σ detection.



(a) broad-bands (b) narrow-bands

Figure 4.14: Transmissions of QUIRC filters; (*a*) broad-band and (*b*) narrow-band filters.

4.3 Subaru Telescope

4.3.1 Telescope

The Subaru telescope is also at the summit of Mauna Kea. It celebrated the first light in early 1999 and has been operated by National Astronomy Observatory of Japan (NAOJ).

Table 4.2: Comparison of NIR imagers used in this thesis

telescopes	UH88	Subaru	IRTF
instruments	QUIRC	IRCS	NSFCam
format (pixels)	1024×1024	1024×1024	256×256
pixel scale ($''$ pixel^{-1})	0.189	0.058	0.15
FOV ($''$)	193×193	59.4×59.4	76.8×76.8
filters used	J, H, K, H_2, K-cont.	J, H, K, H_2, K-cont.	L'
sensitivity (mag)[a]	∼16	∼18	∼17

[a] The 5 σ detection limit in the K band with one minute on-source integration time.

The telescope has the alto-azimuth mounting with four foci; the prime focus, the Cassegrain focus, and two Nasmyth foci for optical and infrared detectors (Fig. 4.15). It has the largest primary mirror among all the single mirrors facilitated in the telescopes around the world, with the diameter of 8.2 m and the focal length of 15 m.

4.3.2 Instrument (IRCS)

Infrared Camera and Spectrograph (IRCS) is a NIR imager and spectrometer for wide variety of usage (Tokunaga et al. 1998[184]; Kobayashi et al. 2000[105]). It has two ALADDIN II 1024×1024 pixel InSb arrays; one is for the imaging and the grism spectroscopy and the other for the echelle spectroscopy observations.

As an imager, it has two pixel scales of 0.058$''$ pixel^{-1} and 0.023$''$ pixel^{-1} with the corresponding FOV of 59.4$''$×59.4$''$ and 23.6$''$×23.6$''$. They are optimized to be used with the tip-tilt and adaptive optics, respectively. The detector is sensitive at 0.9–5.5 μm with the filters of J, H, K', K, L', and M' for the broad-band, and Fe II, Br α, Br γ, H_2 $v = 1-0$ S(1), 2–1 S(1), etc. and their neighboring continuum for the narrow-band. The transmissions of broad-band filters are shown in Figure 4.16.

The detector gain is 12.2 electrons ADU^{-1} and the linearity is kept up to ∼136000 ADUs. The dark current and the readout noise are 0.2–0.3 electrons s^{-1} and 67 electrons (r.m.s.), respectively. We used non-destructive readouts to decrease readout noise. The camera attains $J \sim 23.6$ mag, $H \sim 22.8$ mag, and $K \sim 22.3$ mag with one-hour integration time, a PSF of 0.5$''$ FWHM, and 5 σ detection. In this observation, we used the ALLADIN II

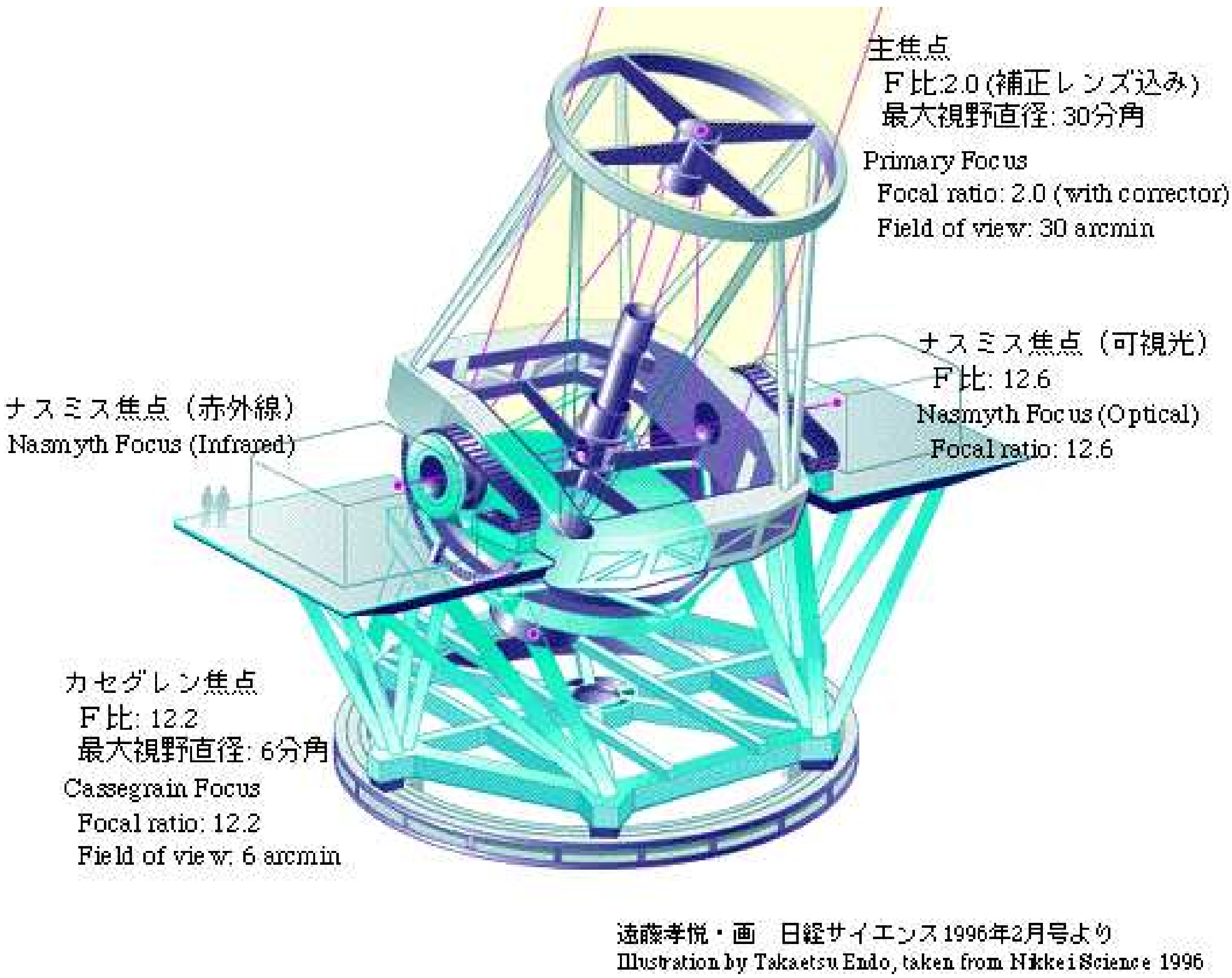


Figure 4.15: Schematic view of the Subaru telescope. IRCS is mounted on the Cassegrain focus of the telescope.

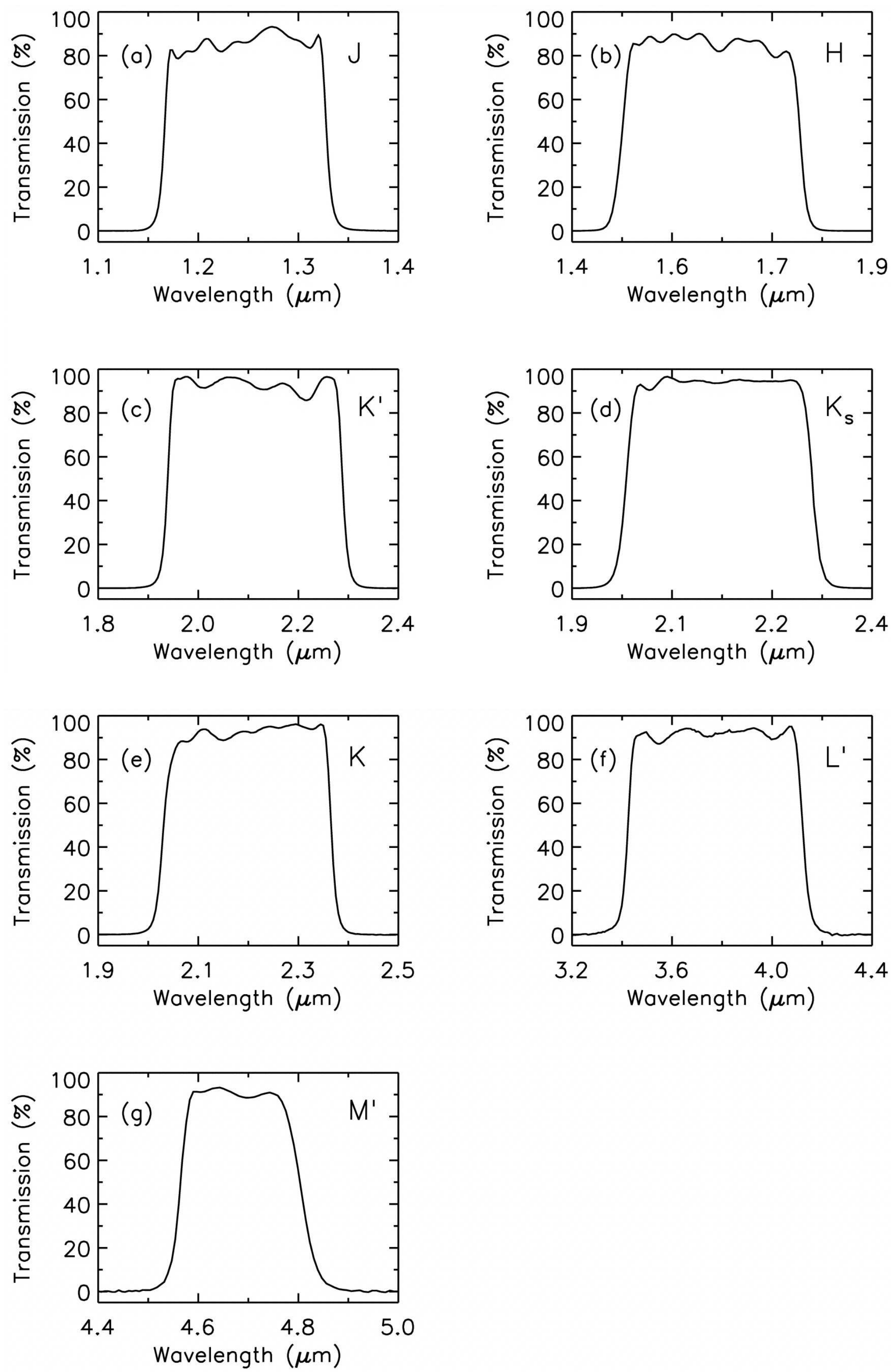


Figure 4.16: Broad-band filter transmissions of the Mauna Kea system advocated by Tokunaga, Simons, & Vacca (2002)[186]. Subaru IRCS and IRTF NSFCam comply with the system.

for the camera array, which was replaced with the ALLADIN III array with much superior sensitivity in September 2001.

IRCS camera also serves as a grism spectrograph providing a spectral resolving power of $R = 100$–2000. IRCS also has a separated echelle spectrograph providing a spectral resolving power of $R = 5000$–20000.

4.4 Infrared Telescope Facility

4.4.1 Telescope

The Infrared Telescope Facility (IRTF) is located at the summit of Mauna Kea. It is constructed in 1979 and is operated for NASA by the University of Hawaii. IRTF is optimized for NIR observations. It has the primary mirror of the 3.0 m in aperture and the Cassegrain focus.

4.4.2 Instrument (NSFCam)

NSFCam (Leggett & Denault 1996[115]) is a NIR camera mounted on the Cassegrain focus of the telescope. It employs an InSb array with the format of 256×256 pixels. Three different magnifications can be selected that respectively yield the pixel scale of $0.3''$ pixel^{-1}, $0.15''$ pixel^{-1} and $0.06''$ pixel^{-1} with the corresponding FOVs of $76.8''\times76.8''$, $37.9''\times37.9''$, and $14.1''\times14.1''$.

NSFCam has the sensitivity in 1–5 μm. By rotating its filter wheel, broad-band (J, H, K, K', L, L', and M) and narrow-band (He I, Fe II, H_2 $v = 1 - 0$ S(1), $v = 2 - 1$ S(1), CO band-head and their neighboring continuum) filters can be selected. We obtained the L'-band images with the pixel scale of $0.3''$ pixel^{-1} to complement the J-, H-, and K-band images taken with the Subaru telescope. The filters also comply with the Mauna Kea filter system (Fig. 4.16).

The camera can detect $J \sim 20.6$ mag, $H \sim 18.9$ mag, $K \sim 18.8$ mag, and $L' = 13.6$ mag with one-minute integration time, an aperture of $2''\times2''$box, and 3 σ detection.

4.5 Very Large Array

4.5.1 Telescope

Very Large Array (VLA) is the radio interferometer on Plains of San Agustin, New Mexico, U.S.A. It celebrated its first fringe in February 1976 and has been operated by National Radio Astronomy Observatory (NRAO). The interferometer consists of 27 antennae, which collaborate together for interferometry observations. The antennae have the size of 25 m in diameter and are positioned in the "Y" shape with nine antennae on each arm (Fig. 4.17).

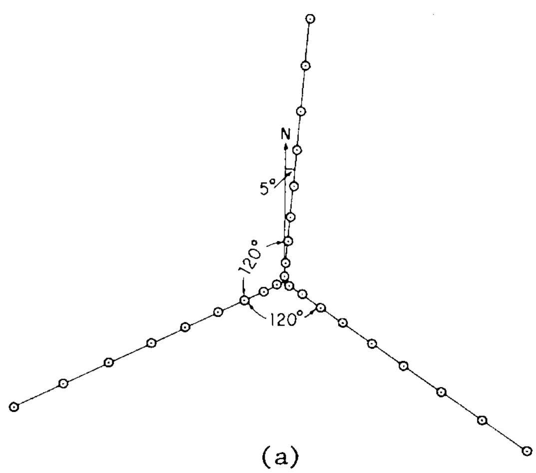


Figure 4.17: Configuration of VLA. The 27 antennae are laid out in the "Y" shape with the appropriate spacing to optimize the efficiency of $u - v$ coverage.

Four configurations (A, B, C, and D) with different arm lengths are available and the configuration changes once in four months. The longer arm length configurations have longer baselines (distances between two antennae), thus have higher spatial resolution but are less sensitive to extended structures (Table 4.3). At the Cassegrain focus of each antenna, receivers in eight frequency bands are equipped (Table 4.4). Our observation was performed in 3.6 cm with the A configuration. Details on the radio interferometry and VLA can be found in Clark (1999)[36], Thompson (1999)[180], Napier (1999)[140], Rommey (1999)[166], and Perley & Taylor (2002)[151].

Table 4.3: VLA configurations

configuration	A	B	C	D
min. baseline (km)	0.68	0.21	0.035	0.035
max. baseline (km)	36.4	11.4	3.4	1.03
synthesized beam width (″) [a]	0.24	0.7	2.3	8.4

[a] The half power beam width at 3.6 cm.

Table 4.4: VLA frequency bands

band	frequency (Hz)	wavelength (cm)	resolution[a] (″)
4	0.073–0.0745	400	24.0
P	0.30–0.34	90	6.0
L	1.34–1.73	20	1.4
C	4.5–5.0	6	0.4
X	8.0–8.8	3.6	0.24
U	14.4–15.4	2	0.14
K	22–24	1.3	0.08
Q	40–50	0.7	0.05

[a] The synthesized beam width of the A configuration.

4.5.2 Radio Interferometry

Let us suppose that the surface brightness of a celestial source is $\epsilon_\nu(\mathbf{R})$ at the position of $\mathbf{R}$ and at the frequency of ν. The quasi-monochromatic component of the electric field that we receive at the position of $\mathbf{r}$ can be expressed as

$$E_\nu(\mathbf{r}) = \int_S d\mathbf{R}\ \epsilon_\nu(\mathbf{R}) \frac{\exp(2\pi i\nu|\mathbf{R}-\mathbf{r}|/c)}{|\mathbf{R}-\mathbf{r}|}, \tag{4.1}$$

by integrating $\epsilon_\nu(\mathbf{R})$ over the source region (S). The spatial coherence function is defined as the product of $E_\nu(\mathbf{r})$ at two different positions ($\mathbf{r_1}$ and $\mathbf{r_2}$);

$$\begin{aligned} V_\nu(\mathbf{r_1},\mathbf{r_2}) &= < E_\nu(\mathbf{r_1})E_\nu(\mathbf{r_2}) > \\ &= \Big\langle \int_S d\mathbf{R_1} \int_S d\mathbf{R_2}\ \epsilon_\nu(\mathbf{R_1})\epsilon^*_\nu(\mathbf{R_2}) \\ &\quad \times \left(\frac{\exp\left(2\pi i\nu|\mathbf{R_1}-\mathbf{r_1}|/c\right)}{|\mathbf{R_1}-\mathbf{r_1}|}\right)\left(\frac{\exp\left(-2\pi i\nu|\mathbf{R_2}-\mathbf{r_2}|/c\right)}{|\mathbf{R_2}-\mathbf{r_2}|}\right)\Big\rangle. \end{aligned} \tag{4.2}$$

The surface brightness of the celestial source can be assumed to be coherent with each other at different positions with $< \epsilon(\mathbf{R_1})\epsilon(\mathbf{R_2}) >\propto \delta(\mathbf{R_1}-\mathbf{R_2})$. We can make a further approximation of $1/|\mathbf{R}-\mathbf{r}| \approx 1/|\mathbf{R}|$ to transform the spatial coherence function into

$$V_\nu(\mathbf{r_1},\mathbf{r_2}) = \int_S d\Omega\ I_\nu(\mathbf{s}) \exp\left(-2\pi i\nu \mathbf{s}\cdot(\mathbf{r_1}-\mathbf{r_2})/c\right), \tag{4.3}$$

where $\mathbf{s} = \mathbf{R}/|\mathbf{R}|$, $d\mathbf{R} = |\mathbf{R}|^2 d\Omega$, and $I_\nu(\mathbf{s}) = |\mathbf{R}|^2 < |\epsilon_\nu(\mathbf{s})|^2 >$. We introduce a new metric so that $\mathbf{r_1}-\mathbf{r_2} = \lambda(u,v,w)$ in order to measure lengths in the unit of the wavelength; $\lambda = c/\nu$. By setting the w-axis in the direction of $\mathbf{r_1}\times\mathbf{r_2}$, $w=0$. Then,

$$\begin{aligned} V_\nu(u,v) &= \int du' \int dv'\ I_\nu(u',v') \frac{\exp\left(-2\pi i(uu'+vv')\right)}{\sqrt{(1-u'^2-v'^2)}} \\ &\approx \int du' \int dv'\ I_\nu(u',v') \exp\left(-2\pi i(uu'+vv')\right). \end{aligned} \tag{4.4}$$

Here, $\sqrt{(1-u'^2-v'^2)} \approx 1$ is assumed for simplicity. By considering the vignetting effect $A_\nu(u',v')$ of the telescope,

$$V_\nu(u,v) = \int du' \int dv'\ A_\nu(u',v') I_\nu(u',v') \exp\left(-2\pi i(uu'+vv')\right). \tag{4.5}$$

The surface brightness distribution of the source $I_\nu(u,v)$ is thus derived by Fourier-converting $V_\nu(u,v)$.

When the spatial coherence function is applied to the simplest interferometer with two antennae at $\mathbf{r_1}$ and $\mathbf{r_2}$, the visibility function is obtained as

$$V = |V| \exp(i\phi) \tag{4.6}$$

$$= \int_S d\Omega \; A(\mathbf{s}')I(\mathbf{s}') \exp(-2\pi i\nu \mathbf{b} \cdot \mathbf{s}'/c), \tag{4.7}$$

where $\mathbf{b} = \mathbf{r_1} - \mathbf{r_2}$ is the baseline vector in the u-v plane and $\mathbf{s}' = \mathbf{s} - \mathbf{s_0}$ with $\mathbf{s_0}$ being the unit vector pointing at the phase tracking center.

The value of the visibility function over the u-v plane can be obtained through the correlater outputs. The correlater multiplies the voltages from two antennae ($V_1(t)$ and $V_2(t)$) and integrates it for a given period of time (T). As $V_1(t)$ and $V_2(t)$ have different phases stemming from their geometrical distance to the source;

$$V_1(t) = V_1 \cos(2\pi\nu(t - \mathbf{b} \cdot \mathbf{s}/c) \tag{4.8}$$

$$V_2(t) = V_2 \cos(2\pi\nu t). \tag{4.9}$$

The correlater output is described as

$$r \propto \frac{1}{T}\int_0^T dt \; V_1(t)V_2(t)$$

$$\longrightarrow V_1 V_2 \cos(2\pi\nu \mathbf{b} \cdot \mathbf{s}/c) \quad (T \longrightarrow \infty) \tag{4.10}$$

The output is also proportional to the power that an antenna receives; $A(\mathbf{s})I(\mathbf{s})\Delta\nu\Delta\Omega$. Therefore, by integrating $\mathbf{s}$ over the antenna beam width, the correlater output is

$$r = \Delta\nu \int_S A(\mathbf{s})I(\mathbf{s}) \cos(2\pi\nu \mathbf{b} \cdot \mathbf{s}/c) \tag{4.11}$$

When we substitute $\mathbf{s}$ with $\mathbf{s}' + \mathbf{s_0}$, we obtain

$$r = \Delta\nu \cos(2\pi\nu \mathbf{b} \cdot \mathbf{s_0}/c) \int_S d\Omega \; A(\mathbf{s}')I(\mathbf{s}') \cos(2\pi\nu \mathbf{b} \cdot \mathbf{s}'/c)$$

$$-\Delta\nu \sin(2\pi\nu \mathbf{b} \cdot \mathbf{s_0}/c) \int_S d\Omega \; A(\mathbf{s}')I(\mathbf{s}') \sin(2\pi\nu \mathbf{b} \cdot \mathbf{s}'/c) \tag{4.12}$$

The real and imaginary part of the visibility function;

$$\Re[V] = |V| \cos(\phi) = \int_S d\Omega \; A(\mathbf{s}')I(\mathbf{s}') \cos(2\pi\nu \mathbf{b} \cdot \mathbf{s}'/c) \tag{4.13}$$

$$\Im[V] = |V| \sin(\phi) = -\int_S d\Omega \; A(\mathbf{s}')I(\mathbf{s}') \sin(2\pi\nu \mathbf{b} \cdot \mathbf{s}'/c) \tag{4.14}$$

$$\tag{4.15}$$

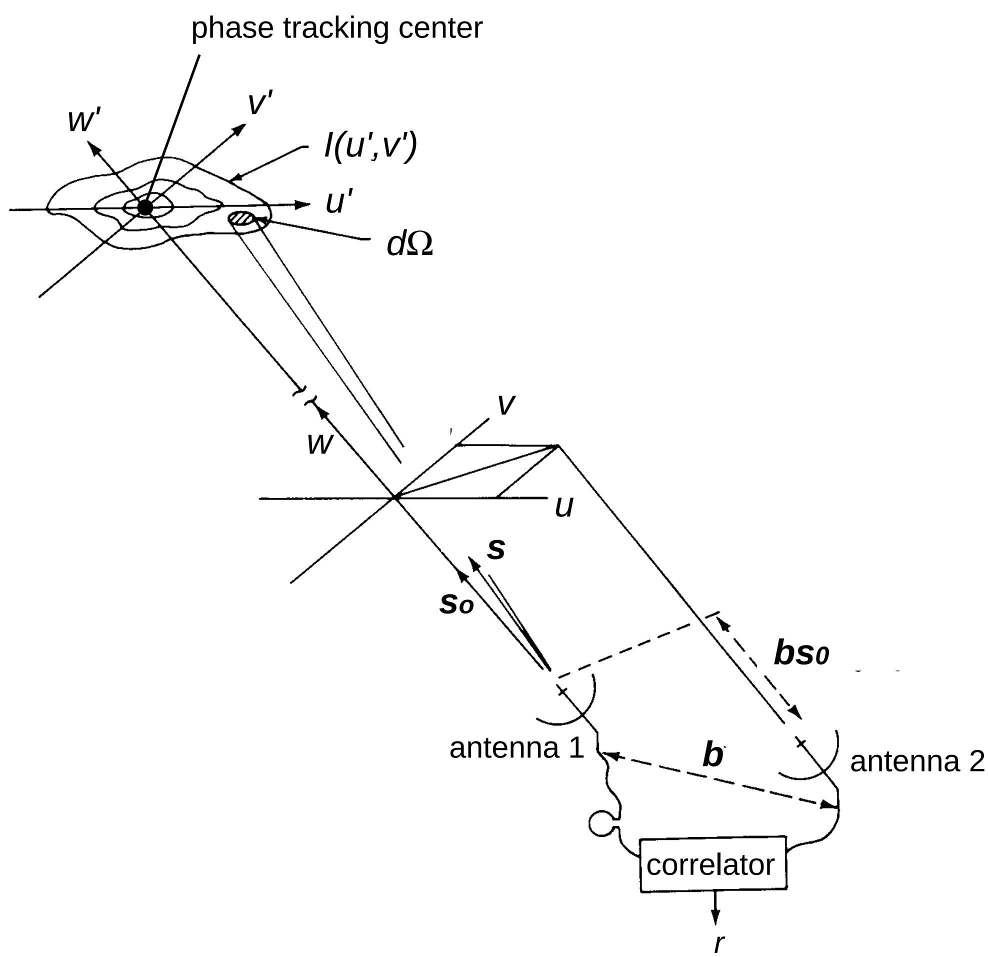


Figure 4.18: Schematic view of the radio interferometry (Rohlfs & Wilson 1999[165]).

appears in r, which connects r and V as

$$r = \Delta\nu|V|\cos(2\pi\nu\mathbf{b}\cdot\mathbf{s_0}/c-\phi). \quad (4.16)$$

By measuring the amplitude and the phase of r, the value of the visibility function is derived at a given point of the u-v plane. With N antennae, ${}_NC_2$ correlations can be obtained, which are all used to derive the value of the visibility function at ${}_NC_2$ points. These points move around the u-v plane as the earth rotates and $\mathbf{s_0}$ changes from time to time. The visibility function $V(u,v)$ is obtained after gridding the data points over the u-v plane, then is Fourier-converted to derive the surface brightness distribution.

Chapter 5

Chandra and QUIRC Observations

Contents

In this chapter, we describe our X-ray and NIR observations on OMC-2 and OMC-3, and their data reduction and results. We conducted a deep-exposure *Chandra*/ACIS observation and detected ∼400 X-ray sources in this field. These X-ray sources were first correlated with the 2MASS catalog to identify the NIR counterpart (Sect. 5.1). The 2MASS catalog, which covers brighter NIR sources than $K = 14$ mag, was found not to be deep enough to match the *Chandra* depth. In order to complement 2MASS with fainter sources, we made a follow-up NIR observation deeper on the *Chandra* field (Sect. 5.2) using QUIRC mounted on the University of Hawaii 88 inch (2.2 m) telescope. The *Chandra* sources that have no 2MASS counterpart were further correlated with the QUIRC sources to find their NIR counterpart. The X-ray and NIR source lists are separately given in Tables A.1 and B.1.

5.1 X-ray Observation with *Chandra*/ACIS

5.1.1 Observation

The *Chandra* observation on OMC-2 and OMC-3 was carried out on January 1–2, 2000 with a nominal exposure time of 88.4 ks. We used the four ACIS-I chips (I0, I1, I2, and I3) and one ACIS-S chip (S2) on the focal plane of the mirror system. The four ACIS-I chips cover a field of $\sim 17' \times 17'$ and the ACIS-S chip of $\sim 8.5' \times 8.5'$. The nominal center was set at R.A. (right ascension) = $05^{\rm h}35^{\rm m}20.792^{\rm s}$ and decl. (declination) = $-05°05'46.95''$ (the equinox J2000.0) so that the ACIS-I FOV covers the whole OMC-2 and OMC-3 (Fig. 5.1). The timed exposure mode was used as the operating mode with the frame time of 3.2 s.

5.1.2 Data Reduction

For data reduction, we used the level 2 data "reprocessed" at the *Chandra* X-ray Center (CXC). This version improves the aspect solution and restores the degradation of the energy gain and resolution due to the increase of the charge transfer inefficiency (CTI) of ACIS[1]. We also applied the acis_process_events program to the data in order to eliminate "spikings" in X-ray spectra reported by CXC[2]. For data manipulation, we used CIAO version 2.2[3] and

[1]See http://asc.harvard.edu/udocs/reprocessing.html.
[2]See http://asc.harvard.edu/ciao/caveats/acis_pi.html.
[3]See http://asc.harvard.edu/ciao/.

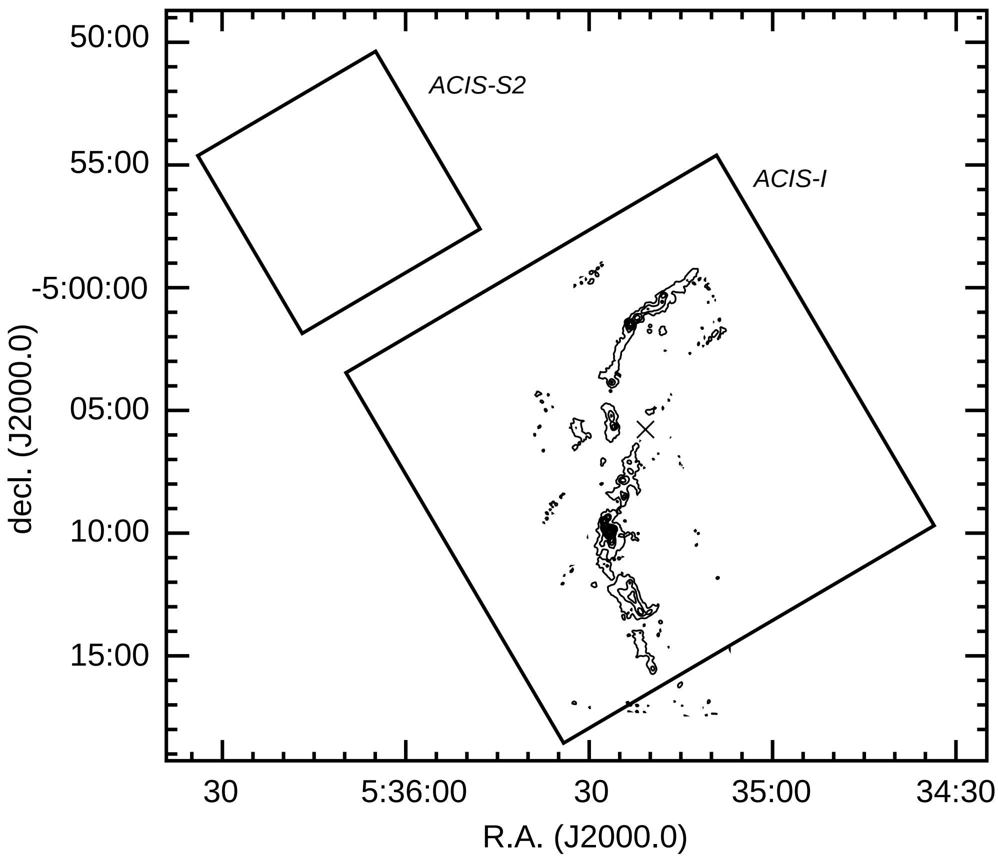


Figure 5.1: FOV of *Chandra*. The larger and smaller squares indicate the FOVs of ACIS-I and ACIS-S2, respectively. The aim point is shown with the cross. The contours are the 1.3 mm intensity (Chini et al. 1997[35]).

HEAsoft version 5.2[4]. Throughout this thesis, we used X-ray photons in the 0.5–8.0 keV energy band unless otherwise noted. Photons of each source were accumulated from an elliptical region. The major and minor axis lengths and the position angles were derived from the wavdetect program. For some sources with bright neighbors, we manually shifted their accumulation region to avoid contamination. Background photons were accumulated from the standard background data provided by the ACIS team[5], which combine some observations of relatively empty fields at high galactic latitudes. The pseudo-color image is shown in Figure 5.2. The *Chandra* PSF radius increases as the off-axis angle becomes larger (Fig. 4.6), which causes off-axis sources to appear extended.

5.1.3 Source Extraction

For source detection, we used the wavdetect program with the significance threshold of $1{\times}10^{-5}$ (one false recognition of event pixels is expected in a 10^5 pixel image) and the wavelet scales ranging from 1 to 16 pixels in multiples of $\sqrt{2}$. We removed spurious sources through careful inspection by eye. We then detected 365 sources in the 0.5–8.0 keV band image (I1–I354 from ACIS-I and S1–S11 from ACIS-S). In order to pick up either highly absorbed (hard) or less absorbed (soft) sources more effectively, we also applied the same detection algorithm to the 2.0–8.0 keV (hard) and 0.5–2.0 keV (soft) band images. Then, 17 (I355–I369 and S11-S12) and 16 (I370–I385) sources were additionally found, respectively. In total, we detected 398 *Chandra* sources in the ACIS-I and ACIS-S FOVs. For each detected source, we calculated the X-ray photon counts (0.5–8.0 keV) and the hardness ratio (HR), which is defined as $(H-S)/(H+S)$, where H and S are the photon counts in the hard and soft band, respectively (Table A.1).

5.1.4 2MASS Counterpart of *Chandra* Sources

2MASS Database

The Two Micron All Sky Survey (2MASS)[6] is a joint project among NASA Infrared Processing and Analysis Center (IPAC), Jet Propulsion Laboratory (JPL), the University of

[4]See http://heasarc.gsfc.nasa.gov/docs/software/lheasoft/.
[5]See http://cxc.harvard.edu/contrib/maxim/bg/index.html.
[6]See http://www.ipac.caltech.edu/2mass/ for more details.

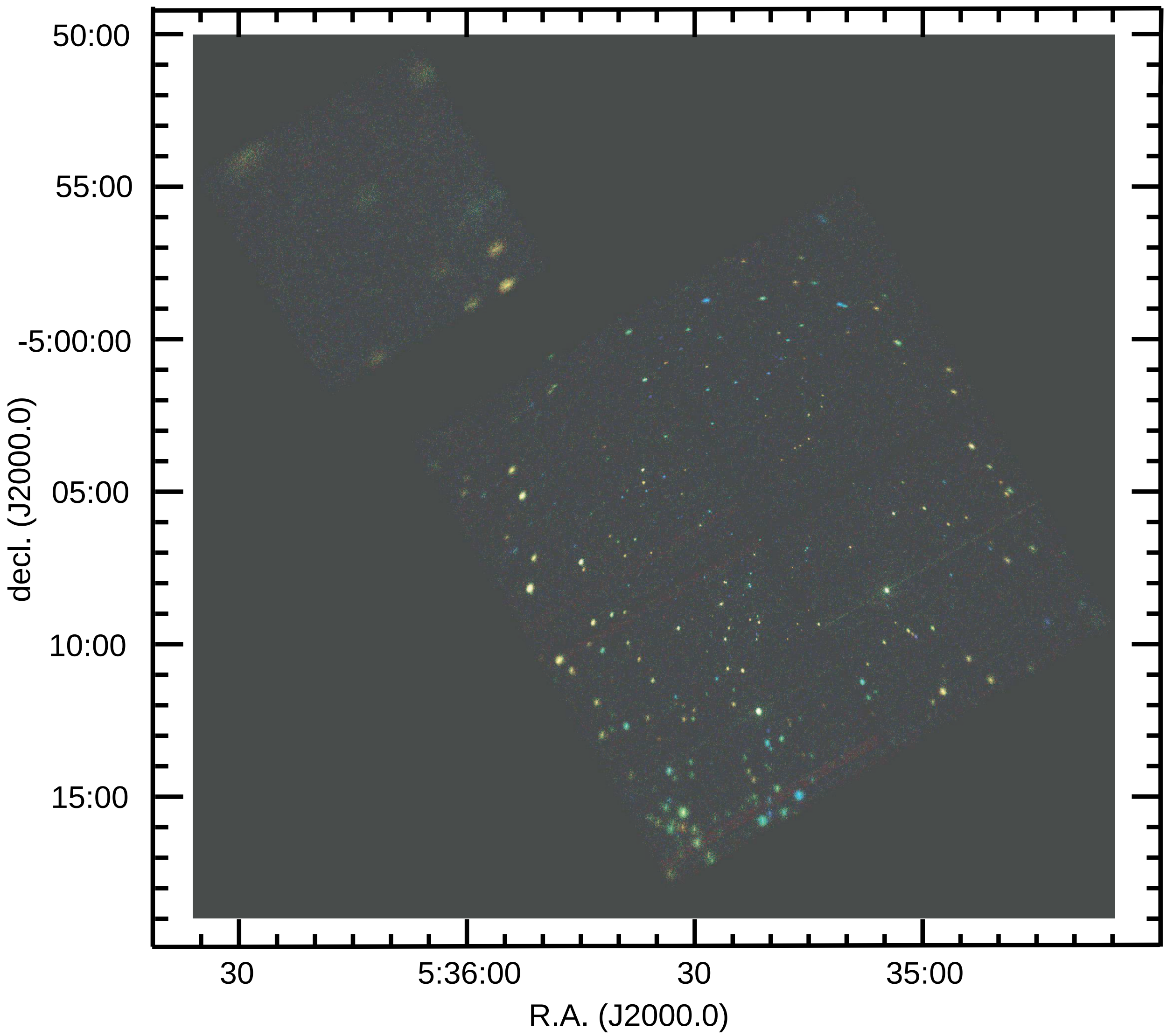


Figure 5.2: Three-color image of *Chandra* observations. Red, green, and blue are for photons in the 0.2–1.0 keV, 1.0–2.5 keV, and 2.5–8.0 keV, respectively.

Massachusetts, and California Institute of Technology (CIT), which aims to provide imaging data of the all sky in the J (1.25 μm), H (1.65 μm), and K_s (2.17 μm) bands. Two telescopes; one at Mt. Hopkins, Arizona, U.S.A. and another at Cerro Tololo, Chile, are exclusively devoted for this project to cover the northern and southern sky, respectively. The identical NICMOS3 arrays are used for three bands at two telescopes. The J-, H-, and K_s-band observations are conducted simultaneously with a 7.8 s exposure per frame. The FOV is $8.5'\times8.5'$ with the pixel scale of $1''$ pixel^{-1}.

The Second Incremental Data Release of 2MASS, which was released in March 2000, consists of two data sets; the Point Source Catalog (PSC) and the Extended Source Catalog (XSC). The release covers 19680.8 square degrees (~47% of the all sky; Fig. 5.3). OMC-2 and OMC-3 are fully covered.

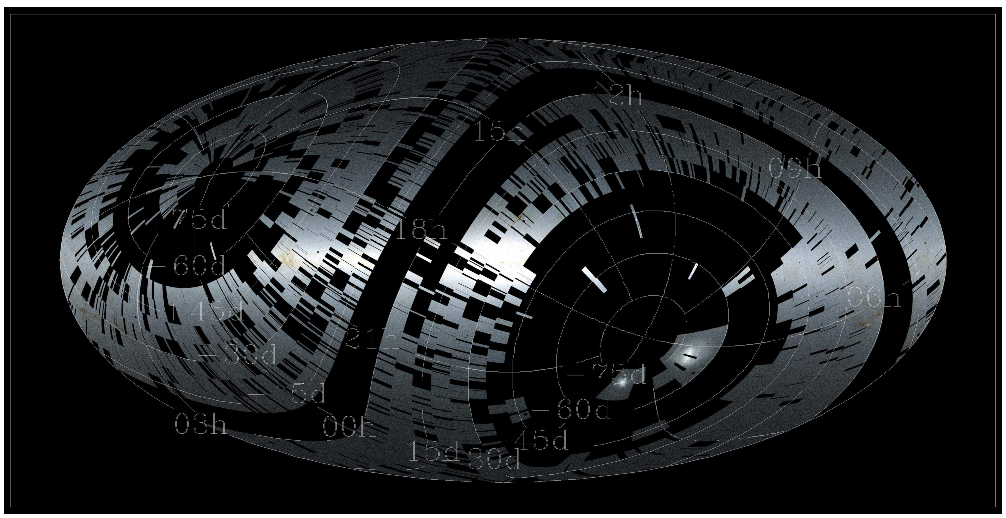


Figure 5.3: Sky coverage by the 2MASS Second Incremental Data Release. It covers ~47% of the all sky (the shining parts of the map; http://www.ipac.caltech.edu/2mass). OMC-2 and OMC-3 are fully covered at R.A. $\sim 5^{\rm h}30^{\rm m}$ and decl. $\sim -5^\circ$.

The PSC contains the position and the magnitude of 157,820,597, 149,650,034, and 130,337,158 sources in the J, H, and K_s band, respectively. The astrometric accuracy is $\sim0.1''$ both in the R.A. and decl. directions with a negligible systematic offset from the Astrographic Catalog/Tyco (ACT) catalog (Fig. 5.4). The photometric accuracy is given in Figure 5.5, where the uncertainties in magnitudes are less than 0.1 mag for the J-, H-, and K_s-band sources brighter than 15.8 mag, 15.1 mag, and 14.3 mag, respectively.

We consulted the PSC and found that ~600 sources are in the FOV of *Chandra*. These sources are correlated with QUIRC and *Chandra* sources, and were used to evaluate their astrometric and photometric accuracy.

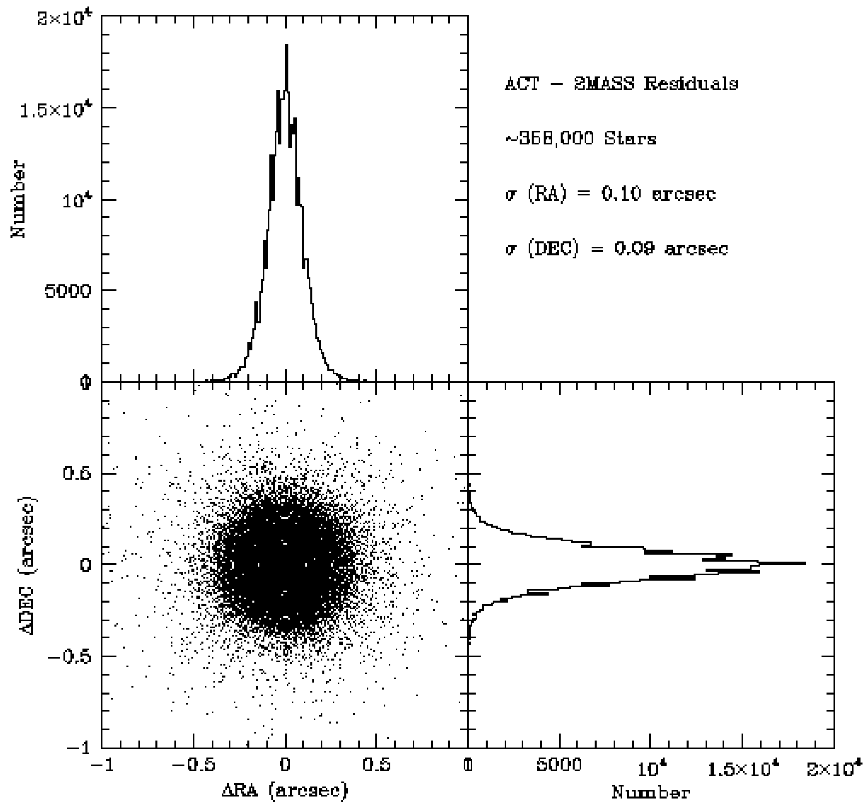


Figure 5.4: Distribution of the R.A. and decl. differences of ∼358,000 sources between 2MASS PSC and ACT (*bottom left*) and its profiles on R.A. (*top left*) and decl. (*bottom right*) axes. The 1 σ values for both profiles (=the typical astrometric accuracy of 2MASS PSC) are ∼0.1″ (http://www.ipac.caltech.edu/2mass).

2MASS Counterpart of *Chandra* Sources

We searched for the 2MASS counterpart of the *Chandra* sources in the following manner. In the ACIS-I and ACIS-S FOVs, we found 638 2MASS sources. First, we searched for the 2MASS source closest to each *Chandra* source within a 3″ radius. Second, we conversely searched for the *Chandra* source closest to each 2MASS source within a 3″ radius. Thus we picked up the closest *Chandra*-2MASS pairs. The systematic position offset of the *Chandra* sources from their 2MASS counterpart was found to be −0.18″ and 0.21″ in the direction

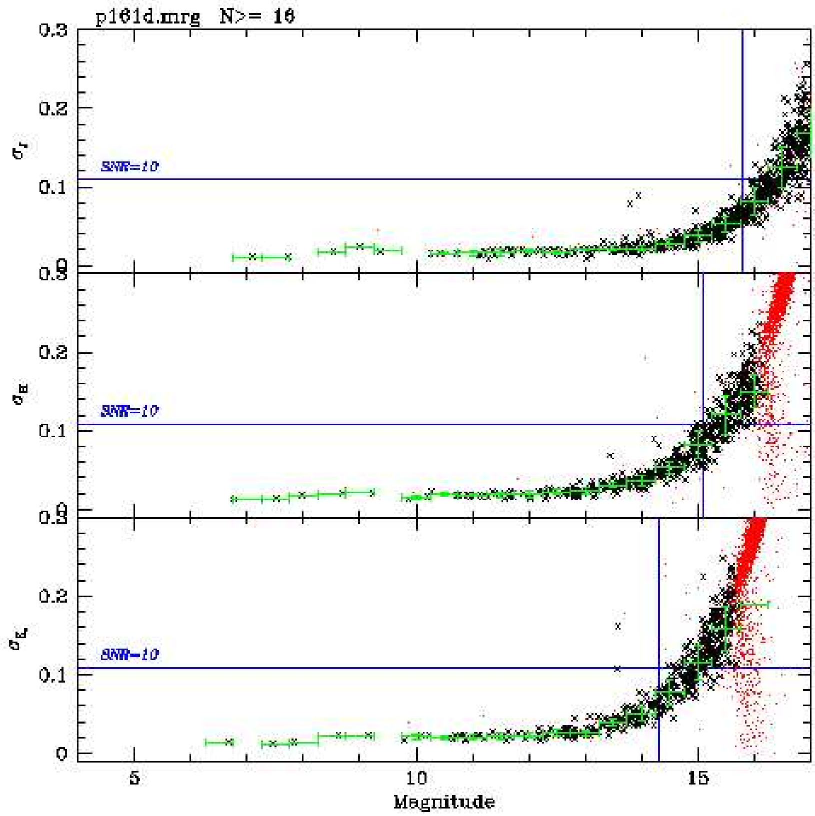


Figure 5.5: Photometric accuracy of sources detected in a calibration field. The uncertainty in magnitudes (*vertical axis*) versus the mean magnitude (*horizontal axis*) is plotted in black (sources detected at least 16 out of the 18 trials) and in red (sources detected fewer than 16 times) separately for the J (*top*), H (*middle*), and K_s (*bottom*) band. The dotted bars indicate the r.m.s. averaged in each 0.5 mag bin (http://www.ipac.caltech.edu/2mass).

of R.A. and decl., respectively. After correcting the *Chandra* positions for the systematic offsets, we repeated the same procedure for the 2MASS counterpart search. Finally, we found that 237 out of 398 (~60%) *Chandra* sources have the 2MASS counterpart.

5.1.5 Results

Source List

The results of the X-ray imaging analysis are compiled in Table A.1 with the X-ray source numbers, positions corrected for the systematic offset from the 2MASS frame, detector raw counts, and HRs. The X-ray sources can be referred following the International Astronomical Union (IAU) convention; e.g., CXOU J05343860−0508428 for the source No. 1.

Astrometric Accuracy

Figure 5.6 shows the differences between the *Chandra* and 2MASS positions both in the R.A. and decl. directions. We found that ΔR.A. = $0.00''\pm0.40''$ (1 σ) and Δdecl. = $0.01''\pm0.36''$ (1 σ), which indicates that the *Chandra* positions are determined with the accuracy of the size of a *Chandra* pixel ($0.492''$) and the systematic offset between *Chandra* and 2MASS positions is negligible.

Survey Depth

Figure 5.7 shows the histogram of X-ray counts. The peak of the histogram ($10^{1.5} = 32$ counts) roughly corresponds to the completeness limit of this observation, while the minimum number of counts (3 counts) gives the flux of the faintest detected sources. To convert the X-ray counts into the X-ray flux or luminosity, we made an empirical relation between these two parameters using the result of the spectral fitting of bright *Chandra* sources (Figs. 7.22 and 7.23). The completeness limit was estimated to be $\sim10^{-14.5}$ ergs s^{-1} cm^{-2} in flux and $\sim10^{29.1}$ ergs s^{-1} in luminosity in the 0.5–8.0 keV range. The X-ray flux and the luminosity of the faintest detected source were $\sim10^{-15.5}$ ergs s^{-1} cm^{-2} and $\sim10^{28.1}$ ergs s^{-1}, respectively.

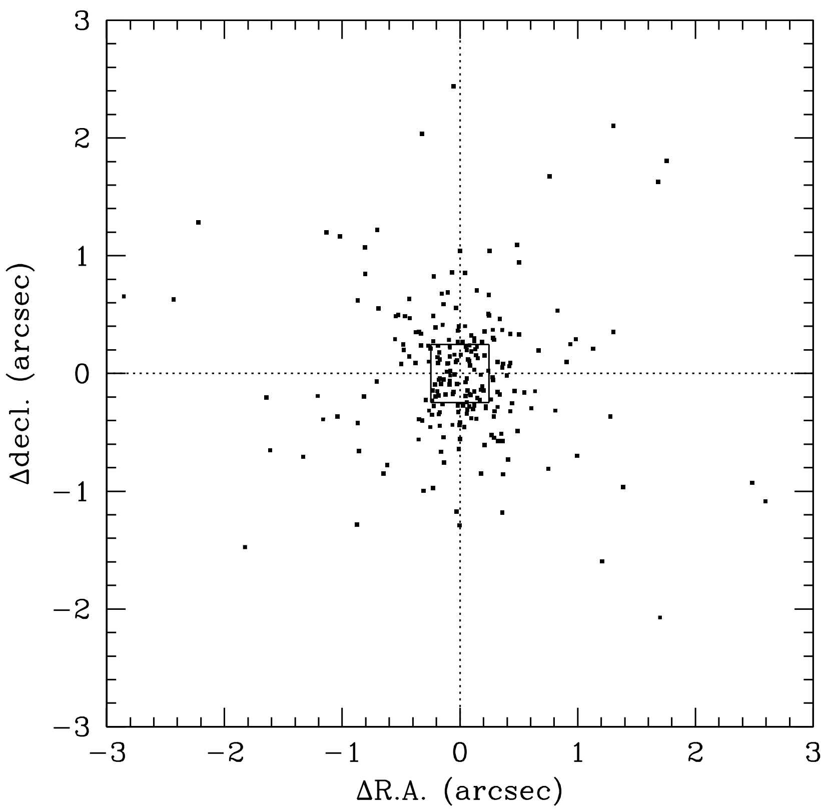


Figure 5.6: Astrometric accuracy of *Chandra* sources. The difference of R.A. and decl. in the *Chandra* and 2MASS positions (ΔR.A. and Δdecl., respectively) is plotted for *Chandra*–2MASS counterpart pairs. The solid square at the center represents the size of a *Chandra* pixel of $0.492'' \times 0.492''$.

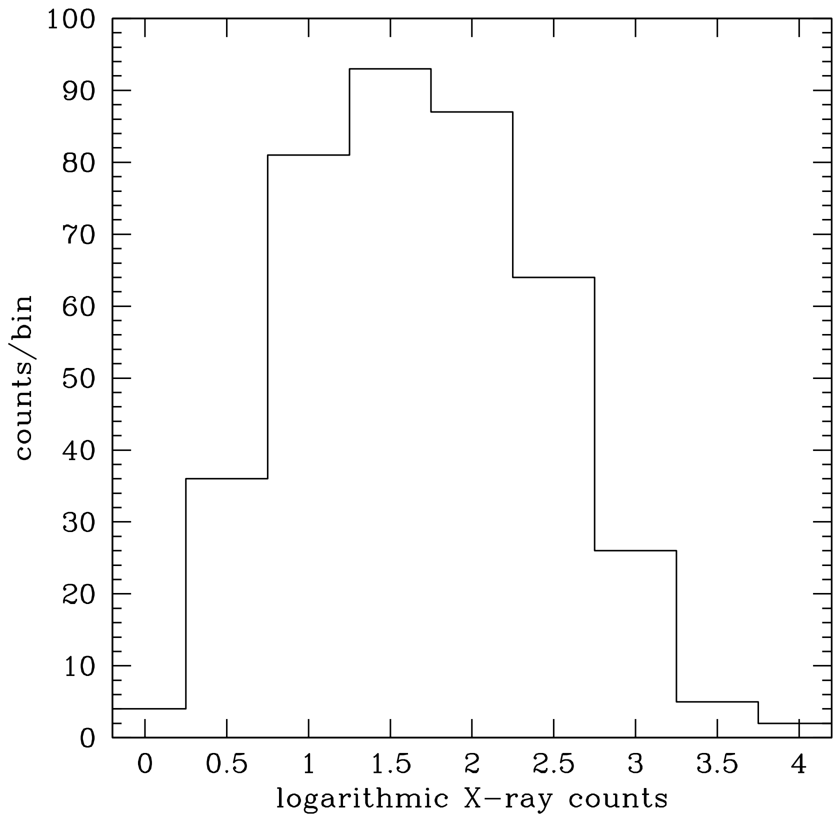


Figure 5.7: Histogram of *Chandra* counts of all detected sources.

5.2 NIR Observation with UH88/QUIRC

5.2.1 Observation

In two points, the 2MASS data are not sufficient to find the NIR counterpart of our X-ray sources and to identify their nature. One point is that the 2MASS data sometimes lack J-, H-, or both band detections in star-forming regions, where the extinction by dense ISM makes sources heavily reddened. Another point is that the 2MASS depth is not deep enough to match the depth of our *Chandra* observation. Casanova et al. (1995)[29] derived an empirical relation between the un-dereddened J-band magnitude (J) and the X-ray flux (F_{X}) in the 1.0–2.4 keV band using T Tauri star samples in the ρ Ophiuchi cloud ($D =$ 160 pc) detected by *ROSAT*, which is expressed as

$$\log F_X (4\pi D)^2 \sim -0.30J + 32. \tag{5.1}$$

The faintest detected sources of our *Chandra* observation have $F_{\mathrm{X}} \sim 10^{-15.5}$ ergs s^{-1} cm^{-2} in 0.5–8.0 keV at the distance of $D = 450$ pc, which can be converted to $F_{\mathrm{X}} \sim 2{\times}10^{-16}$ ergs s^{-1} cm^{-2} in the 1.0–2.4 keV band assuming a thin-thermal plasma spectrum of 1 keV temperature. If the equation (5.1) can be applied to other YSOs, the required NIR detection limit would be $J \sim 17$ mag. In order to complement the 2MASS data with fainter sources than the 2MASS limit, we conducted deeper NIR observations in OMC-2 and OMC-3 with the J-band integration time twice longer than that of the H and K bands.

We used QUIRC mounted on the Cassegrain focus of the University of Hawaii 88 inch (2.2 m) telescope (UH88; Hodapp et al. 1997[89]). QUIRC provides a $3.2'{\times}3.2'$ FOV with the pixel scale of $0.189''$ pixel^{-1}. The smaller pixel scale than that of the 2MASS detectors is more appropriate to pick up sources contaminated by diffuse emissions, particularly in the southern half of our study field (Fig. 5.9).

We conducted a mosaic mapping observation to sweep the *Chandra* ACIS-I FOV (Fig. 5.8). We first covered the uppermost row (frame Nos .1–5) of the mosaic by shifting the field center by $1.6'$ (half of the side of the QUIRC FOV) from east to west. Then, we moved $1.6'$ southward and covered the second uppermost row (frame Nos. 6–13) by inversely shifting from west to east. We continued these raster scans to reach the lowermost row of the mosaic. In this way, we can minimize dead times due to the shifts of field center and reduce the effect of ghost signals caused by residual electrons in detector pixels. In total, we swept the *Chandra* FOV with 169 QUIRC FOVs that amount to $\sim$512 arcmin^2.

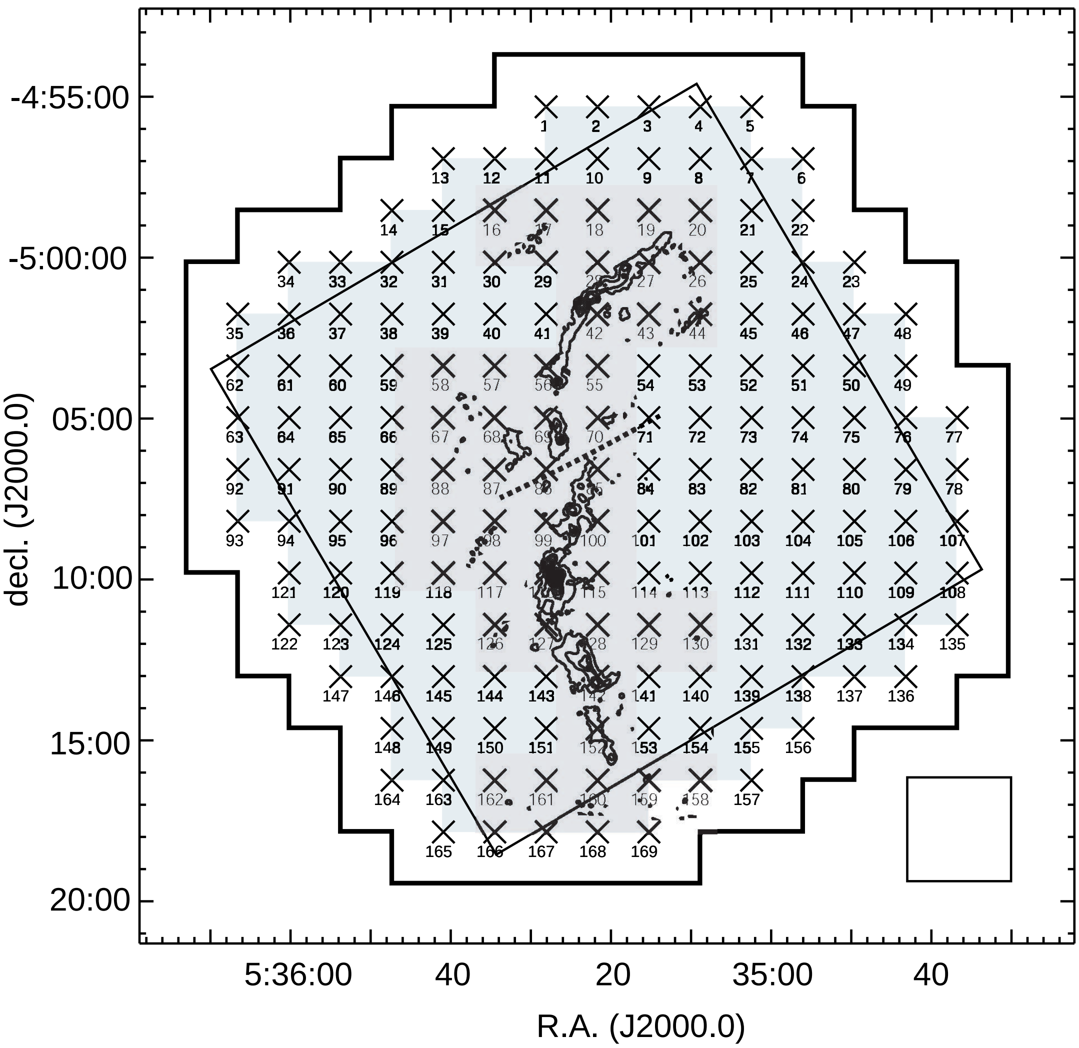


Figure 5.8: Configuration of our QUIRC mosaic mapping observations. The positions of 169 frame centers are marked with crosses with the relevant frame number. The frames are tiled at the intervals of 1.6′ to sweep the *Chandra* ACIS-I FOV (*solid oblique square*). Each frame covers a 3.2′×3.2′ square region (the size is shown at the right bottom) centered on each cross. When the all frames are combined, any point in the thick lines (∼512 arcmin^2) is covered at least by one frame, while that in the gray region (∼360 arcmin^2) is covered by four frames. The contours show the 1.3 mm intensity (Chini et al. 1997[35]). OMC-2 and OMC-3 are separated by the dashed line into the southern and northern part of the integral-shaped ridge, respectively.

Each 169 frame was exposed for 60 s. We conducted the mosaic mapping once in the H and K band and twice in the J band, spending six half-nights of February 4–6 and March 11–13, 2001 (Table 5.1). All nights were photometric with the seeing of $0.7''$–$1.1''$. In one sweep, any region inside the *Chandra* field except for edges (the gray region in Fig. 5.8) was covered four times, which makes the nominal exposure time to be 240 s in the H and K band and 480 s in the J band.

Table 5.1: QUIRC observation log

date	band	frame numbers[a]
2001 Feb. 04	K	Nos. 1–169
2001 Feb. 05	J	Nos. 1–52
2001 Feb. 06	H	Nos. 1–169
2001 Mar. 11	J	Nos. 1, 51–169
2001 Mar. 12	J	Nos. 1–129
2001 Mar. 13	J	Nos. 1, 107–169

[a] The frame numbers (see Fig. 5.8) covered in each night. In the J-band sweep, for which we spent four nights, some frames are duplicated; i.e., the frame No. 1 at the beginning of every run to check the telescope positioning accuracy, and frame Nos. 51–52 and 107–129 to check any difference of image quality in different nights. All these duplicated frames were confirmed to have similar quality and were combined equally into the final J-band image with the correction for the exposure time.

5.2.2 Data Reduction

All QUIRC frames were reduced following the standard procedures using IRAF[7]; i.e., dark subtraction, flat fielding, sky subtraction, and bad pixel removal. Dark frames with a 60 s integration time were taken at the end of each night for each filter, and were used to subtract dark current signals of source images. Flat frames were also taken at each night for each filter by observing the telescope dome, and were used to correct for the pixel-to-pixel variation in quantum efficiency. Sky frames were constructed for each frame set of a mosaic image by adopting the median of 169 ADU values at a given pixel. SExtractor (Bertin & Arnouts 1996[21]) was used for source extraction and photometry.

For the purpose of the astrometric and photometric calibration of all frames, we consulted the 2MASS catalog. We tentatively extracted sources from each frame and correlated

[7] See http://iraf.noao.edu/.

them with the 2MASS sources using WCSTools[8]. We found that all QUIRC frames have 18–145 sources including 4–51 sources with the 2MASS counterpart. Using the QUIRC–2MASS counterpart pairs, we first shifted each QUIRC frame so that the mean separation between QUIRC sources and their 2MASS counterpart in the frame reaches the minimum. Second, we multiplied each QUIRC frame with a constant value that was derived by the least-square method in order to match the QUIRC photometry with the 2MASS photometry of QUIRC–2MASS counterpart pairs. In this way, we used the 2MASS sources as standard stars. Sources brighter than 11 mag or fainter than 16 mag were not used for this procedure because of the unguaranteed linearity of QUIRC or because of the large uncertainty in the 2MASS photometry.

All the frames, which were thus corrected for astrometry and photometry, were combined into three large mosaic images of the J, H, and K bands. The pseudo-color image is shown in Figure 5.9. Some discontinuities are left in the image because of the difficulty in determining the background intensity level, particularly in regions contaminated by diffuse emissions.

5.2.3 Source Extraction and Photometry

Prior to source detection, all the mosaic images were binned with neighboring 2×2 pixels and smoothed with a Gaussian function to attain better signal-to-noise ratio. For the K-band mosaic image thus improved, we extracted NIR sources above 3 σ level. In SExtractor, we can choose any convolution masks $\phi(x, y)$ in peak finding procedures to increase detectability of faint sources. The two dimensional Gaussian mask $\phi_G(x, y)$ is routinely used, where

$$\phi_G(x,y) = \frac{1}{2\pi\sigma_x\sigma_y} \exp\left(-\frac{1}{2\sigma_x^2} - \frac{1}{2\sigma_y^2}\right). \tag{5.2}$$

However, we adopted the Mexican hat function $\phi_{MH}(x, y)$ for the mask; i.e.,

$$\begin{aligned} \phi_{MH}(x,y) &= \left[\left(x\frac{\partial}{\partial x}+1\right)+\left(y\frac{\partial}{\partial y}+1\right)\right]\phi_G(x,y) \\ &= \frac{1}{2\pi\sigma_x\sigma_y}\left(2-\frac{x^2}{\sigma_x^2}-\frac{y^2}{\sigma_y^2}\right)\exp\left(-\frac{x^2}{2\sigma_x}-\frac{y^2}{2\sigma_y}\right). \end{aligned} \tag{5.3}$$

This function has a positive kernel surrounded by a negative annulus. The limited spatial extent of the kernel favors sources with intrinsically or instrumentally broadened size of

[8] See http://tdc-www.harvard.edu/software/wcstools/ for more details.

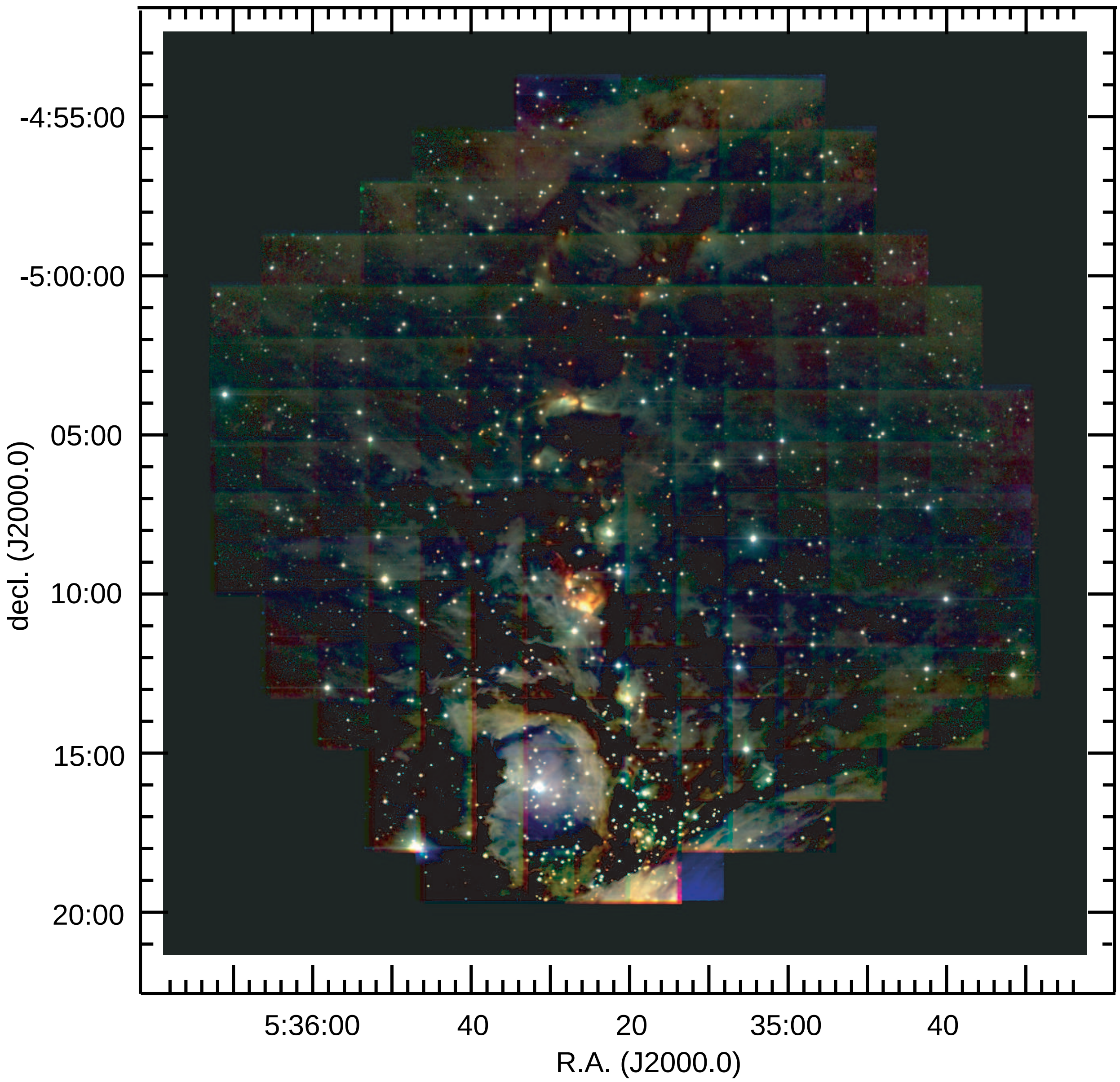


Figure 5.9: Three-color image of QUIRC observations. The neighboring 8×8 pixels are binned. Red, green, and blue are for the *K*-, *H*-, and *J*-band intensity, respectively.

about (σ_x, σ_y) to be detected. We confirmed that a Mexican hat function of the seeing size radius gives the most robust result among several masks with various radii we examined, particularly in regions contaminated with diffuse emissions. With careful visual inspections on the output, we removed (1) sources at the edge of the mosaic images, (2) ghosts of bright sources, and (3) spurious detections (in most cases, diffuse structures were identified as point-like sources). As a result, we picked up 1448 K-band sources.

For each K-band detected source, we derived the J-, H-, and K-band magnitude with the adaptive aperture photometry. For source with less than a 3 σ detection in the J, H, or both bands, we calculated the 3 σ upper limit of their magnitudes.

5.2.4 2MASS Counterpart of QUIRC Sources

We correlated all the QUIRC sources with the 2MASS catalog with the imtmc command in the WCSTools package and found that 692 (∼48%) have the 2MASS counterpart.

5.2.5 Results

Source List

Table B.1 lists the QUIRC sources with their source number, position, J-, H-, and K-band magnitudes, and their 2MASS counterpart. Hereafter, we refer these sources following the IAU convention; e.g., TKK J05342894−0508387 for the source No. 1. Note that all the magnitudes in the list are in the 2MASS color system. For sources that lack the J, H, or both band detections, we listed the 3 σ upper limit of the flux (lower limit of the magnitude) and labeled them with ">". For QUIRC magnitudes brighter than 11 mag, we replaced them with the 2MASS magnitudes with the label "†" if they have the 2MASS counterpart. If they do not have the 2MASS counterpart, we labeled the magnitudes with "<" and recognized them as the lower limit of the flux (the upper limit of the magnitude).

Astrometric and Photometric Accuracy

Using the QUIRC–2MASS counterpart pairs, we evaluated the accuracy of the QUIRC astrometry and photometry in the following way.

In Figure 5.10, we plotted differences between the QUIRC and 2MASS positions both in the R.A. and decl. directions. We found that ΔR.A. = 0.048″±0.161″ (1 σ) and Δdecl. = 0.022″±0.148″ (1 σ), which indicates that the QUIRC positions are determined with the accuracy of the size of a QUIRC pixel (0.189″) and the systematic offset between QUIRC and 2MASS positions is negligible.

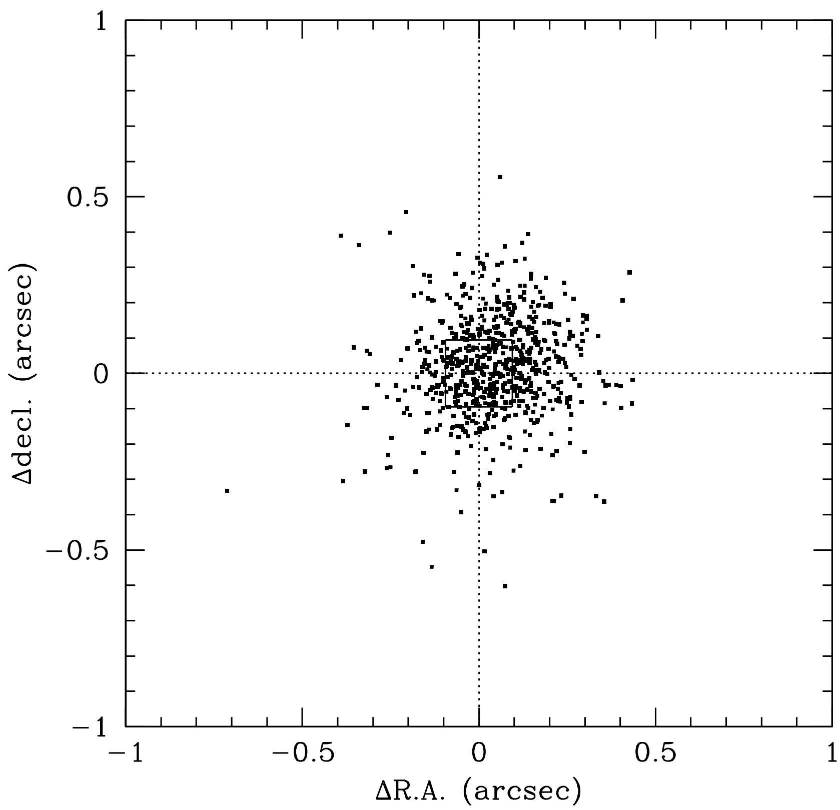


Figure 5.10: Astrometric accuracy of QUIRC sources. The difference of R.A. and decl. in the QUIRC and 2MASS positions (ΔR.A. and Δdecl., respectively) are plotted for each QUIRC–2MASS pair. The solid square at the center represents a QUIRC pixel size of 0.189″×0.189″.

In Figure 5.11, we plotted the QUIRC and 2MASS magnitudes of QUIRC–2MASS counterpart pairs separately for each band. The linear relations (2MASS magnitudes equal QUIRC magnitudes) are violated at the brighter side than ∼11 mag due to the unguaranteed linearity of QUIRC, where its saturation limit of 44000 ADU counts corresponds to ∼11 mag in our observations. For sources fainter than 11 mag in the QUIRC and 2MASS magnitudes, we calculated the difference between these magnitudes for each band (Δm_J, Δm_H, and

Δm_K). We found that $\Delta m_J = -0.06 \pm 0.22$ mag (1 σ), $\Delta m_H = -0.07 \pm 0.19$ mag (1 σ), and $\Delta m_K = -0.04 \pm 0.18$ mag (1 σ), indicating that the QUIRC photometry fainter than 11 mag is consistent with the 2MASS photometry within ∼0.2 mag.

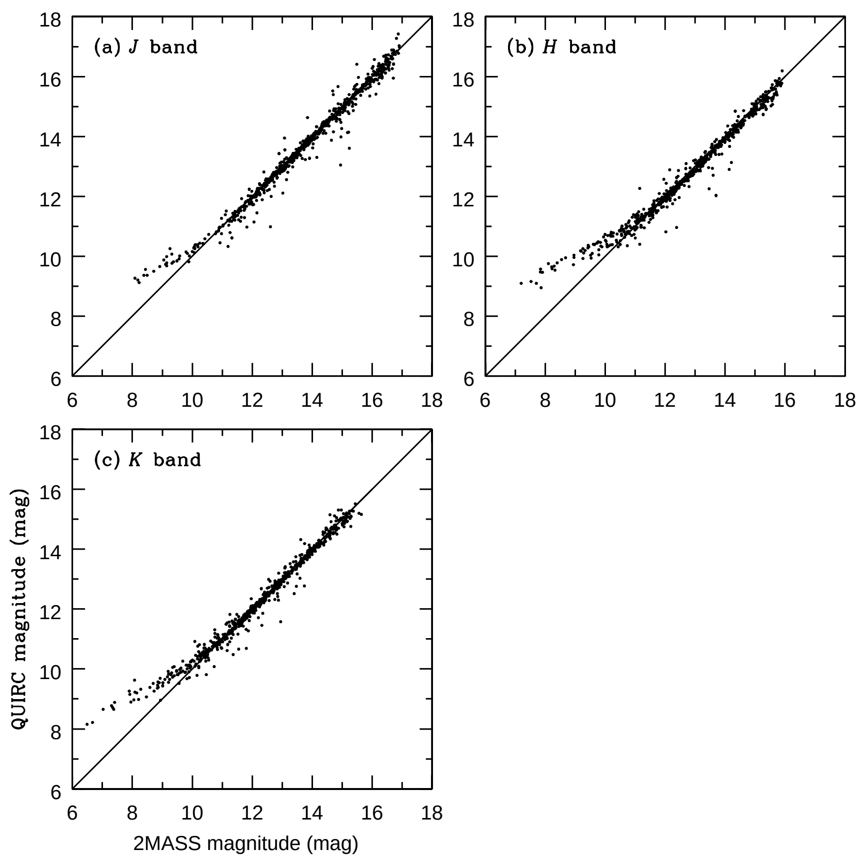


Figure 5.11: Photometric accuracy of QUIRC sources. The 2MASS and QUIRC magnitudes of the counterpart pairs are plotted separately for the (*a*) *J*, (*b*) *H*, and (*c*) *K* bands. Solid lines represent that the 2MASS magnitudes equal to the QUIRC magnitudes).

Survey Depth

We estimated the survey depth of our QUIRC observations in the following manner. First, we embedded 500 artificial sources with 13.0–13.5 mag in the *J*-, *H*-, and *K*-band mosaic images. The same source detection algorithm was employed to detect these artificial sources, then the detection rate of sources with 13.0–13.5 mag was derived. The same procedure was repeated for sources of different magnitudes from 13.0 to 20.0 mag with 0.5 mag bins. The detection rate at each magnitude bin is given in Figure 5.12 for the K (*solid*), H (*long-dashed*), and J (*short-dashed*) band, respectively. The 90% completeness limit (3 σ) was thus estimated to be $K \sim 16.0$ mag, $H \sim 16.5$ mag, and $J \sim 17.5$ mag.

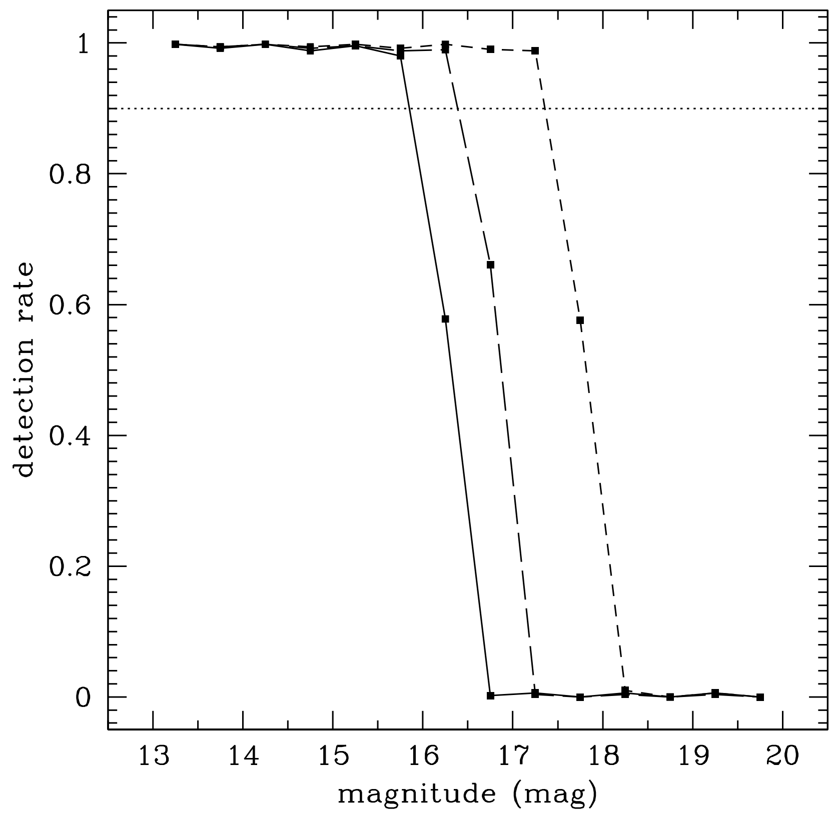


Figure 5.12: Completeness limit of QUIRC observations. The fractions of detected artificial sources are shown with the solid (K), long-dashed (H), and short-dashed (J) lines. The 90% (*dotted line*) completeness limit is $K \sim 16.0$ mag, $H \sim 16.5$ mag, and $J \sim 17.5$ mag.

In order to compare the depth of our observation with that of 2MASS, we made a histogram of the number counts of the K-band detected sources at each magnitude (Fig. 5.13). The short- and long-dashed histograms respectively represent the number of the QUIRC sources with and without the 2MASS counterpart, while the total is given in the solid histogram. The 2MASS catalog fails to detect some bright sources even in $K < 14$ mag, which is mainly due to the contamination by diffuse emission or to their binarity. This was confirmed by plotting the 2MASS source list on the 2MASS image.

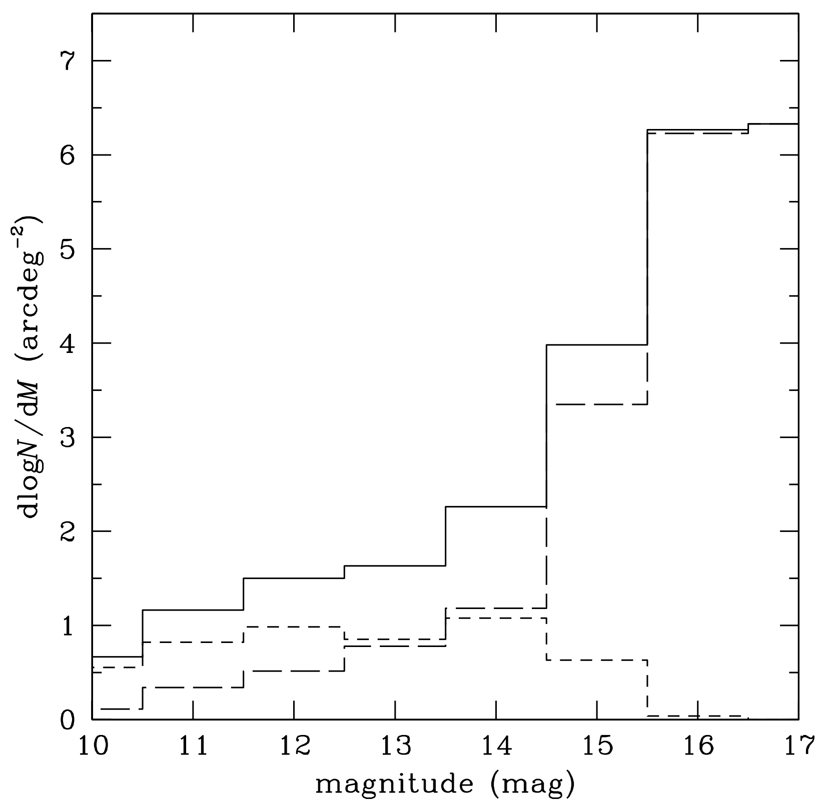


Figure 5.13: Source counts ($d\log N/dM$) of QUIRC K-band detected sources, where M is the magnitude and N is the number of sources fainter than M mag. The solid histogram is the counts of all the QUIRC K-band sources, while the short- and long-dashed histograms are those with and without the 2MASS counterpart, respectively.

5.2.6 QUIRC Counterpart of *Chandra* Sources

Among 237 *Chandra* sources with the 2MASS counterpart, 30 sources lack either 2MASS J-, H-, or K_s-band detections. Together with 161 *Chandra* sources that have no 2MASS counterpart, we hereafter call them the "2MASS-unIDed" *Chandra* sources. There are 183 and 9 2MASS-unIDed sources in the ACIS-I and ACIS-S2 FOVs. Using the QUIRC source list (Table B.1), we searched for the NIR counterpart of 183 2MASS-unIDed ACIS-I sources with a more elaborated procedure as follows.

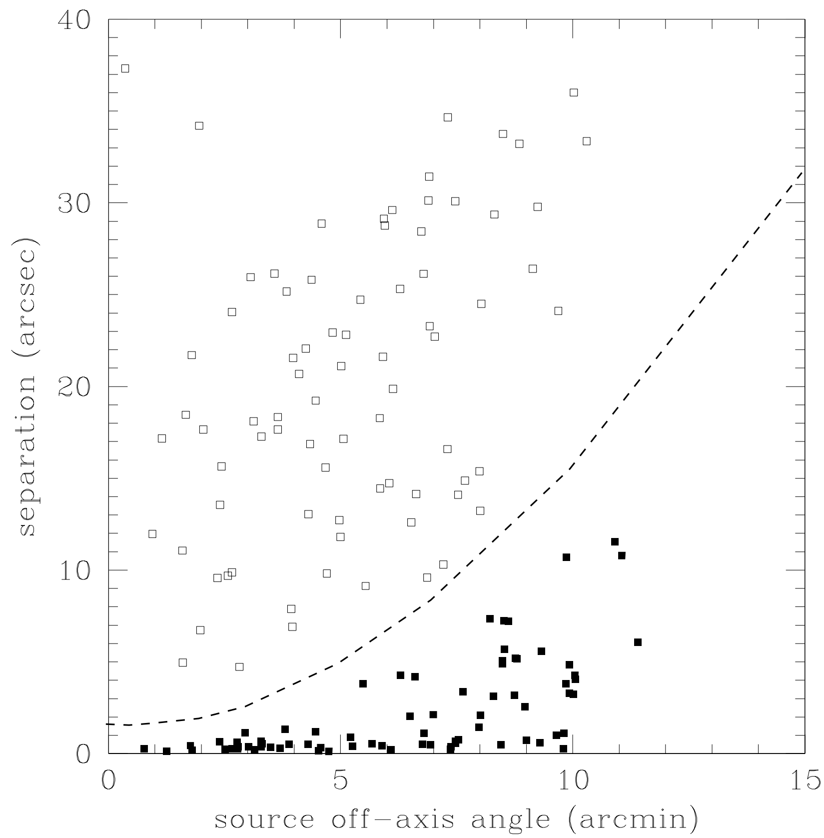


Figure 5.14: Closest QUIRC–*Chandra* pairs with their separation and the X-ray off-axis angle on the vertical and the horizontal axis, respectively. The dotted curve is the 90% encircled energy radius (r_{90}) of 1.49 keV X-rays at a given off-axis angle. Filled squares, which are below the dotted curve, are recognized as the counterpart pairs. Open squares, which are above the dotted curve, are recognized as non-associated pairs.

The *Chandra* PSF radii differ by more than an order of magnitude between on-axis sources and those at the field edge, which deteriorates the position accuracy of sources at large off-axis angles. Therefore, in identifying the QUIRC counterpart of the *Chandra* sources, we took the off-axis angle into account. First, we searched for the QUIRC source closest to each *Chandra* source. Second, we conversely searched for the *Chandra* source closest to each QUIRC source. Then, we picked up 159 QUIRC–*Chandra* pairs that are the closest to each other. These pairs include those of physically associated ("counterpart pairs") and those of no physical association ("non-associated pairs"). In Figure 5.14, we plotted the separation between the closest QUIRC–*Chandra* pairs as a function of the off-axis angle of the *Chandra* source. We showed two groups by the filled and open squares, which are well separated by the dotted curve indicating the 90% encircled energy radius (r_{90}; ~ 0.9 times FWHM of a Gaussian PSF) of 1.49 keV X-rays as a function of the off-axis angle[9]. Since r_{90} is the radius in which the 90% of incident X-ray photons are accumulated, it also represents the position accuracy of the *Chandra* sources at each off-axis angle. We therefore regard the closest QUIRC–*Chandra* pairs with the separation angle less than r_{90} (*filled squares*) to be counterpart pairs, while those with larger separation angle of more than r_{90} (*open squares*) as non-associated pairs. As a result, 75 2MASS-unIDed *Chandra* sources were newly found to have the QUIRC counterpart (hereafter we call them "QUIRC-IDed" sources). The 2MASS and QUIRC counterpart of *Chandra* sources and their NIR colors are given in Table A.1. Note that the NIR colors are converted to the CIT color system (Sect. 6.2).

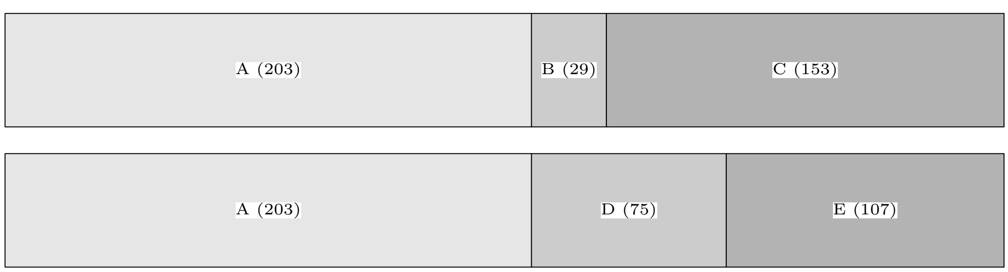


Figure 5.15: NIR identifications of *Chandra* ACIS-I sources. (*top*) 2MASS identifications. Those with the 2MASS J-, H-, and K_s-band detections ("2MASS-IDed") are in A, those lack at least one detection in these three bands are in B, and those with no 2MASS counterpart are in C. B and C are called "2MASS-unIDed" sources, for which the QUIRC counterpart is searched. (*bottom*) QUIRC identifications. Those with the QUIRC counterpart ("QUIRC-IDed") are in D, while those without 2MASS nor QUIRC counterpart ("NIR-unIDed") are in E. 2MASS-IDed and QUIRC-IDed sources (A and D) are collectively called "NIR-IDed" sources. The number of sources are shown in parentheses.

[9]See http://asc.harvard.edu/udocs/docs/POG/MPOG/index.html.

Chapter 6

NIR-IDed X-ray Sources: (1) NIR Properties

Contents

In this chapter, the NIR properties of the QUIRC sources and NIR-IDed X-ray sources are discussed. Using color-color and color-magnitude diagrams of the J, H, and K bands, ample information on physical parameters of YSOs can be obtained such as the evolutional class, disk existence, stellar mass, bolometric luminosity and amount of extinction. In Sect. 6.1, we examine the NIR colors of all QUIRC sources and describe the overall features of OMC-2 and OMC-3. In Sect. 6.2, we zoom in to the NIR-IDed X-ray sources (=NIR sources with the X-ray counterpart), and estimate their mass and evolutional class to classify them into several groups.

6.1 NIR Sources

6.1.1 *K*-band Luminosity Function

We made a histogram of the K-band-detected QUIRC sources at each K-band magnitude (K-band luminosity function; KLF) in Figure 6.1. The short- and long-dashed histograms respectively represent the KLF of QUIRC sources with and without the 2MASS counterpart, while the total is given in the solid histogram. The dashed-and-dotted histogram is the KLF of NIR-IDed X-ray sources normalized to the area of QUIRC FOV. In the KLF, we can estimate the back- and foreground source contamination. The dotted curve shows the back- and foreground source counts predicted by a Galactic star count model (the SKY model; Cohen 1993[38]; Cohen 1994[39]; Cohen, Sasseen, & Bowyer 1994[40]; Cohen 1995[41]) assuming no interstellar extinction. There is an overall extinction of $A_K \sim 1$ mag in this region (Fig. 6.4; see the relevant discussion in the text), which shifts the dotted curve rightward by $\sim$1 mag to the dashed curve. The contamination is not negligible for sources fainter than $K \sim 15$ mag. For brighter sources, however, we can assume that most of them are cloud members.

The peak of KLF for all the QUIRC sources (*solid histogram*) at 12–13 mag is a real feature when the completeness limit ($K \sim 16$ mag) and back- and foreground contamination are taken into consideration. This has been inferred by Jones et al. (1994)[101], who studied the KLF of OMC-2 sources. The peak of KLF is often seen in the young associations like the Orion Nebula Cluster (Ali & DePoy 1995[2]; Hillenbrand & Carpenter 2000[86]) and can be explained in terms of difference in the mass-to-luminosity relation between main sequence and pre-main-sequence sources (Muench, Lada, & Lada 2000[137]). The peak magnitude is consistent with the cloud age of ∼1 Myr (Ali & DePoy 1995[2]), which was also confirmed by other methods (Sect. 3.1).

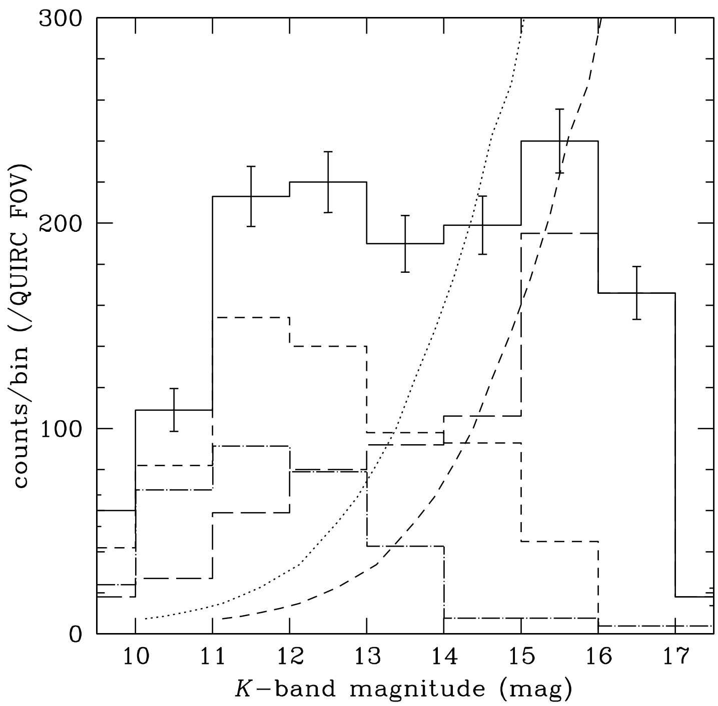


Figure 6.1: QUIRC K-band-detected source counts at each magnitude in the 2MASS color system. The solid histogram is the counts of all the QUIRC K-band sources, while the short- and long-dashed histograms are those with and without the 2MASS counterpart, respectively. The dashed-and-dotted histogram is the KLF of the NIR-IDed X-ray sources normalized to the area of QUIRC FOV. The uncertainty of $\sqrt{\text{counts bin}^{-1}}$ is given for the solid histogram. The dotted curve shows the back- and foreground source counts predicted by a Galactic star count model assuming no extinction. When a uniform extinction of 1 mag is assumed, this curve should be shifted rightward by 1 mag to the dashed curve.

6.1.2 Color-color Diagram of QUIRC Sources

We first converted the J-, H-, and K-band magnitudes in the 2MASS color system in Table B.1 ($J_{\rm 2MASS}$, $H_{\rm 2MASS}$, and $K_{\rm 2MASS}$) to the CIT color system ($J_{\rm CIT}$, $H_{\rm CIT}$, and $K_{\rm CIT}$) using the following conversion formula (Carpenter 2001[27]):

$$J_{\rm CIT} = 0.947 J_{\rm 2MASS} + 0.053 K_{\rm 2MASS} + 0.036 \qquad (6.1)$$

$$(J-H)_{\rm CIT} = 0.929(J-H)_{\rm 2MASS} + 0.040 \tag{6.2}$$

$$(H-K)_{\rm CIT} = 0.975(H-K)_{\rm 2MASS} - 0.027 \tag{6.3}$$

The color-color diagram (Lada & Adams 1992[114]) was made using all QUIRC sources with significant J-, H-. and K-band detections (1305 out of 1448 sources). Figure 6.2 shows the $(J–H)/(H–K)$ diagram, where QUIRC sources with and without the 2MASS counterpart are respectively shown with open and filled symbols. The intrinsic colors of giants and dwarfs are given by the thick solid curves (Tokunaga 2000[185]). The emission from the circumstellar disks of cTTSs gives NIR excess on the giant and dwarf colors, hence cTTSs are aligned from bottom left to top right along the cTTS locus in the thick solid line (Meyer, Calvet, & Hillenbrand 1997[130]). These colors change their position along the reddening vector as increasing interstellar and circumstellar medium. Therefore, sources between the right and middle reddening lines are reddened cTTSs (*diamonds*), while those between the middle and left reddening lines are reddened wTTSs and some fraction of cTTSs (*triangles*). Sources located to the right of the right reddening line are surrounded by extended envelopes in addition to the disks, hence they are reddened more than disks alone (Strom, Kepner, & Strom 1995[177]). These sources (*hexagons*) are classified to be class I protostars.

Cares should be taken that back- and foreground sources are included in our NIR source samples. However, sources with NIR excess (protostars and cTTSs; hexagons and diamonds in Fig. 6.2) can be safely assumed to be YSOs, namely cloud members.

6.1.3 Color-magnitude Diagram of QUIRC Sources

In order to estimate the stellar mass and extinction of NIR sources, we employed the $J/(J–H)$ color-magnitude diagram (Fig. 6.3), where QUIRC sources with significant J- and H-band detections are shown in squares. Among them, filled squares are those with NIR excess. All sources are assumed to be at the distance of 450 pc. The intrinsic colors of sources at the age of 1 Myr are shown in the solid curves for the mass range of 0.002–1.4 $M_\odot$ (Baraffe et al. 1998[16]) and 1.4–7.0 $M_\odot$ (Siess, Dufour, & Forestini 2000[171]). As a result of the extinction, these intrinsic colors are reddened in the direction of the reddening vector at the top right of Figure 6.3. By moving the position of squares backward along the reddening vector to the 1 Myr isochrone curves, the mass and amount of extinction (A_V) can be estimated for each source.

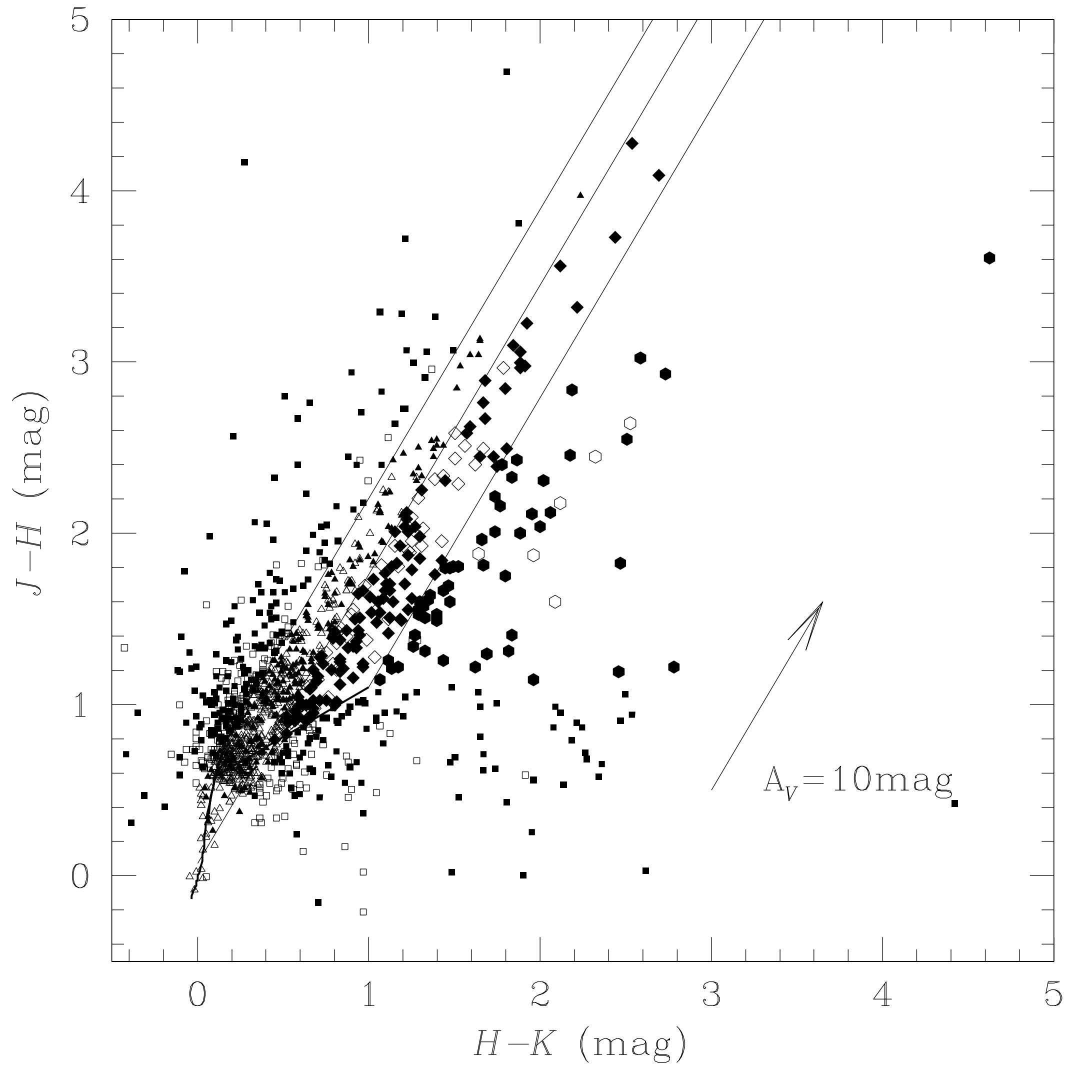


Figure 6.2: Color-color diagram of QUIRC sources with significant J-, H- and K-band detections. Open and filled symbols are QUIRC sources with and without 2MASS counterpart, respectively. Hexagons and diamonds are classified as protostars and cTTSs, while triangles are the mixture of main sequence, wTTSs, and cTTSs. Squares are in none of these classifications. The intrinsic colors of dwarfs and giants are given with thick solid curves, while the cTTS locus is with the thick solid line. The arrow at the bottom right gives the reddening vector of A_V =10 mag. The slope of the reddening lines is assumed to be $E(J\text{–}H)_{\text{reddening}}/E(H\text{–}K)_{\text{reddening}} = 1.69$ (Meyer et al. 1997[130]). The typical uncertainty of colors is roughly ±0.1 mag.

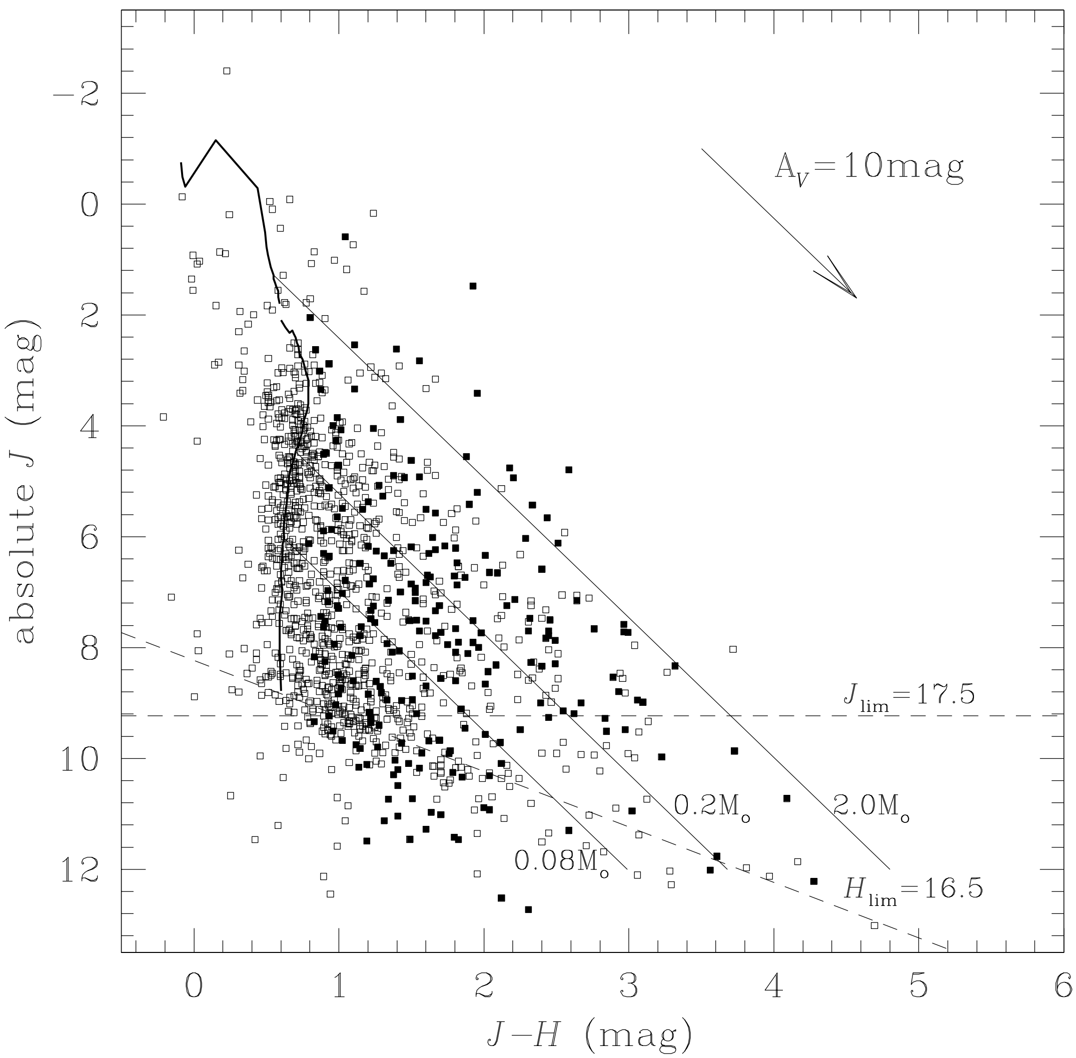


Figure 6.3: Color-magnitude diagram of QUIRC sources with significant J- and H-band detections (*squares*). Among them, filled are those with NIR excess. The 1 Myr isochrone curves (*thick solid curves*) are from Baraffe et al. (1998)[16] for $0.002\ M_\odot \leq M \leq 1.4\ M_\odot$ and Siess et al. (2000)[171] for $1.4\ M_\odot \leq M \leq 7.0\ M_\odot$. Two dashed lines show the detection limit of J = 17.5 mag and H = 16.5 mag. The arrow at the top right indicates the reddening vector of A_V =10 mag. Reddening lines for 2.0 $M_\odot$, 0.2 $M_\odot$, and 0.08 $M_\odot$ are given with solid lines. The typical uncertainty of colors and magnitudes is roughly ±0.1 mag.

To examine the distribution of NIR extinction, we made the histogram of the K-band extinction (A_K) in Figure 6.4. The A_K value of each source was converted from A_V using $A_K/A_V = 0.125$ (Mathis 2000[123]). Sources in the mass range of $0.2\ M_\odot \leq M \leq 2.0\ M_\odot$ and $J > 17.5$ mag (Fig. 6.3) are used to construct the histogram. The observational bias in favor of low extinction sources is thus avoided for $A_V < 20$ mag ($A_K < 2.5$ mag). The typical A_K was estimated to be ~1 mag.

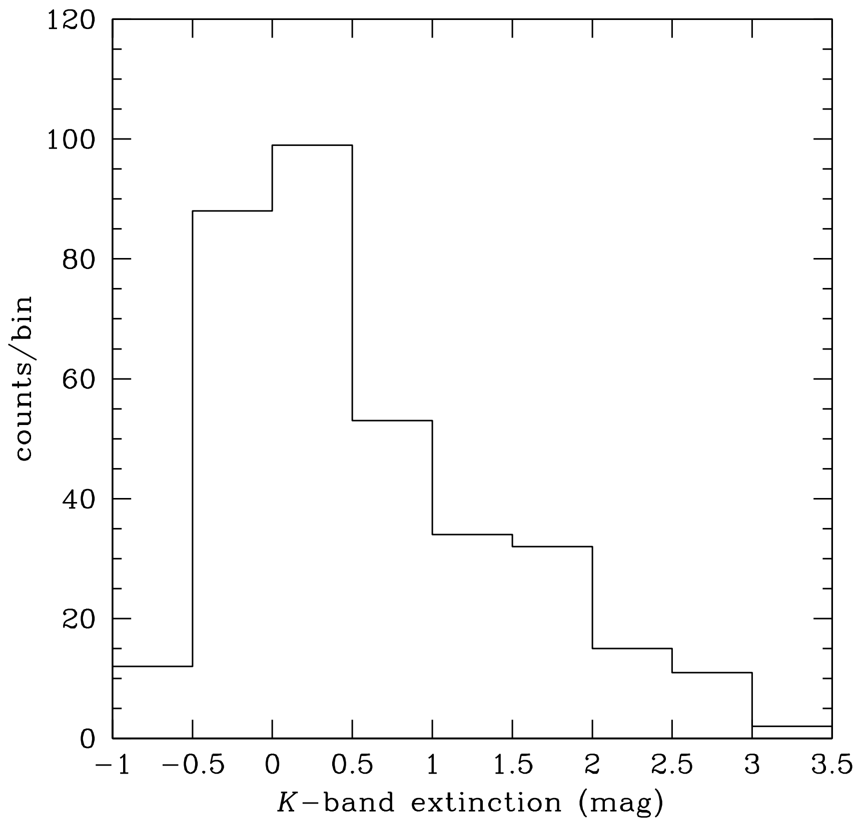


Figure 6.4: Histogram of the K-band extinction (A_K) derived from NIR sources in the mass range of $0.2\ M_\odot < M < 2.0\ M_\odot$ and $J > 17.5$ mag. The A_V value of each source was estimated from Figure 6.3, and converted to A_K using $A_K/A_V = 0.125$ (Mathis 2000[123]). The histogram was fitted with an exponential function in the positive A_K region to derive the typical value of $A_K \sim 1$ mag.

6.2 The NIR-IDed X-ray Sources

6.2.1 Cloud Membership

Krishnamurthi et al. (2001)[108] observed the core of the Pleiades star cluster with *Chandra* for 36 ks and found that a significant fraction of X-ray sources are likely to be AGNs. We examined the contamination of our X-ray samples by AGNs and concluded that there is only a negligible number of AGNs among NIR-IDed X-ray sources for the following reason.

The number count of galaxies (N) per square degree per magnitude at a certain K-band magnitude (K) is given by

$$\frac{dN}{dK} = 4000 \times 10^{\alpha(K-17)}, \tag{6.4}$$

where $\alpha = 0.67$ for 10 mag $< K <$ 17 mag (Tokunaga 2000[185]). The number of galaxies in the range of K_{min} mag $< K < K_{max}$ mag is then estimated by

$$\int_{K_{min}}^{K_{max}} \frac{dN}{dK}\, dK = \frac{4000}{\alpha \log 10}\{10^{\alpha(K_{max}-17)} - 10^{\alpha(K_{min}-17)}\}. \tag{6.5}$$

The NIR counterparts of the *Chandra* sources have a K-band magnitude in 6 mag $< K <$ 15 mag. We substituted $K_{min} = 6$ and $K_{max} = 15$ for simplicity, though the equation (6.5) is valid only for $K_{min} > 10$ mag. Still, this gives us a good estimate since the second term of the right-hand side is negligible compared to the first term in this case. Considering the *Chandra* FOV, the estimated number of galaxies in the range of 6 mag $< K <$ 15 mag is ∼12. This is only ∼0.8% of all the NIR sources in the same magnitude range. Moreover, the background galaxies in this direction suffer a significant extinction due to dense ISM, which makes the contribution of extragalactic sources to our NIR samples even smaller. This excludes the possibility of NIR-IDed X-ray sources to be extragalactic AGNs.

The KLF of NIR-IDed X-ray sources (Fig. 6.1) indicates that most of the X-ray sources have the K-band magnitude of $K <$ 14 mag, where the contamination by back- and foreground galactic sources is negligible. We hereafter regard all the NIR-IDed X-ray sources as cloud members.

6.2.2 Evolutional Class Estimates

Table A.1 lists the J-band magnitude and (J–H) and (H–K) colors of NIR-IDed *Chandra* sources in the CIT color system. Using these colors, we first made the (J–H)/(H–K)

color-color diagram (Fig. 6.5), where NIR-IDed *Chandra* sources are plotted. Filled and open symbols are QUIRC-IDed and 2MASS-IDed *Chandra* sources. The arrows on symbols indicate the lower limit of colors; i.e., the upward arrows are due to the lack of the J-band detections and the rightward arrows are due to the saturation in the K band. Following the classification scheme addressed in the previous section, we picked up protostars, cTTSs, and wTTSs (*hexagons*, *diamonds*, and *triangles*). Some sources can not be classified in any of these classes either because they are out of the three regions (*squares*) or they have no J-, H-, or both band detections.

We have to keep in mind that the K-band magnitude is not always sensitive enough to detect excess emissions from circumstellar disks. Recent *Infrared Space Observatory* observations on R Coronae Australis, Chamaeleon, and ρ Ophiuchi dark clouds indicate that the classification based only on the J, H, and K-band data underestimates the number of YSOs with disks (Olofsson et al. 1999[144]; Persi et al. 2000[152]; Bontemps et al. 2001[24]). Haisch et al. (2001)[76] combined the MIR with NIR photometry on NGC 2024 and reported that about one-third of class II sources are recognized not to have excess emissions with the $(J–H)/(H–K)$ diagram alone, hence are classified as class III sources.

In classifying the NIR-IDed X-ray sources, therefore, we additionally used H_α emission (Herbig & Bell 1988[83]) and UV excess (Rebull et al. 2000[160]) data to complement the NIR excess data. The H_α emission line is directly related to disks and its equivalent width is used to discriminate cTTSs (>10 Å) and wTTSs (<10 Å). The NIR and UV excess data work complementarily. The former is sensitive to disks of earlier-type sources with higher photospheric temperature, while the latter is sensitive to later-type sources with lower temperature, because of the contrast between the photosphere and disk colors. The UV excess of cTTSs is considered to originate from the boundary layer of an accretion disk (Hartigan et al. 1991[79]), making it another indicator of disk existence.

Among sources classified as wTTSs in the $(J–H)/(H–K)$ color-color diagram, four (I13, I248, I324, and I334) and eleven (I111, I122, I150, I200, I237, I308, I319, I322, I328, I332, and S4) sources are reclassified into cTTSs based respectively on their H_α and UV excess emissions. Consequently, we identified 13 protostars, 59 cTTSs, and 170 wTTSs among 278 NIR-IDed X-ray sources. The result of the classification is summarized in Table A.1.

We examined the position of *Chandra* sources separately for each class (Fig. 6.6). Protostars are clearly concentrated along the 1.3 mm ridge. Since cores seen in the millimeter continuum are the sites of on-going star formation, this reinforces the idea that they are

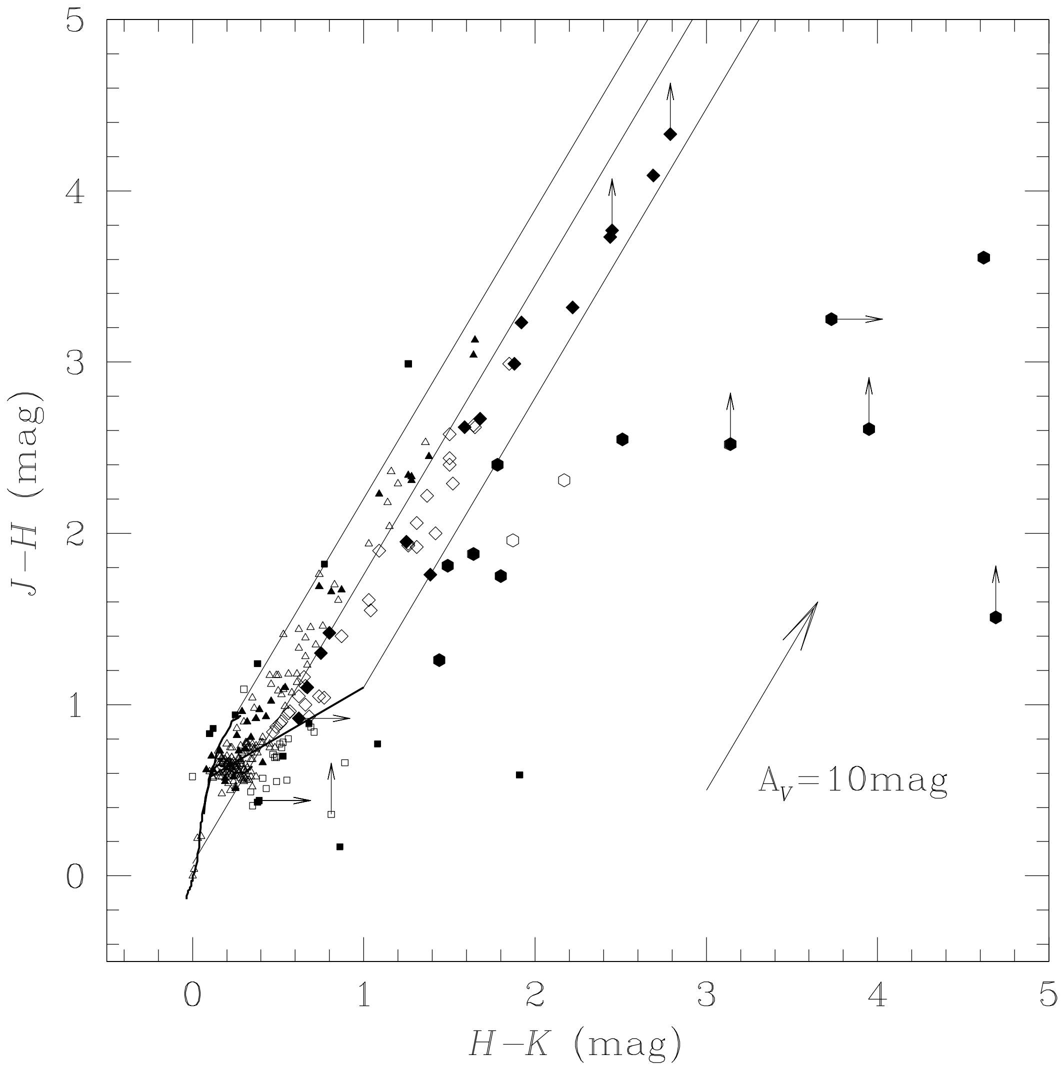


Figure 6.5: Color-color diagram of NIR-IDed X-ray sources. The 2MASS-IDed and QUIRC-IDed *Chandra* sources are plotted with open and filled symbols, respectively. Hexagons, diamonds, and triangles are classified to be protostars, cTTSs, and wTTSs, while squares are in none of these classes. The intrinsic colors of dwarfs and giants are given with solid curves, while the cTTS locus is with solid line. The arrow at the bottom right gives the reddening vector of A_V =10 mag. The slope of the reddening lines is assumed to be $E(J–H)_{\rm reddening}/E(H–K)_{\rm reddening}$ = 1.69 (Meyer et al. 1997[130]). The typical uncertainty is roughly ±0.1 mag.

YSOs at the very early stage. Most of protostars also accompany apparent NIR nebulosity in the K-band image, which indicates that they are deeply embedded sources. The spatial distribution of cTTSs is also correlated with the ridge, although wTTSs and X-ray sources with no NIR counterpart (Fig. 6.8) are not.

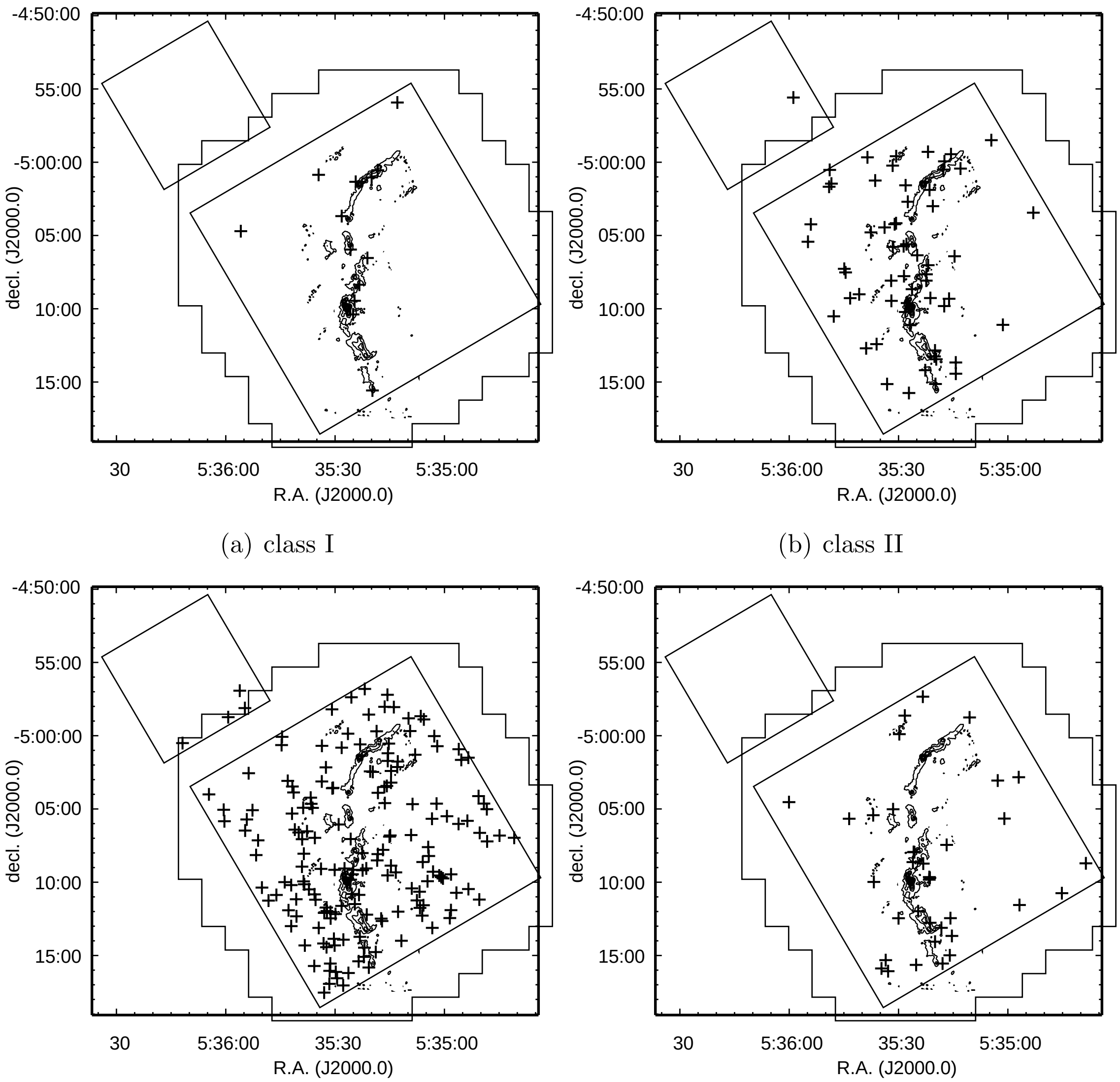
6.2.3 Mass and Bolometric Luminosity Estimates

We next made the $J/(J–H)$ color-magnitude diagram (Fig. 6.7), where NIR-IDed *Chandra* sources are shown in squares. Among them, filled squares are those with NIR excess. Squares with the rightward arrows indicate the saturation in the H band, and those with both the rightward and the downward arrows have H- but not J-band detection.

We categorized the NIR-IDed X-ray sources into four groups based on their mass. Using this diagram, we can estimate the mass with the J-band magnitude because nearly all the emission in this band is photospheric for late-type pre-main-sequence stars (Strom & Strom 1994[176]). Magnitudes in the longer wavelengths are affected by NIR excess emissions from circumstellar dust, while those in the shorter wavelengths are by UV excess emissions from the disk-boundary layer, giving a significant overestimate of the mass (Gagné et al. 1995[61]).

Using the theoretical calculations of isochrone curves by Baraffe et al. (1998)[16] for sources in 0.002 $M_{\odot} \leq M \leq 1.4\ M_{\odot}$ and Siess et al. (2000)[171] in 1.4 $M_{\odot} \leq M \leq 7.0\ M_{\odot}$, we estimated the mass (Table A.1), which is quantized depending on the grid of these computations. The bolometric luminosity is also derived for each source in the same way.

Among 278 NIR-IDed X-ray sources, 268 have significant J- and H-band detections. These sources are separated into high mass (HM) with $M > 10\ M_{\odot}$, intermediate mass (IM) with 10 $M_{\odot} > M > 2.0\ M_{\odot}$, low mass (LM) with 2.0 $M_{\odot} > M > 0.2\ M_{\odot}$, and very low mass (VLM) sources with $M < 0.2\ M_{\odot}$. For sources with $M > 4.0\ M_{\odot}$, where the J-band magnitude is insensitive to the mass, we assumed that they have $M = 4.0\ M_{\odot}$ except for I242, which we classified into HM based on its spectroscopy observations (Kukarkin et al. 1971[111]). As a consequence, we found 1 HM, 21 IM, 139 LM, and 107 VLM sources. The spatial distributions of these sources are given separately for each mass range in Figure 6.8.



(a) class I

(b) class II

(c) class III

(d) others

Figure 6.6: Spatial distribution of *Chandra* sources (*pluses*) separately for (*a*) class I (protostars), (*b*) class II (cTTSs), (*c*) class III (wTTSs), and (*d*) NIR-IDed sources that are classified into none of these three classes. The FOVs of ACIS and QUIRC are shown with solid lines. The contours in each panel are the 1.3 mm intensity (Chini et al. 1997[35]).

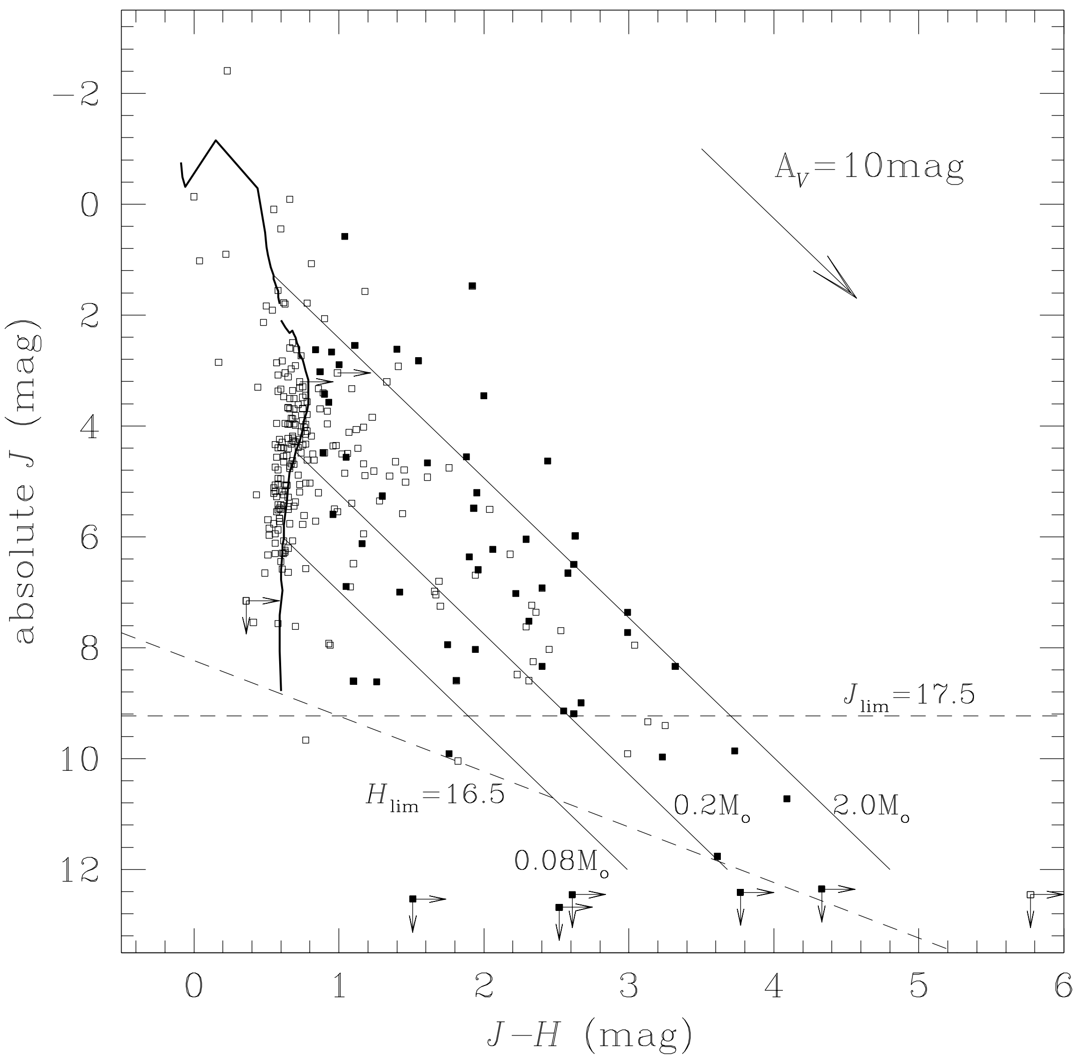
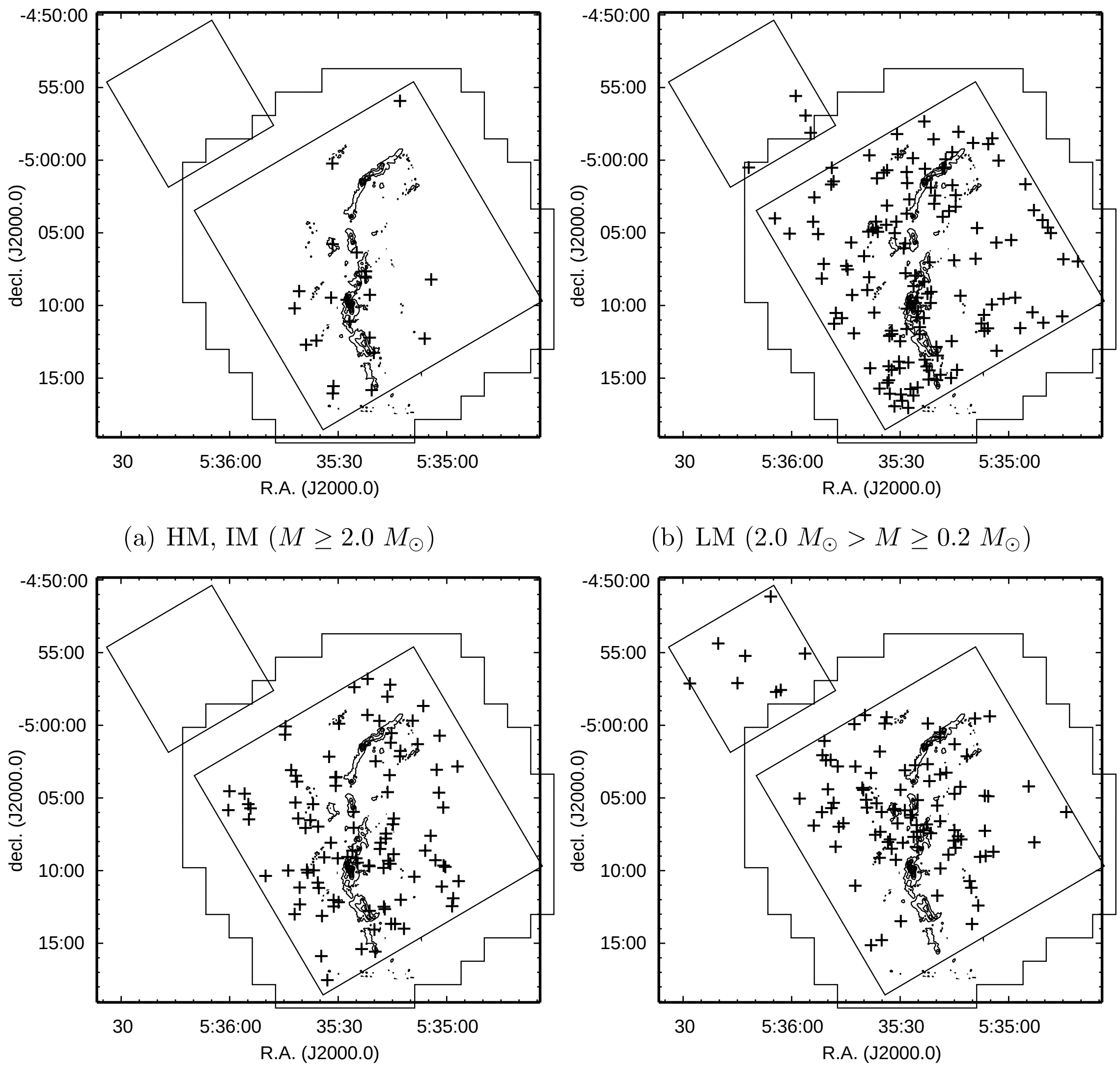


Figure 6.7: Color-magnitude diagram of NIR-IDed *Chandra* sources (*squares*). Among them, filled ones are those with NIR excess. The 1 Myr isochrone curves (*thick solid curves*) are from Baraffe et al. (1998)[16] for 0.002 $M_{\odot} \leq M \leq 1.4\ M_{\odot}$ and from Siess et al. (2000)[171] for 1.4 $M_{\odot} \leq M \leq 7.0\ M_{\odot}$. Two dashed lines show the detection limit of $J = 17.5$ mag and $H = 16.5$ mag. The arrow at the top right indicates the reddening vector of A_V =10 mag. The reddening lines for 2.0 $M_{\odot}$, 0.2 $M_{\odot}$, and 0.08 $M_{\odot}$ are given with solid lines. The typical uncertainty is roughly ±0.1 mag.



(a) HM, IM ($M \geq 2.0\ M_\odot$)

(b) LM ($2.0\ M_\odot > M \geq 0.2\ M_\odot$)

(c) VLM ($M < 0.2\ M_\odot$)

(d) NIR-unIDed

Figure 6.8: Spatial distribution of *Chandra* sources (*pluses*) separately for (*a*) HM and IM ($M \geq 2.0\ M_\odot$) (*b*) LM ($2.0\ M_\odot > M \geq 0.2\ M_\odot$), (*c*) VLM ($M < 0.2\ M_\odot$), and (*d*) NIR-unIDed X-ray sources. The FOVs of ACIS and QUIRC are shown with solid lines. The contours in each panel are the 1.3 mm intensity (Chini et al. 1997[35]).

Chapter 7

NIR-IDed X-ray Sources: (2) X-ray Properties

Contents

In this chapter, we discuss the X-ray properties of NIR-IDed *Chandra* sources. In order to concentrate on sources with enough statistics, we deal with bright ACIS-I sources in this and the following chapters. In Sect. 7.1, we conduct temporal analysis and illustrate that about half of them show flux variability including flares. In Sect. 7.2, we perform spectral analysis. By fitting with one-temperature thin-thermal plasma model, we derive the plasma temperature ($k_{\rm B}T$), X-ray luminosity ($L_{\rm X}$), and the amount of absorption ($N_{\rm H}$). When rejected, we apply two-temperature plasma model to fit the spectra. About 90% of sources are fit with either by one-temperature or two-temperature plasma model. Finally, in Sect. 7.3, we examine the relations among these parameters.

7.1 Temporal Analysis

One of the most notable characteristics of the X-ray emissions from YSOs is rapid variability of their X-ray light curves. Flare episodes of fast rise and slow decay are often observed. In order to pick up X-ray sources with variability, we applied a simple temporal analysis for the *Chandra* sources. The XRONOS package[1] was used for the following procedures.

We concentrated our temporal analysis on bright ACIS-I sources so that we can ignore the backgrounds. Figure 7.1 shows the distribution of X-ray counts (0.5–8.0 keV) and S/N, where S and N are the source and background counts in the 0.5–8.0 keV range. The background counts were derived using the values in Table 4.1 and normalized by the area of source accumulation region and the exposure time. We picked up sources with more counts than 100 and higher S/N than 10 as the samples for temporal analysis. We hereafter call them "bright (T)" sources, while the remaining sources are called "faint (T)" sources. Among 278 NIR-IDed ACIS-I sources, 120 are bright (T) and 158 are faint (T).

The X-ray counts of bright (T) sources were binned with three binning sizes (100 s bin^{-1}, 1000 s bin^{-1}, and 10000 s bin^{-1}) to construct light curves. These light curves were fitted with a constant flux model. The χ^2 value was derived for each fitting with

$$\chi^2 = \sum_{i=1}^{N} \left(\frac{F_{data}^{(i)} - F_{model}^{(i)}}{\Delta F_{data}^{(i)}} \right)^2 , \tag{7.1}$$

where $F_{data}^{(i)}$ and $F_{model}^{(i)}$ are the count rate of the data and the model in the i'th time bin, and $\Delta F_{data}^{(i)}$ is the uncertainty of $F_{data}^{(i)}$. The χ^2 value follows the chi square distribution of

[1] See http://heasarc.gsfc.nasa.gov/docs/xanadu/xronos/xronos.html.

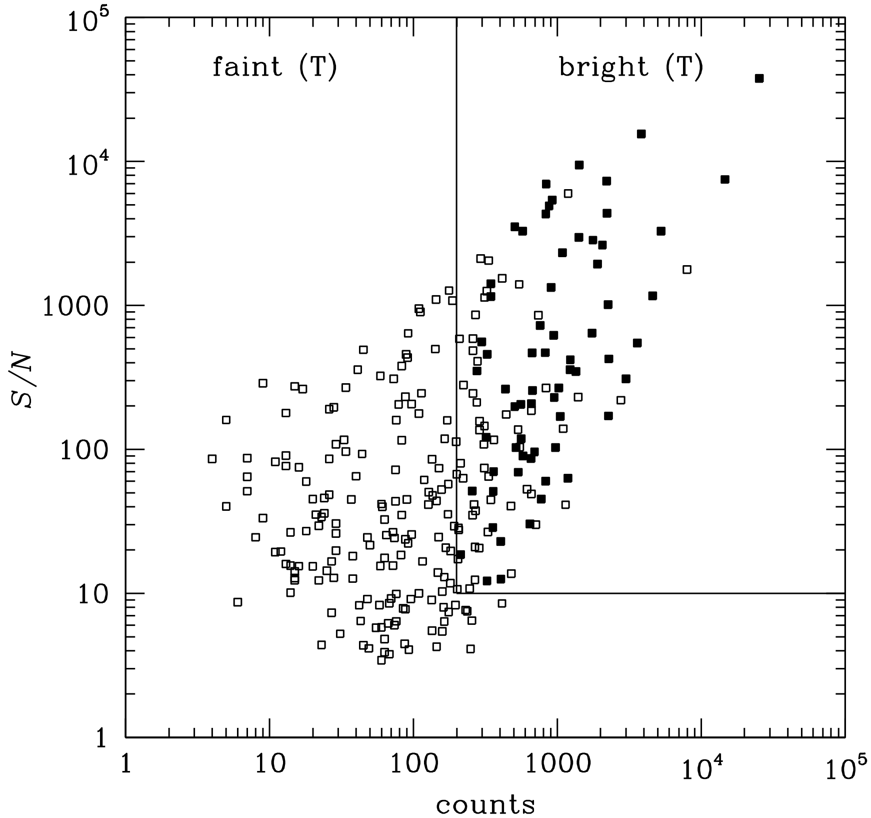


Figure 7.1: X-ray counts (0.5–8.0 keV) and S/N (0.5–8.0 keV) of NIR-IDed ACIS-I sources. Sources with more counts than 200 and higher S/N than 10 (120 sources) are bright (T) sources, for which the temporal analysis was conducted. The remaining 158 sources are faint (T) sources. Variable sources among bright (T) sources are marked filled.

$(N-1)$ degrees of freedom; $\chi^2_{N-1}(x)$. The upper probability for the null hypothesis (α) was thus calculated as

$$\alpha = \int_{\chi^2}^{\infty} \chi^2_{N-1}(x)dx. \tag{7.2}$$

We recognized the null hypothesis (the constant count rate) is rejected; i.e., the light curve is variable, if $\alpha < 0.01$ in at least one binning size. Consequently, we picked up 66 (55%) variable light curves among bright (T) sources.

Figures 7.2–7.6 show the background-unsubtracted light curves of all the variable X-ray sources. Note that the background count rate is less than the total count rate by more than 10 times. Many of the light curves show typical flare-like events.

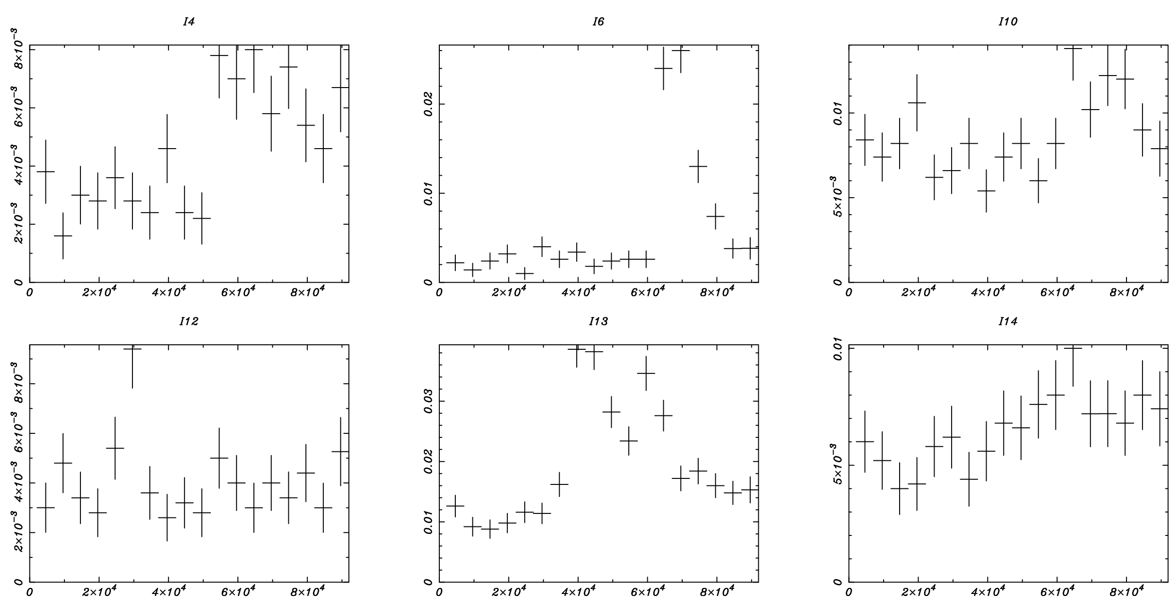


Figure 7.2: Light curves of variable bright (T) NIR-IDed *Chandra* sources (I4–I14) with the binning of 5000 s bin^{-1}. Light curves are plotted over the count rate (s^{-1}; *vertical axis*) versus the time from the start of the observation (s; *horizontal axis*) plane.

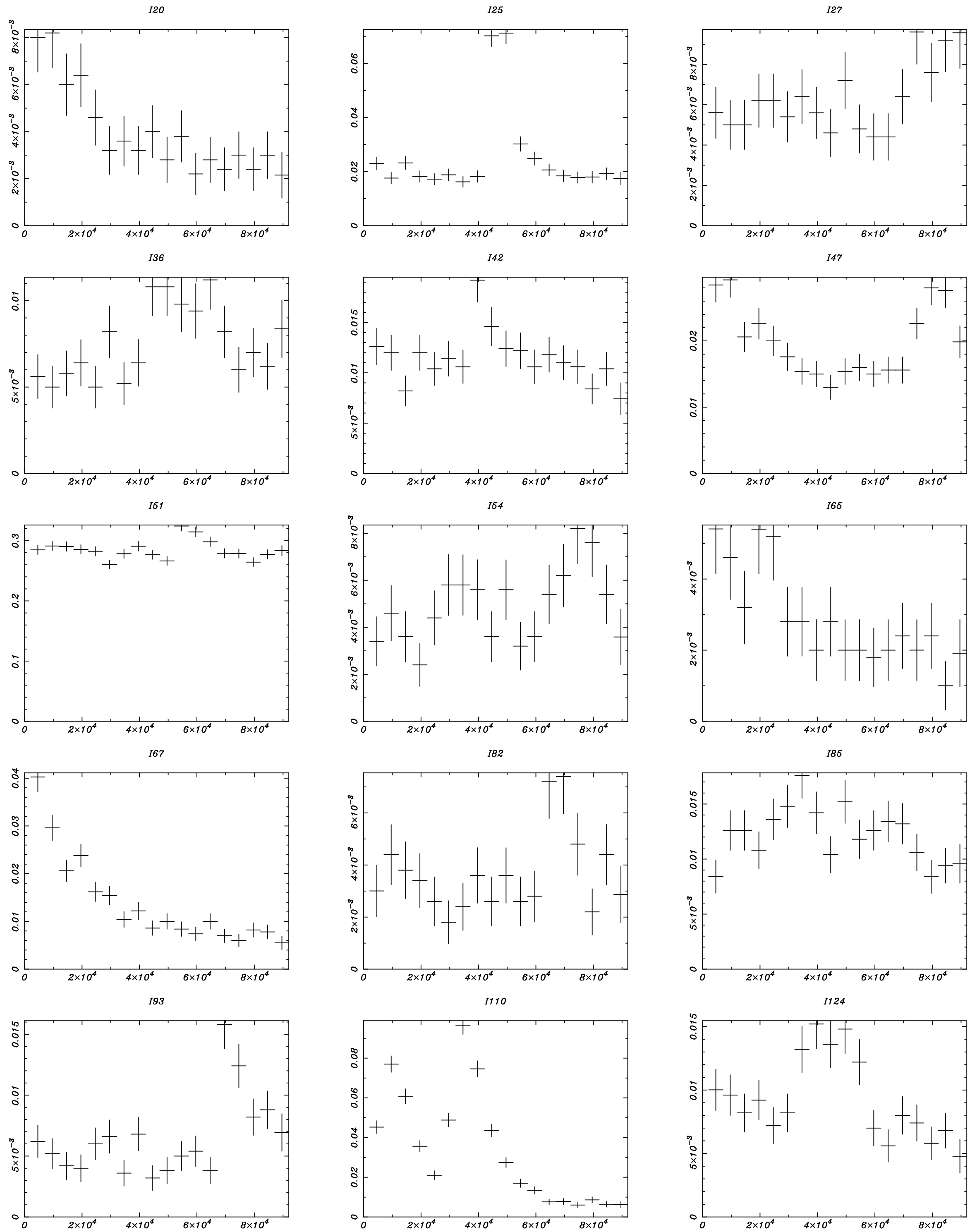


Figure 7.3: Light curves of variable bright (T) NIR-IDed *Chandra* sources (I20–I124).

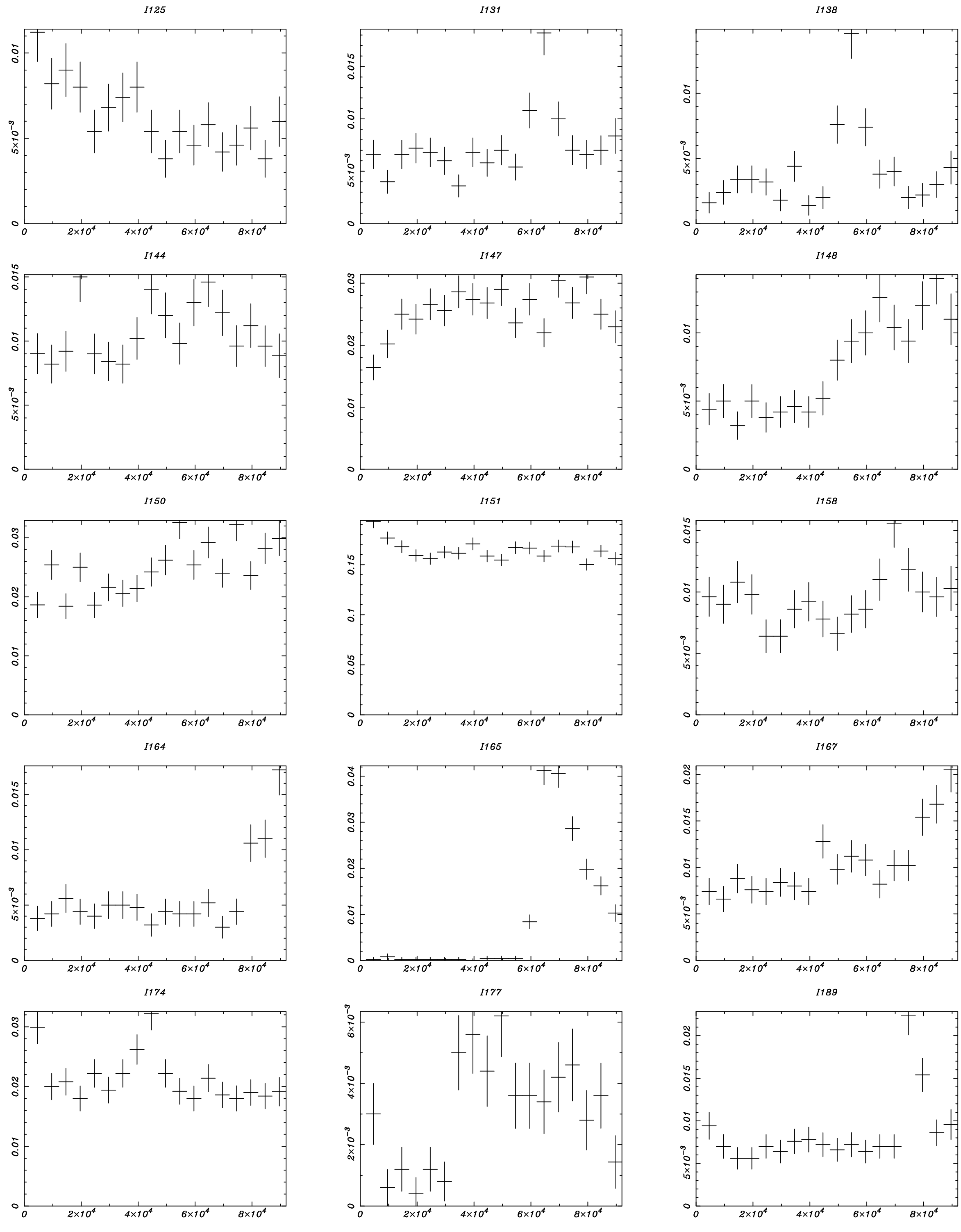


Figure 7.4: Light curves of variable bright (T) NIR-IDed *Chandra* sources (I125–I189).

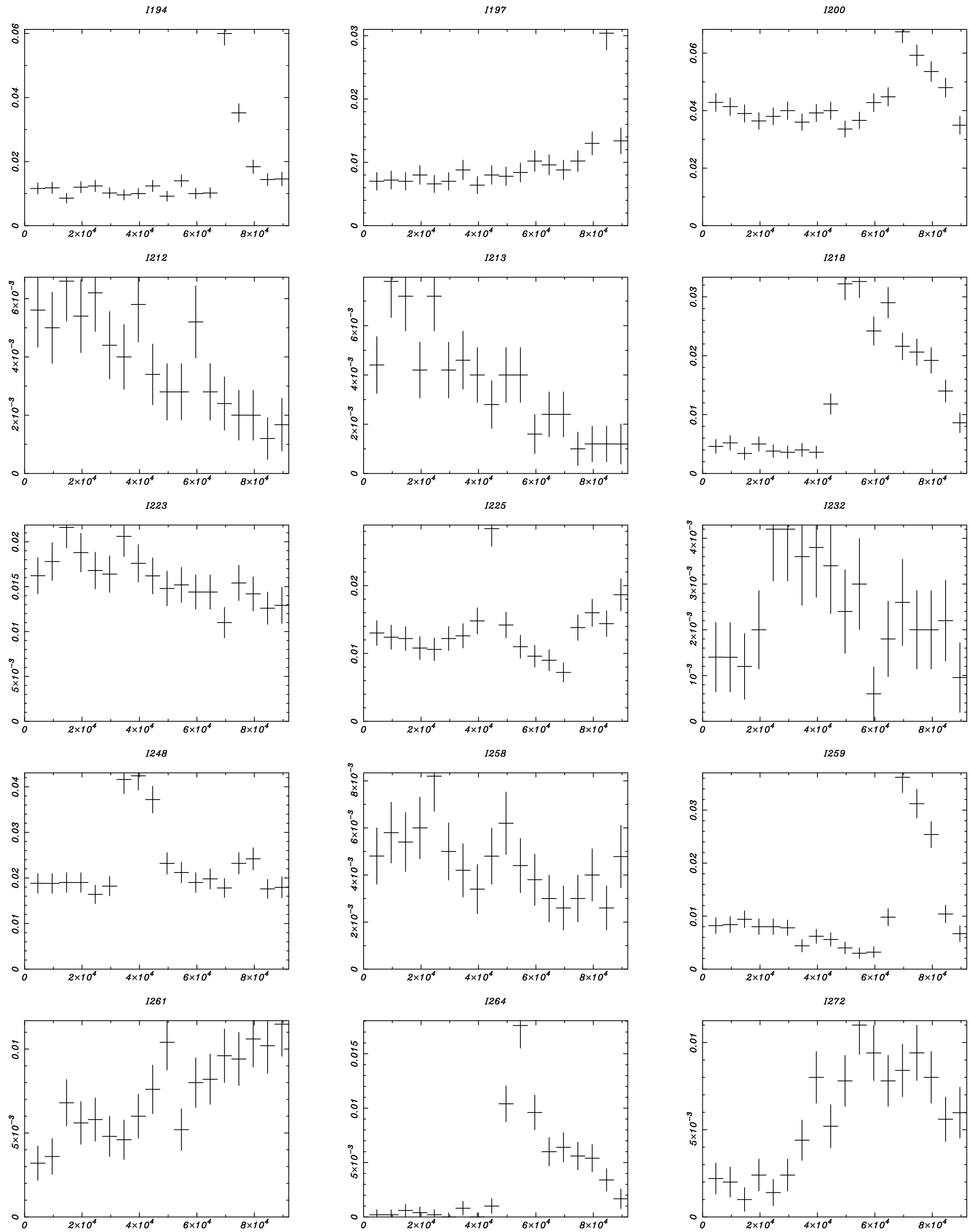


Figure 7.5: Light curves of variable bright (T) NIR-IDed *Chandra* sources (I194–I272).

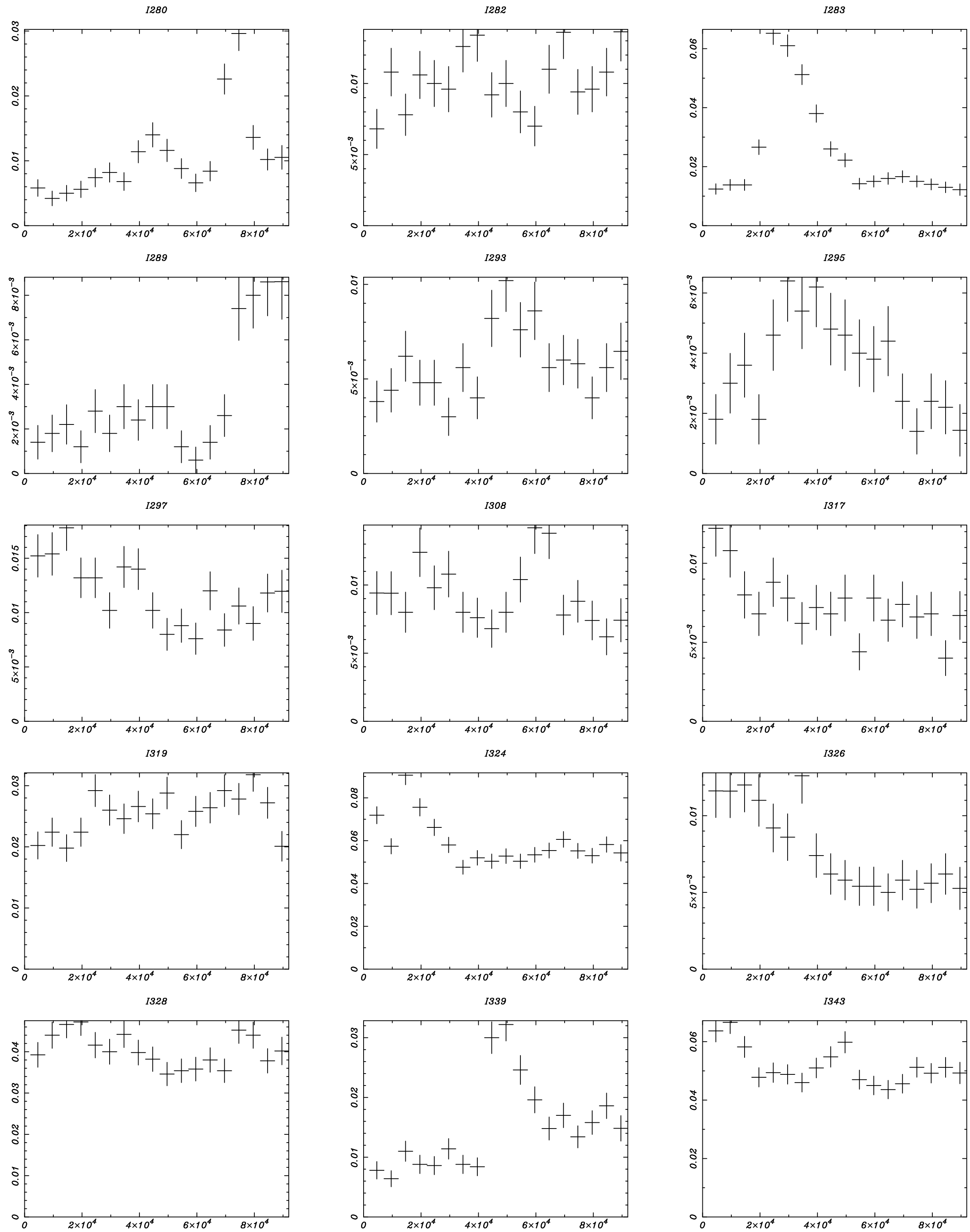


Figure 7.6: Light curves of variable bright (T) NIR-IDed *Chandra* sources (I280–I343).

7.2 Spectral Analysis

7.2.1 Spectrum Models and Fittings

As almost all the NIR-IDed X-ray sources are considered to be YSOs and the X-ray emissions from YSOs are of plasma origin (Feigelson & Montmerle 1999[53]), we fitted the X-ray spectra of NIR-IDed sources with the mekal model convolved with the wabs model.

The mekal model implements the X-ray emissions from thermal plasma at the optically thin limit; i.e., the plasma is transparent to its own X-ray radiation. The plasma consists of electrons and collisionally ionized and excited atoms. When an electron is accelerated by the electric field around an atom, it emits X-rays known as "bremsstrahlung". When the bremsstrahlung is integrated over all electrons at a thermal equilibrium with a Maxwellian velocity distribution, the intensity of the radiation $I(E,T)$ at the energy E and the plasma temperature T is expressed as

$$I(E,T) \propto g(E,T) Z^2 n_{\mathrm{e}} n_{\mathrm{i}} (k_B T)^{-\frac{1}{2}} \exp(-\frac{E}{k_B T}), \tag{7.3}$$

where n_e and n_i are the number density of the electrons and the ions, Z is the charge of the ion, and $g(E,T)$ is the Gaunt factor having a weak dependence on E and T. Line emissions from excited and ionized atoms are added on the continuum radiation. The calculation of thin-thermal plasma spectrum in the mekal model is based on Mewe, Gronenschild, & van den Oord (1985)[127], Mewe, Lemen, & van den Oord (1986)[128], and Kaastra (1992)[102] with Fe L-shell line emissions improved by Liedahl, Osterheld, & Goldstein (1995)[116]. The continuum emissions as well as line emissions from atoms (C, N, O, Ne, Na, Mg, Al, Si, S, Ar, Ca, Fe, and Ni) are included. Free parameters of the model are the electron temperature ($k_{\mathrm{B}}T$ keV), metallicity of elements (Z_{C}, Z_{N}, Z_{O}, and so on), and emission measure ($EM = \int n_{\mathrm{e}} n_{\mathrm{H}} dV$ cm^{-3} where n_{e} and n_{H} are the electron and hydrogen density and V is the emitting volume). Given the values of $k_{\mathrm{B}}T$, EM, and the distance to the source (450 pc), the X-ray luminosity (L_{X} ergs s^{-1}) is determined.

The wabs model implements the photoelectric absorption by cold interstellar medium and is used to derive the amount of interstellar absorption using absorption lines in the X-ray spectra. In most cases, C, N, and O lines in the soft X-ray band dominate fittings. The calculation is based on the Wisconsin cross-sections that include photoelectric absorption cross-sections of major cosmic elements (Morrison & McCammon 1983[136]). Neither

Thomson nor Compton scattering is considered. The spectral intensity of the model (M) at a given energy (E) is described as

$$M(E) = \exp(-p_1 \sigma(E)), \tag{7.4}$$

where $\sigma(E)$ is the effective photoelectric absorption cross-section that is derived by converting all the cross-sections of major elements into the hydrogen equivalent and averaging them out. The amount of absorption (p_1) is thus determined as the equivalent hydrogen column density ($N_{\rm H}$ cm^{-2}).

The cosmic abundances used in the mekal and wabs models are respectively based on the results by Anders & Grevesse (1989)[4] and Anders & Ebihara (1982)[3], who derived values with meteorites and chondrites in the solar system and the solar photosphere and corona.

The combined model of mekal and wabs is further convolved with the optics and detector responses (Auxiliary Response Function [ARF] and Redistribution Matrix Function [RMF], respectively). Both of ARF and RMF for *Chandra*/ACIS are provided by *Chandra* X-ray Center. This convolution is fitted to the spectra that were constructed by binning the X-ray photons along the energy with 20 counts bin^{-1} (for sources of more than 200 counts) or 10 counts bin^{-1} (for sources of less than 200 counts). Again, the χ^2 value was calculated for each fitting trial as

$$\chi^2 = \sum_{i=1}^{N} \left(\frac{E_{data}^{(i)} - E_{model \otimes ARF \otimes RMF}^{(i)}}{\Delta E_{data}^{(i)}} \right)^2, \tag{7.5}$$

where $E_{data}^{(i)}$ and $E_{model \otimes ARF \otimes RMF}^{(i)}$ are the spectrum intensity (s^{-1} keV^{-1}) of background-subtracted data and the convolution of the model, ARF, and RMF in the i'th energy bin. $\Delta E_{data}^{(i)}$ is the uncertainty of $E_{data}^{(i)}$ and is calculated as $\sqrt{E_{data}^{(i)}}$. The upper probability for the null hypothesis (α) for the minimum χ^2 value was used to discriminate whether a fitting is acceptable or not. The XSPEC package[2] was used for all these procedures.

7.2.2 One-temperature Plasma Fittings

We concentrated our spectral analysis on bright ACIS-I sources in the same way as our temporal analysis. Figure 7.7 shows the distribution of X-ray counts (0.5–8.0 keV) and S/N (2.0–8.0 keV), which is defined as the ratio of the source and the background counts in

[2] See http://heasarc.gsfc.nasa.gov/docs/xanadu/xspec/index.html.

the 2.0–8.0 keV range. Unlike Figure 7.1, we used the S/N values in the hard band (2.0–8.0 keV). This is because the spectrum in this energy range is important to see whether a source has hard X-ray component in addition to the soft component, nevertheless the source spectrum can be more easily contaminated by the background spectrum in this band. We defined "bright (S)" sources for those with more counts than 50 and larger S/N than 10 and "faint (S)" for the remaining sources among the 278 NIR-IDed ACIS-I sources. The former consists of 142 (51%) samples while the latter of 136 (49%).

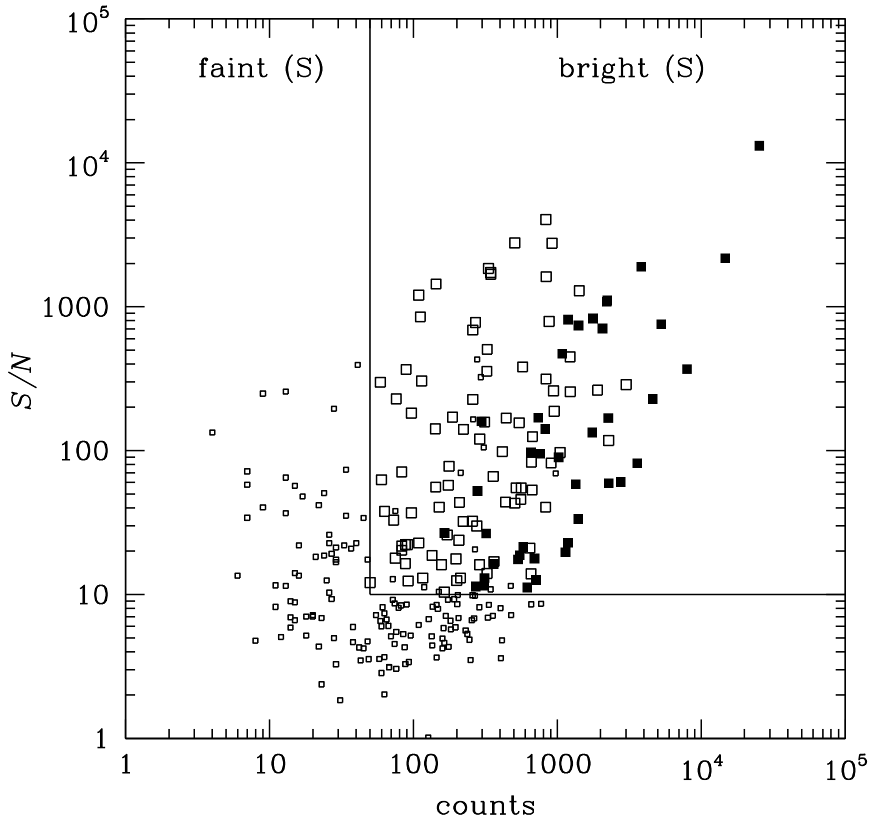


Figure 7.7: X-ray counts (0.5–8.0 keV) and S/N (2.0–8.0 keV) of NIR-IDed ACIS-I sources. Sources with more counts than 50 and higher S/N than 10 (142 sources) are bright (S) sources, for which the spectral analysis was conducted. Larger squares indicate sources fitted by thin-thermal plasma models; open for one-temperature model and filled for two temperature model.

First, for bright (S) sources, we applied one-temperature thin-thermal plasma model (the mekal model) convolved with the ISM absorption (the wabs model). The metallicity of all elements is fixed to 0.3 solar based on previous works (e.g., Imanishi et al. 2001[91]).

When α is less than 0.05 or the best-fit parameters are unphysical (e.g., $k_{\rm B}T > 10$ keV), we recognized that the model was rejected. Among fitted samples, 87 (61%) sources had an acceptable fit. Using the best-fit parameters, we compiled $N_{\rm H}$, $k_{\rm B}T$, and $L_{\rm X}$ in Table 7.1. The spectra and the best-fit models of these sources are tiled in Figures 7.8–7.13.

Table 7.1: One-temperature plasma fittings of bright NIR-IDed *Chandra* sources

ID	counts[a]	S/N[b]	$N_{\rm H}$[c] (10^{22} cm^{-2})	$k_{\rm B}T$[c] (keV)	$L_{\rm X}$[a] (ergs s^{-1})
I21	276	29.9	0.00 (0.00–0.07)	1.18 (1.00–1.31)	3.65e+29
I27	558	45.7	0.03 (0.00–0.15)	1.58 (1.34–1.79)	9.84e+29
I34	135	18.7	0.00 (0.00–0.07)	1.26 (0.99–1.67)	1.70e+29
I36	670	53.2	0.07 (0.01–0.15)	1.09 (1.01–1.22)	1.04e+30
I38	83	21.7	0.72 (0.39–1.24)	0.63 (0.22–0.97)	5.32e+29
I54	437	43.9	0.45 (0.00–0.31)	0.84 (0.98–1.42)	7.04e+29
I58	92	12.4	0.40 (0.00–1.15)	6.39 (2.42–79.9)	3.62e+29
I65	257	32.3	0.77 (0.41–1.25)	2.72 (1.67–5.06)	1.26e+30
I66	200	12.5	0.09 (0.00–0.53)	0.97 (0.60–1.17)	3.23e+29
I67	1232	256.5	0.81 (0.66–0.93)	2.74 (2.35–3.45)	6.36e+30
I74	442	168.5	2.25 (1.84–2.77)	2.21 (1.74–2.82)	4.28e+30
I77	835	314.7	3.20 (2.83–3.63)	3.16 (2.51–4.13)	1.02e+31
I82	325	14.0	2.67 (2.13–3.44)	4.19 (2.41–6.80)	3.47e+30
I83	109	22.8	0.71 (0.46–0.99)	0.46 (0.30–0.67)	9.44e+29
I87	207	23.8	1.93 (1.04–2.47)	2.06 (1.32–3.93)	1.97e+30
I90	116	13.0	0.65 (0.00–1.39)	7.09 (2.79–79.9)	5.64e+29
I92	313	157.9	0.55 (0.36–0.72)	0.60 (0.45–0.83)	1.55e+30
I93	574	382.2	0.00 (0.00–0.07)	1.20 (1.00–1.12)	7.16e+29
I96	333	1845.7	1.93 (1.38–2.53)	2.48 (1.83–3.66)	2.62e+30
I99	112	849.6	2.79 (1.55–4.30)	1.91 (1.19–3.75)	1.25e+30
I106	151	40.4	1.46 (1.15–2.08)	1.42 (1.00–1.76)	1.02e+30
I109	177	77.9	0.65 (0.35–0.93)	0.37 (0.22–0.69)	1.70e+30
I110	3013	287.9	1.23 (1.12–1.35)	7.69 (6.01–10.4)	2.24e+31
I111	88	16.4	0.44 (0.00–1.10)	0.41 (0.11–0.88)	2.04e+29
I121	259	227.2	2.69 (1.77–3.71)	1.70 (1.18–3.00)	3.28e+30
I124	830	40.4	0.72 (0.60–0.87)	4.41 (3.80–6.03)	5.04e+30

(*cont.*)

ID	counts[a]	S/N[b]	N_H[c]	k_BT[c]	L_X[a]
I125	561	54.9	0.33 (0.19–0.51)	3.25 (2.34–4.59)	1.96e+30
I128	60	62.7	12.7 (0.46–49.2)	3.05 (0.72–79.9)	1.55e+31
I131	662	83.3	0.10 (0.00–0.23)	3.19 (2.35–4.48)	1.42e+30
I138	359	66.0	3.22 (2.21–4.28)	7.52 (3.54–79.9)	3.96e+30
I140	259	689.2	6.22 (4.36–12.4)	3.12 (1.32–5.84)	6.32e+30
I142	97	36.9	11.6 (3.75–21.1)	1.07 (0.54–3.56)	4.84e+30
I143	143	55.8	1.60 (0.88–2.41)	2.99 (1.77–6.40)	1.03e+30
I144	953	187.5	1.68 (1.48–1.89)	2.55 (2.13–3.10)	8.08e+30
I145	73	32.9	0.00 (. . . .–. . . .)	1.00 (. . . .–. . . .)	8.76e+28
I147	2276	117.4	1.04 (0.95–1.15)	2.78 (2.45–3.20)	1.46e+31
I148	673	125.0	0.50 (0.32–0.65)	2.76 (2.17–3.93)	2.68e+30
I149	109	1204.9	5.32 (3.23–9.10)	1.48 (0.75–3.02)	5.84e+30
I153	142	141.8	1.38 (1.11–2.21)	0.88 (0.43–1.13)	1.73e+30
I154	97	182.1	1.91 (1.23–3.54)	2.30 (1.18–4.34)	7.64e+29
I155	76	229.1	7.90 (. . . .–. . . .)	1.99 (. . . .–. . . .)	1.27e+30
I158	838	1617.2	0.49 (0.20–0.42)	1.09 (1.19–1.46)	1.96e+30
I164	507	2778.6	2.06 (1.71–2.46)	1.80 (1.48–2.27)	4.52e+30
I165	833	4039.7	1.40 (1.14–1.70)	4.78 (3.34–7.87)	5.48e+30
I166	415	98.2	0.77 (0.65–0.92)	0.30 (0.23–0.48)	7.36e+30
I167	921	2759.8	0.45 (0.31–0.57)	3.12 (2.56–4.13)	3.37e+30
I170	164	10.4	0.50 (0.16–0.99)	7.01 (3.40–36.9)	6.08e+28
I173	59	298.2	1.64 (. . . .–. . . .)	2.92 (. . . .–. . . .)	3.74e+29
I174	1908	263.3	0.00 (0.00–0.02)	1.30 (1.24–1.34)	2.74e+30
I178	83	20.3	0.21 (0.00–1.01)	8.93 (2.16–79.9)	2.82e+29
I179	363	16.9	0.43 (0.28–0.53)	0.61 (0.51–0.74)	1.59e+30
I181	144	1437.5	2.97 (1.39–4.70)	7.54 (2.94–79.9)	1.23e+30
I187	544	156.1	0.00 (0.00–0.06)	1.24 (1.06–1.32)	7.84e+29
I192	324	355.7	0.64 (0.32–0.82)	0.94 (0.78–1.14)	1.34e+30
I197	878	790.1	0.01 (0.00–0.05)	1.38 (1.30–1.60)	1.29e+30
I198	83	71.1	0.18 (0.00–1.02)	0.76 (0.14–1.11)	1.61e+29
I202	88	22.2	2.82 (1.43–4.57)	2.00 (1.07–3.55)	1.04e+30
I204	289	120.3	1.97 (1.50–2.51)	2.62 (1.83–3.40)	2.28e+30
I208	270	775.9	1.46 (1.02–2.30)	4.24 (2.10–21.9)	1.95e+30

(*cont.*)

ID	counts[a]	S/N[b]	$N_{\rm H}$[c]	$k_{\rm B}T$[c]	$L_{\rm X}$[a]
I212	345	1727.6	3.67 (2.60–5.00)	2.35 (1.60–3.89)	4.16e+30
I213	326	504.9	3.16 (2.18–4.27)	2.02 (1.38–3.49)	4.20e+30
I216	172	25.9	0.17 (0.00–0.46)	1.02 (0.88–1.32)	3.06e+29
I217	50	12.1	0.43 (0.00–2.34)	3.15 (0.81–79.9)	1.65e+29
I218	1228	448.2	3.63 (3.24–4.10)	3.10 (2.49–3.95)	1.80e+31
I219	89	366.8	1.31 (0.00–3.11)	1.50 (0.87–7.35)	6.60e+29
I229	157	16.1	0.88 (0.01–0.33)	1.06 (1.70–3.35)	4.28e+28
I223	1424	1291.3	0.07 (0.04–0.11)	1.22 (1.02–1.28)	2.20e+30
I232	213	13.0	1.47 (0.94–2.18)	1.75 (1.14–2.87)	1.52e+30
I234	222	32.2	1.15 (0.85–1.73)	1.49 (1.00–2.84)	1.30e+30
I237	91	22.2	0.46 (0.14–0.99)	0.42 (0.12–0.65)	4.52e+30
I243	187	170.8	0.00 (0.00–0.11)	1.20 (1.04–1.40)	2.23e+29
I246	63	37.8	2.51 (0.46–7.60)	5.68 (1.43–79.9)	6.12e+29
I251	175	57.5	2.08 (1.44–2.84)	3.14 (2.06–5.65)	1.41e+30
I261	646	21.0	0.50 (0.37–0.65)	8.64 (5.85–17.9)	3.08e+30
I263	223	140.6	1.48 (0.50–2.03)	1.26 (0.83–4.09)	1.97e+30
I264	346	1681.7	7.14 (3.82–10.7)	5.48 (3.05–79.9)	5.56e+30
I272	507	43.1	0.00 (0.00–0.04)	1.75 (1.57–2.21)	8.48e+29
I274	209	43.6	0.59 (0.36–0.99)	0.62 (0.80–1.31)	2.78e+29
I279	114	304.4	4.00 (2.16–6.27)	2.52 (1.37–5.85)	1.36e+30
I280	943	260.1	0.33 (0.21–0.43)	2.70 (2.24–3.50)	3.42e+30
I282	908	82.1	0.02 (0.00–0.08)	1.02 (0.95–1.08)	1.24e+30
I293	518	55.1	0.54 (0.30–0.75)	3.01 (2.14–5.36)	2.26e+30
I297	1048	96.8	0.53 (0.43–0.59)	3.45 (2.91–4.36)	4.48e+30
I299	198	17.7	0.03 (0.00–0.14)	1.35 (1.16–1.70)	8.96e+28
I303	75	17.9	0.66 (0.00–1.23)	0.98 (0.46–3.07)	3.28e+30
I317	657	13.9	0.00 (0.00–0.05)	1.30 (1.22–1.38)	1.09e+30
I322	288	16.1	0.34 (0.00–0.16)	0.68 (0.84–1.07)	3.78e+29

[a] Values in the 0.5–8.0 keV range.

[b] Values in the 2.0–8.0 keV range.

[c] The lower and upper limit (1 σ) are given in parentheses. Three sources (I145, I155, and I173) have too few spectral bins to derive the uncertainty of their best-fit parameters.

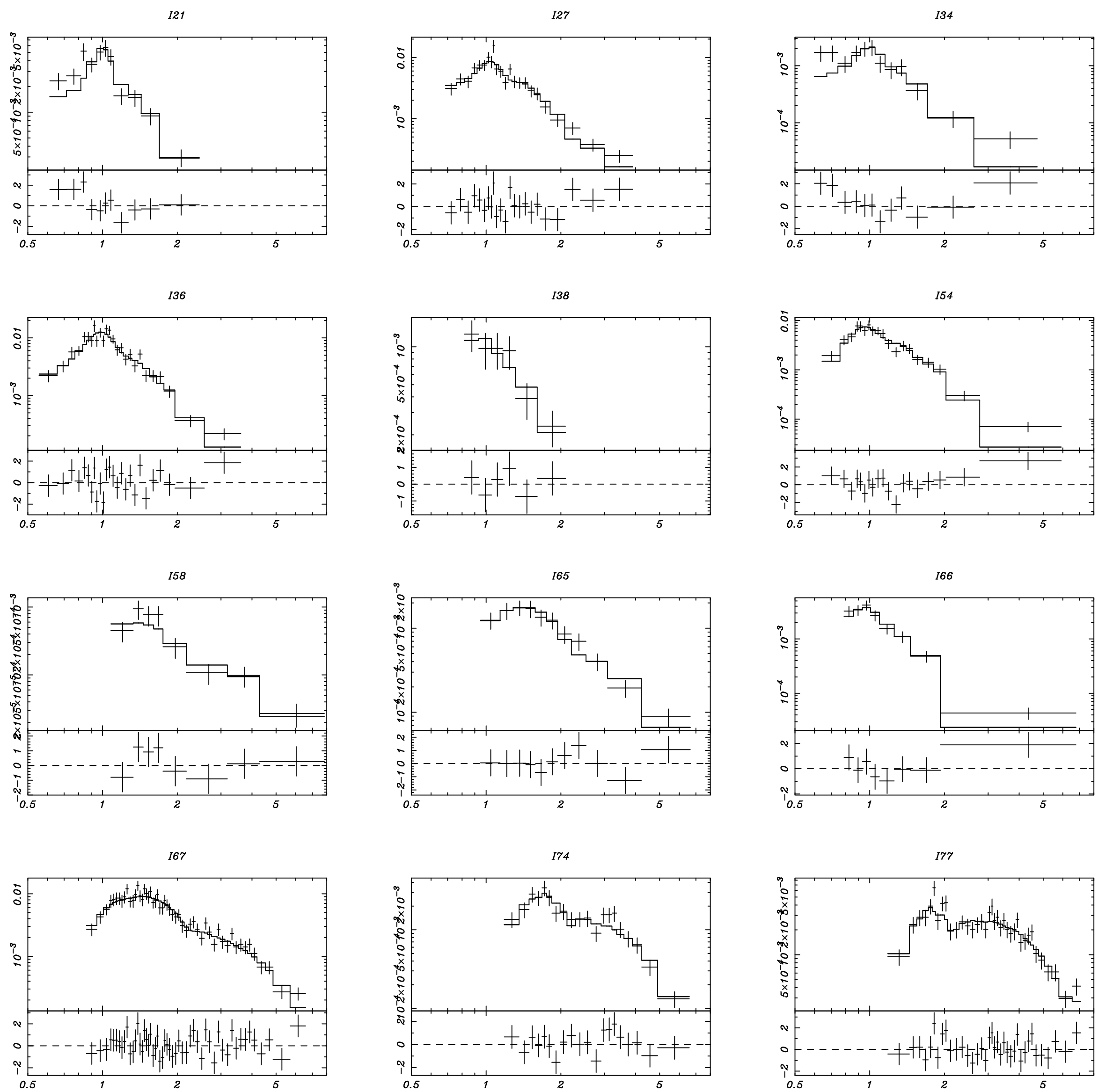


Figure 7.8: Spectra and the best-fit models of one-temperature thin-thermal plasma fittings of bright (S) NIR-IDed *Chandra* sources (I21–I77). The metallicity of all elements is fixed to be 0.3 solar. In the upper panels, the data (*pluses*) and the best-fit model (*solid steps*) are plotted over the energy (keV; *horizontal axis*) versus normalized spectral intensity (count rate keV^{-1}; *vertical axis*) plane. The response functions of the optics and the detector are convolved into the model. In the lower panels, the residuals between the background-subtracted data and the best-fit model are plotted over the energy (keV; *horizontal axis*) versus $\chi^{(i)} = (E^{(i)}_{data} - E^{(i)}_{model \otimes ARF \otimes RMF})/\Delta E^{(i)}_{data}$ (*vertical axis*) plane.

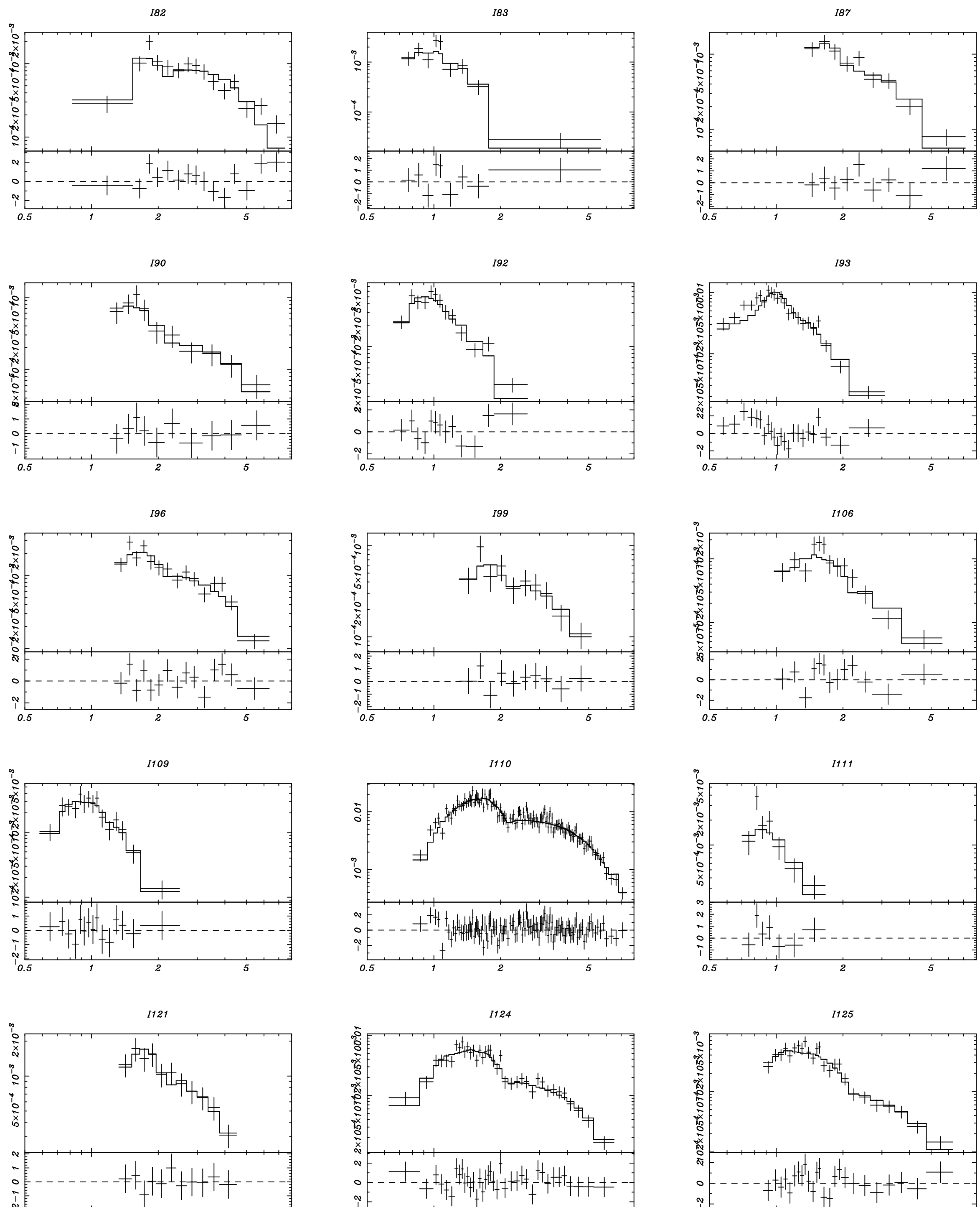


Figure 7.9: Spectra and the best-fit models of one-temperature thin-thermal plasma fittings of bright (S) NIR-IDed *Chandra* sources (I82–I125).

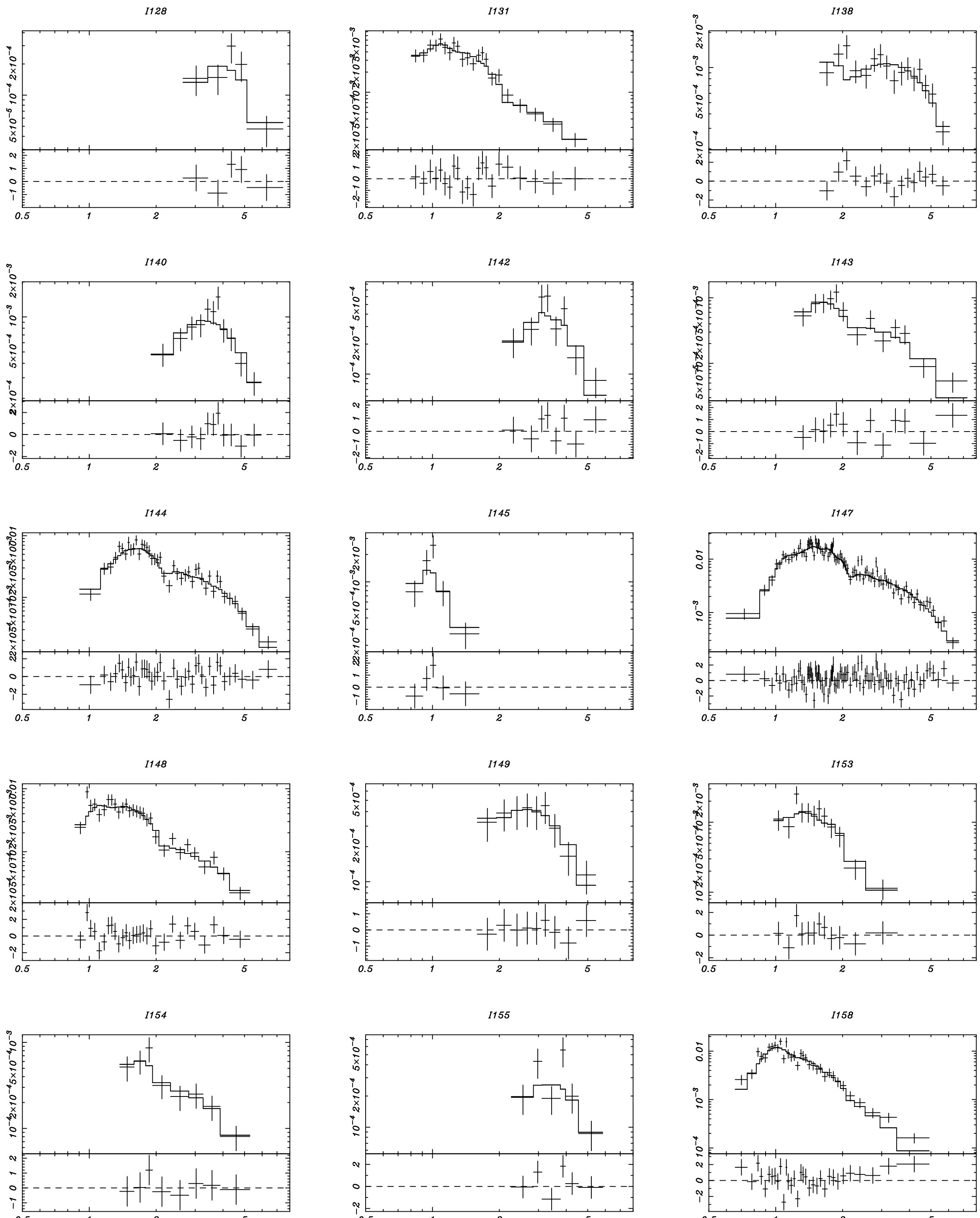


Figure 7.10: Spectra and the best-fit models of one-temperature thin-thermal plasma fittings of bright (S) NIR-IDed *Chandra* sources (I128–I158).

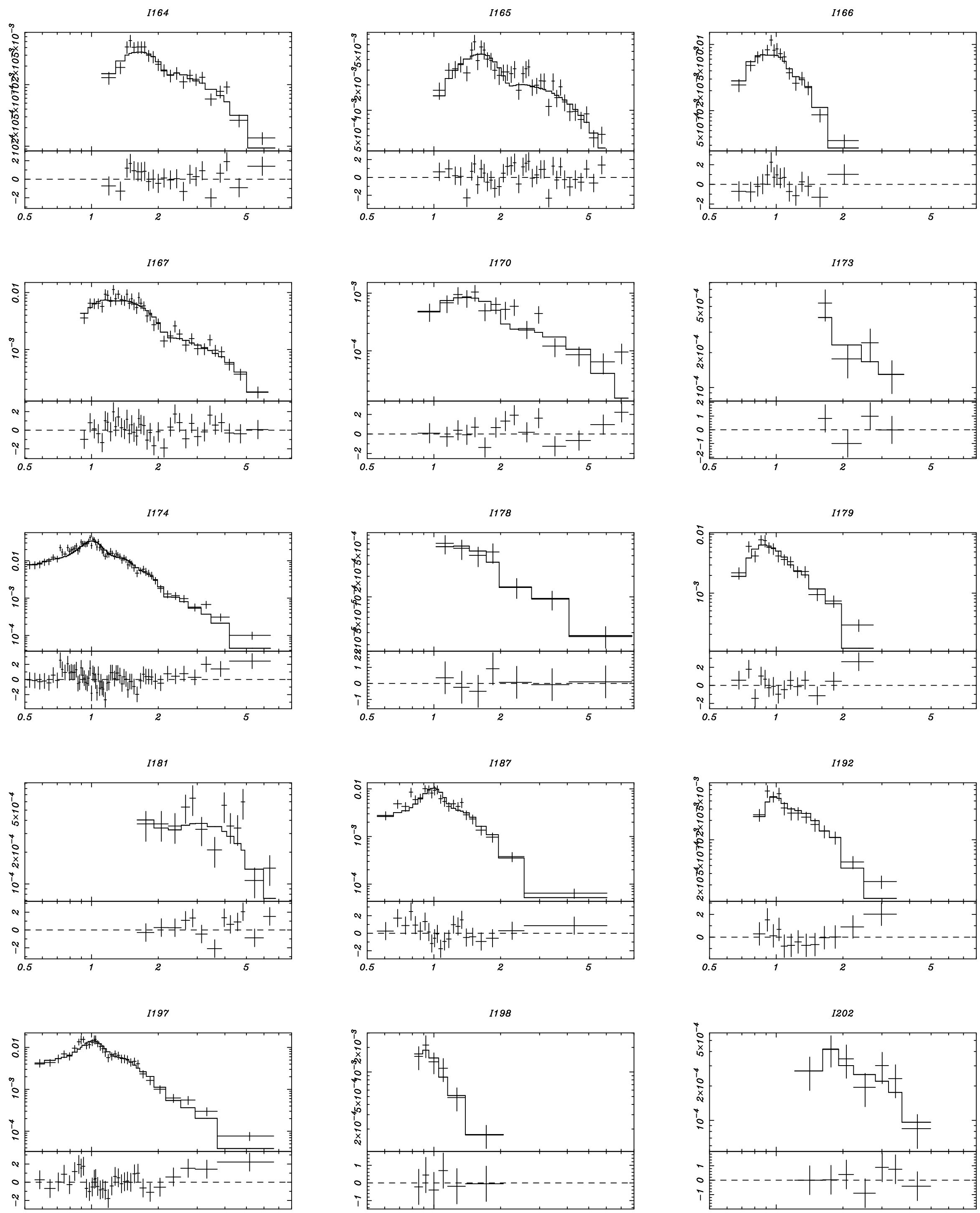


Figure 7.11: Spectra and the best-fit models of one-temperature thin-thermal plasma fittings of bright (S) NIR-IDed *Chandra* sources (I164–I202).

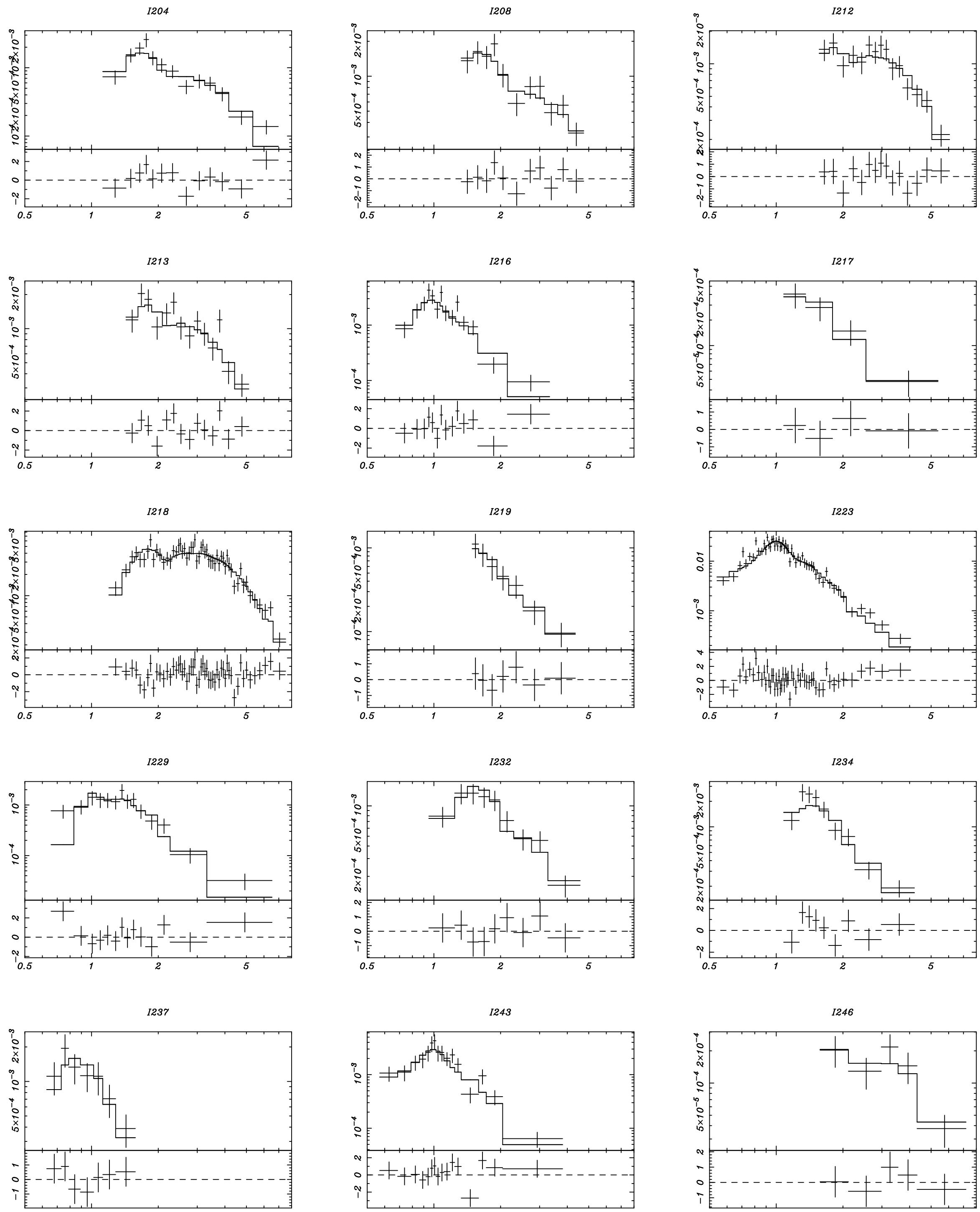


Figure 7.12: Spectra and the best-fit models of one-temperature thin-thermal plasma fittings of bright (S) NIR-IDed *Chandra* sources (I204–I246).

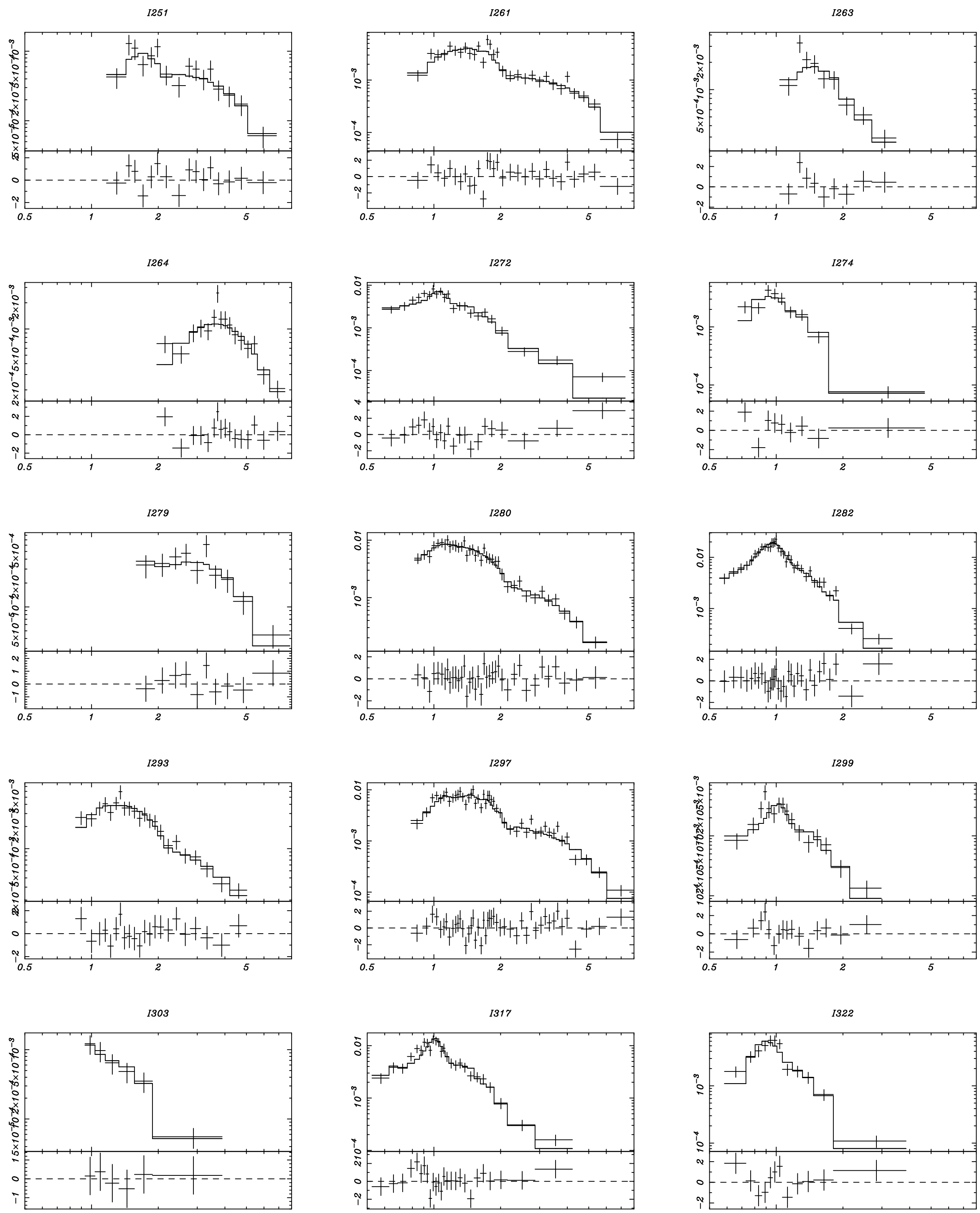


Figure 7.13: Spectra and the best-fit models of one-temperature thin-thermal plasma fittings of bright (S) NIR-IDed *Chandra* sources (I251–I322).

7.2.3 Two-temperature Plasma Fittings

Second, for 55 bright (S) sources that rejected one-temperature plasma model, we applied two-temperature plasma model (two mekal models with different temperatures) convolved with wabs, resulting 32 (23%) sources to have an acceptable fit with $\alpha \geq 0.05$. Both components were assumed to have the same $N_{\rm H}$ value. The metallicity of all elements was again fixed to 0.3 solar. In Table 7.2, we compiled the best-fit values of $k_{\rm B}T$ and $L_{\rm X}$ of the higher temperature component ($k_{\rm B}T_1$ and $L_{\rm X1}$) as well as those of the lower temperature component ($k_{\rm B}T_2$ and $L_{\rm X2}$) along with $N_{\rm H}$ and the total $L_{\rm X}$ of both components. The spectra and the best-fit models are tiled in Figures 7.15–7.17.

We still have 23 bright (S) sources not to be fitted with either one-temperature or two-temperature plasma models with 0.3 solar metallicity. Nine sources, which were globally well fitted with two-temperature plasma model, show excess emission lines that decrease the α value. We therefore fitted these spectra with two-temperature plasma model with variable metallicity for elements with prominent line emissions (two vmekal models), resulting that all of them have an acceptable fit. Table 7.3 lists the best-fit values of continuum emissions ($k_{\rm B}T$ and $L_{\rm X}$ values of the higher and lower temperature component, $N_{\rm H}$ and the total $L_{\rm X}$), while Table 7.4 shows the best-fit values of line emissions (the metallicity of each element). The spectra and the best-fit models are given in Figure 7.18.

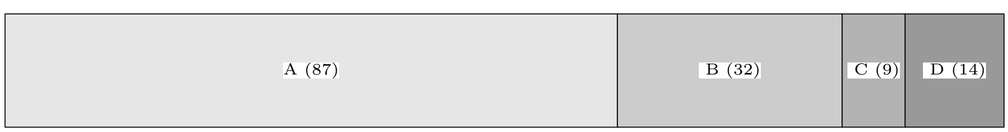


Figure 7.14: Breakdown of best-fit models of 142 bright (S) NIR-IDed ACIS-I sources. A and B are for the spectra fitted with one-temperature and two-temperature plasma model with the metallicity fixed to be 0.3 solar, while C is for the spectra fitted with two-temperature plasma with variable metallicity values. D is for those fitted by none of these models. The number of sources are shown in parentheses.

Table 7.2: Two-temperature plasma fittings of bright NIR-IDed *Chandra* sources (1)

ID	counts[a]	S/N[b]	$N_{\rm H}$[c] (10^{22} cm^{-2})	$k_{\rm B}T_1$[cd] (keV)	$L_{\rm X1}$[ad] (ergs s^{-1})	$k_{\rm B}T_2$[ce] (keV)	$L_{\rm X2}$[ae] (ergs s^{-1})	$L_{\rm X1+2}$[af] (ergs s^{-1})
I7	549	18.8	0.02 (0.00–0.15)	1.68 (1.34–3.26)	4.04e+29	0.74 (0.60–0.86)	2.67e+29	6.71e+29
I12	361	16.3	0.00 (0.00–0.23)	3.15 (0.66–13.8)	5.56e+29	1.04 (0.63–1.37)	1.88e+29	7.44e+29
I14	579	21.3	0.00 (0.00–0.04)	3.49 (2.27–9.44)	7.04e+29	0.88 (0.76–1.06)	3.34e+29	1.04e+30
I17	534	17.5	0.00 (0.00–0.14)	1.94 (1.44–2.77)	5.68e+29	0.77 (0.58–0.87)	2.80e+29	8.48e+29
I25	2287	59.5	0.00 (0.00–0.02)	4.61 (2.48–13.0)	1.61e+30	1.00 (0.87–1.06)	2.24e+30	3.85e+30
I31	740	169.2	0.80 (0.54–1.34)	1.24 (1.01–1.41)	2.51e+30	0.13 (0.09–0.15)	5.68e+31	5.93e+31
I42	1023	89.9	0.90 (0.62–1.07)	1.31 (1.22–1.56)	8.48e+29	0.10 (0.09–0.18)	2.26e+30	3.11e+30
I47	1771	828.8	0.36 (0.27–0.48)	2.45 (2.13–3.07)	3.66e+30	0.84 (0.71–1.03)	1.90e+30	5.56e+30
I71	1191	814.1	0.08 (0.02–0.12)	2.03 (1.15–4.66)	9.44e+29	0.79 (0.67–0.87)	1.06e+30	2.00e+30
I85	1088	471.3	0.02 (0.00–0.07)	4.05 (2.40–6.40)	2.36e+30	1.05 (0.81–1.27)	1.74e+30	4.10e+30
I130	165	26.8	1.00 (0.69–1.15)	2.93 (1.46–50.1)	2.78e+29	0.21 (0.14–1.09)	8.48e+30	8.76e+30
I150	2201	1080.8	0.22 (0.16–0.27)	8.28 (3.86–79.9)	1.26e+30	0.84 (0.78–0.88)	4.00e+30	5.26e+30
I151	14667	2164.2	0.33 (0.31–0.35)	2.80 (2.52–3.34)	2.61e+31	1.10 (1.06–1.21)	1.94e+31	4.55e+31
I189	761	95.0	0.00 (0.00–0.11)	1.47 (1.34–1.75)	8.76e+29	0.55 (0.41–0.70)	2.37e+29	1.11e+30
I194	1411	741.5	0.26 (0.18–0.36)	3.14 (1.73–6.20)	1.80e+30	1.04 (0.79–1.14)	1.76e+30	3.56e+30
I200	3840	1896.2	0.01 (0.00–0.03)	2.45 (2.00–2.75)	2.94e+30	0.87 (0.83–0.91)	3.17e+30	6.11e+30
I211	711	12.6	0.26 (0.17–0.36)	5.67 (3.29–15.8)	2.22e+30	1.21 (1.06–1.34)	1.04e+30	3.26e+30
I225	1189	22.9	0.23 (0.16–0.26)	5.72 (3.76–7.12)	4.36e+30	1.00 (0.84–1.11)	9.32e+29	5.29e+30
I238	311	11.5	0.47 (0.21–1.01)	0.71 (0.45–0.84)	9.88e+29	0.10 (0.08–0.21)	5.88e+30	6.87e+30
I240	2772	60.7	0.43 (0.37–0.50)	2.53 (1.97–3.32)	5.40e+30	1.02 (0.83–1.10)	5.00e+30	1.04e+31
I248	2062	707.9	0.24 (0.19–0.32)	3.01 (2.52–3.61)	3.67e+30	0.85 (0.76–0.97)	1.97e+30	5.64e+30
I256	1143	19.7	0.85 (0.35–1.03)	3.57 (2.80–5.29)	4.04e+30	0.62 (0.48–1.10)	4.60e+30	8.64e+30

(*cont.*)

ID	counts[a]	S/N^b	$N_{\rm H}{}^c$	$k_{\rm B}T_1^{cd}$	$L_{\rm X1}^{ad}$	$k_{\rm B}T_2^{ce}$	$L_{\rm X2}^{ae}$	$L_{\rm X1+2}^{af}$
I278	270	11.4	0.21 (....–....)	6.89 (....–....)	3.67e+29	0.63 (....–....)	3.42e+29	7.09e+29
I287	312	13.0	0.66 (0.35–0.92)	1.11 (0.38–0.63)	4.56e+29	0.21 (0.08–0.30)	1.09e+30	1.55e+30
I289	298	159.8	0.25 (0.00–1.08)	2.21 (1.52–4.26)	7.20e+29	0.18 (0.08–11.2)	6.28e+29	1.35e+30
I298	279	52.4	0.03 (0.00–0.82)	1.38 (1.04–1.80)	3.14e+29	0.35 (0.18–0.68)	1.39e+29	4.53e+29
I308	825	141.5	0.35 (0.23–0.45)	2.80 (1.80–5.16)	2.39e+30	1.06 (0.71–1.41)	1.22e+30	3.61e+30
I313	660	97.8	1.04 (0.74–1.28)	2.82 (2.11–4.45)	2.42e+30	0.85 (0.60–1.06)	2.28e+30	4.70e+30
I314	617	11.2	0.08 (0.03–0.21)	3.79 (1.91–5.78)	8.44e+29	0.95 (0.76–1.05)	5.44e+29	1.39e+30
I319	2262	168.3	0.00 (0.00–0.04)	2.41 (1.90–3.23)	2.99e+30	0.95 (0.80–1.05)	9.96e+29	3.99e+30
I324	5270	757.1	0.09 (0.07–0.11)	2.92 (2.71–3.40)	8.48e+30	0.99 (0.86–1.06)	3.25e+30	1.17e+31
I326	695	17.8	0.21 (0.08–0.41)	2.66 (1.76–10.4)	8.84e+29	0.65 (0.57–0.83)	9.72e+29	1.86e+30

[a] Values in the 0.5–8.0 keV range.

[b] Values in the 2.0–8.0 keV range.

[c] The lower and upper limit (1 σ) are given in parentheses. I278 has too few spectral bins to derive the uncertainty of its best-fit parameters.

[d] Best-fit parameters of the higher temperature component are given.

[e] Best-fit parameters of the lower temperature component are given.

[f] The total luminosity of two components is given.

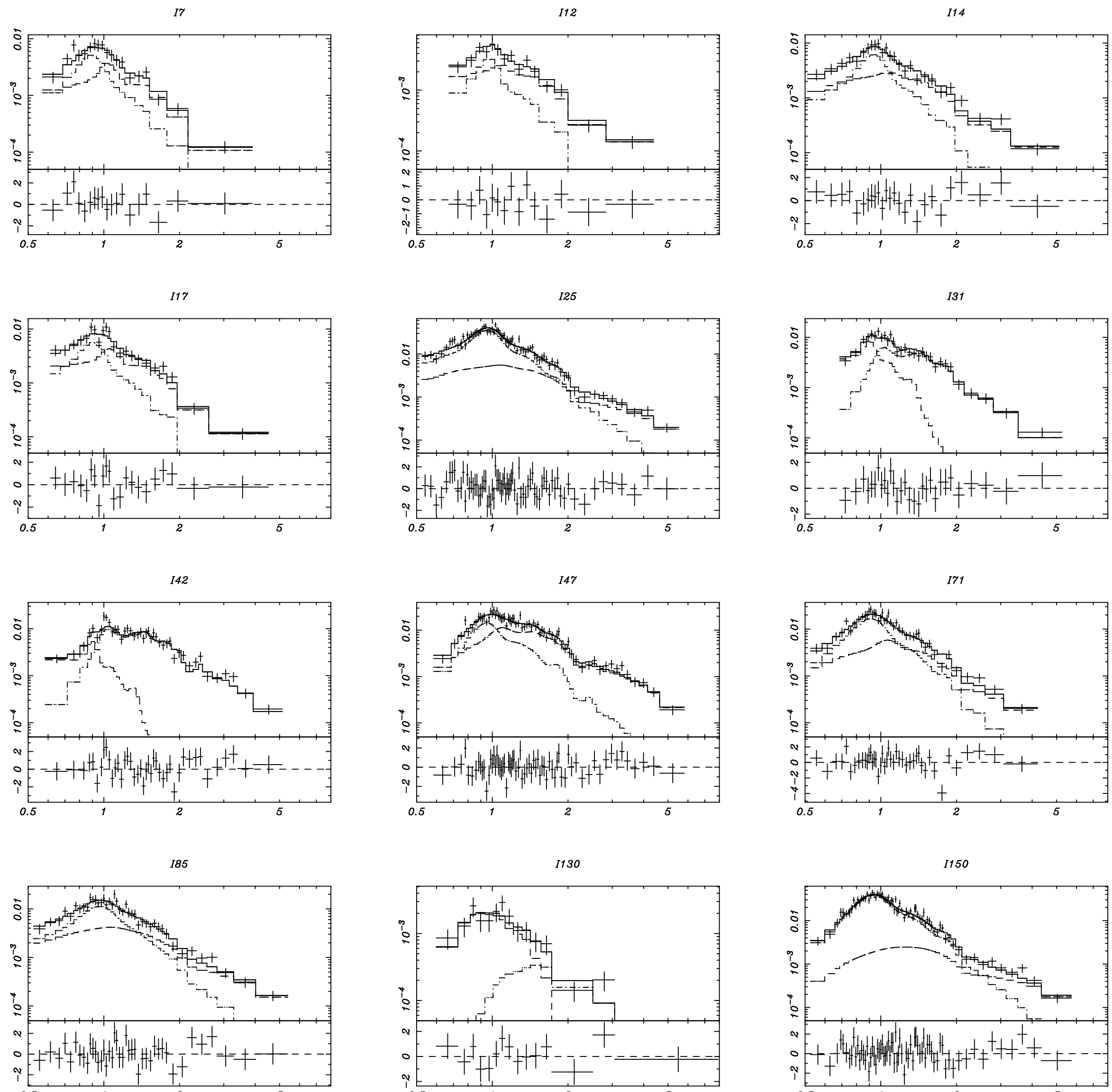


Figure 7.15: Spectra and the best-fit models of two-temperature thin-thermal plasma fittings of bright (S) NIR-IDed *Chandra* sources (I7–I150). The metallicity of all elements is fixed to be 0.3 solar. In the upper panels, the data (*pluses*) and the best-fit model (*solid steps*) are plotted over the energy (keV; *horizontal axis*) versus normalized spectral intensity (count rate keV^{-1}; *vertical axis*) plane. The dashed and dashed-and-dotted steps represent each spectral component. The response functions of the optics and the detector are convolved into the model. In the lower panels, the residuals between the background-subtracted data and the best-fit model are plotted over the energy (keV; *horizontal axis*) versus $\chi^{(i)} = (E^{(i)}_{data} - E^{(i)}_{model \otimes ARF \otimes RMF})/\Delta E^{(i)}_{data}$ (*vertical axis*) plane.

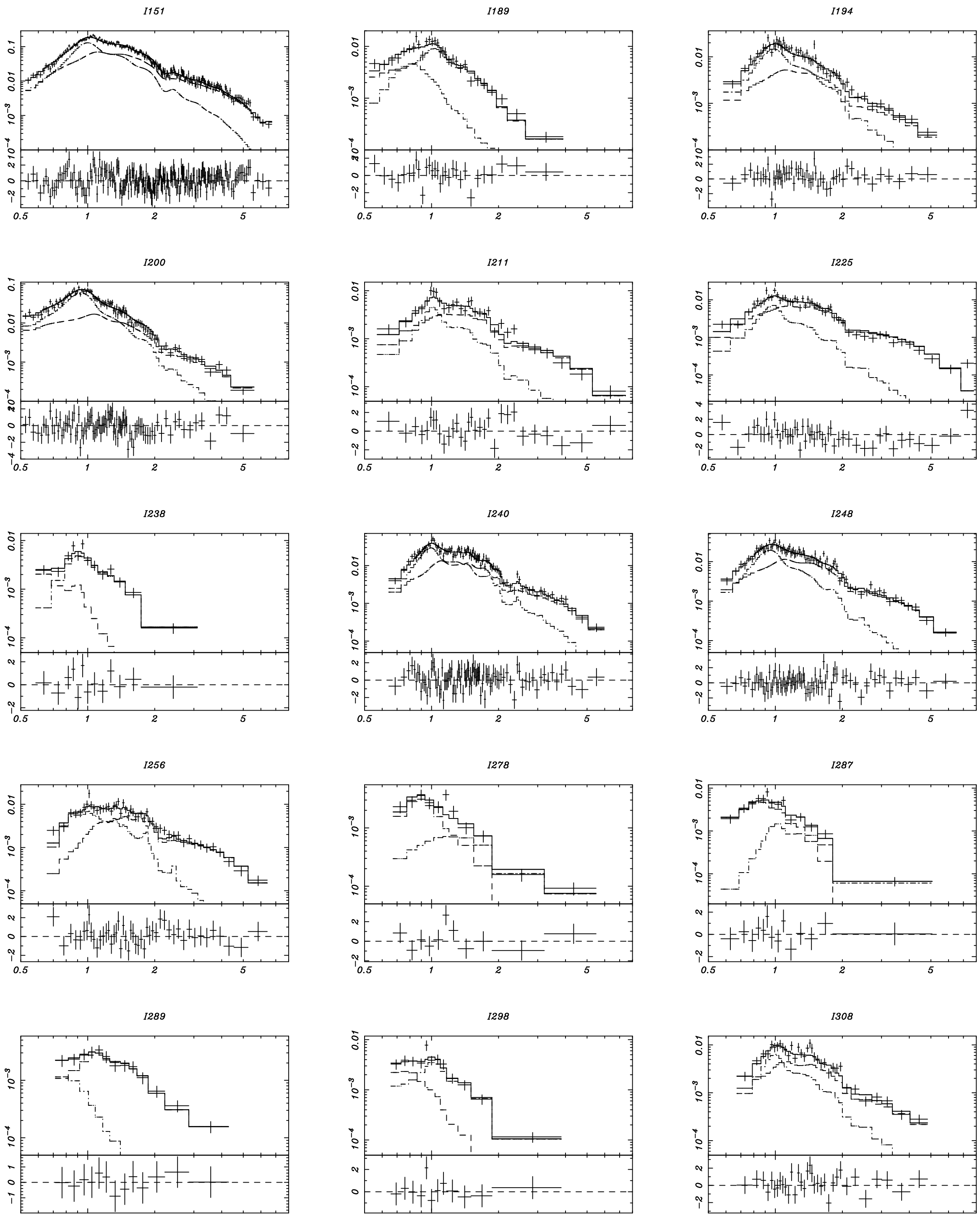


Figure 7.16: Spectra and the best-fit models of two-temperature thin-thermal plasma fittings of bright (S) NIR-IDed *Chandra* sources (I151–I308).

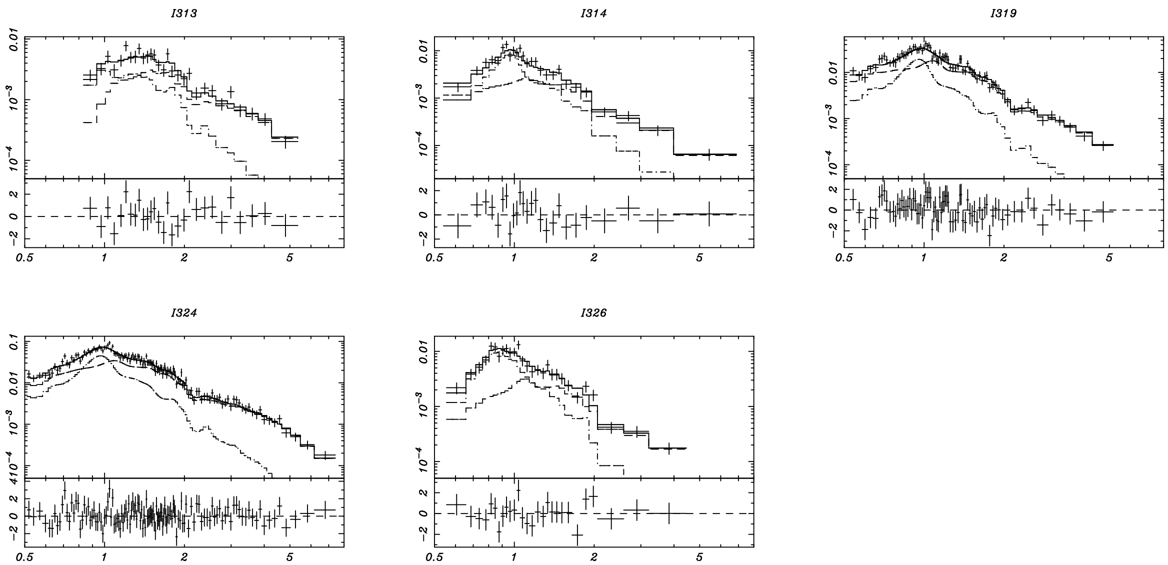


Figure 7.17: Spectra and the best-fit models of two-temperature thin-thermal plasma fittings of bright (S) NIR-IDed *Chandra* sources (I313–I326).

Table 7.3: Two-temperature plasma fittings of bright NIR-IDed *Chandra* sources (2)

ID	counts[a]	S/N[b]	$N_{\rm H}$[c] (10^{22} cm^{-2})	$k_{\rm B}T_1^{cd}$ (keV)	$L_{\rm X}{}_1^{ad}$ (ergs s^{-1})	$k_{\rm B}T_2^{ce}$ (keV)	$L_{\rm X}{}_2^{ae}$ (ergs s^{-1})	$L_{\rm X}{}_{1+2}^{af}$ (ergs s^{-1})
I13	1747	133.5	0.02 (0.00–0.10)	3.29 (2.73–4.06)	2.52e+30	0.70 (0.63–0.79)	8.48e+29	3.37e+30
I51	25400	13166.7	0.20 (0.18–0.22)	2.90 (2.79–3.09)	6.12e+31	0.73 (0.65–0.87)	6.84e+30	6.80e+31
I283	2221	1102.3	0.04 (0.00–0.09)	3.27 (2.55–4.37)	3.83e+30	1.00 (0.81–1.23)	7.04e+29	4.53e+30
I295	320	26.5	0.05 (0.00–0.38)	3.43 (2.48–8.55)	4.92e+29	0.82 (0.60–0.96)	1.43e+29	6.34e+29
I328	3603	82.1	0.07 (0.04–0.12)	1.85 (1.45–1.72)	4.16e+30	0.67 (0.34–0.51)	2.38e+30	6.54e+30
I339	1348	58.4	0.04 (0.00–0.09)	2.38 (1.85–5.58)	1.65e+30	0.95 (0.76–1.11)	8.24e+29	2.47e+30
I342	7968	368.7	0.10 (0.08–0.13)	2.29 (2.12–2.51)	1.44e+31	0.67 (0.61–0.75)	5.12e+30	1.95e+31
I343	4600	228.6	0.21 (0.17–0.27)	1.97 (1.82–2.14)	8.28e+30	0.64 (0.57–0.69)	5.20e+30	1.35e+31
I346	1399	33.4	0.02 (0.00–0.11)	2.11 (1.72–2.55)	1.50e+30	0.73 (0.63–0.85)	8.64e+29	2.36e+30

[a] Values in the 0.5–8.0 keV range.

[b] Values in the 2.0–8.0 keV range.

[c] The lower and upper limit (1 σ) are given in parentheses.

[d] Best-fit parameters of the higher temperature component are given.

[e] Best-fit parameters of the lower temperature component are given.

[f] The total luminosity of two components is given.

Table 7.4: Two-temperature plasma fittings of bright NIR-IDed *Chandra* sources (3)

ID	N[a] (solar)	O[a] (solar)	Ne[a] (solar)	Mg[a] (solar)	Si[a] (solar)	S[a] (solar)	Ar[a] (solar)	Fe[a] (solar)
I13	0.3	0.80 (0.30–1.76)	0.26 (0.00–0.85)	0.3	0.3	0.3	0.3	0.29 (0.15–0.53)
I51	0.3	0.3	1.52 (0.80–2.01)	0.3	0.31 (0.15–0.39)	0.93 (0.69–1.19)	0.78 (0.18–1.38)	0.26 (0.20–0.32)
I283	0.3	0.3	0.3	1.86 (0.41–5.13)	0.3	2.57 (1.17–5.17)	0.3	0.82 (0.30–1.93)
I295	0.3	0.3	0.66 (0.00–)	0.89 (0.00–)	0.3	0.3	0.3	1.05 (0.41–)
I328	0.3	0.21 (0.05–0.28)	1.05 (0.59–1.25)	0.31 (0.10–0.57)	0.3	0.3	0.3	0.13 (0.06–0.16)
I339	0.3	0.3	0.37 (0.00–1.72)	0.3	0.3	0.3	0.3	0.25 (0.18–0.39)
I342	5.15 (0.00–12.7)	0.3	1.05 (0.82–1.31)	0.3	0.3	0.55 (0.29–0.82)	0.40 (0.00–1.14)	0.13 (0.09–0.19)
I343	41.3 (24.3–73.8)	0.3	1.38 (0.88–2.12)	0.3	0.3	0.3	0.3	0.36 (0.26–0.55)
I346	0.3	0.3	1.34 (0.12–2.57)	0.3	0.3	0.3	0.3	0.30 (0.14–0.53)

[a] The lower and upper limit (1 σ) are given in parentheses for free parameters. Others are fixed to be 0.3 solar.

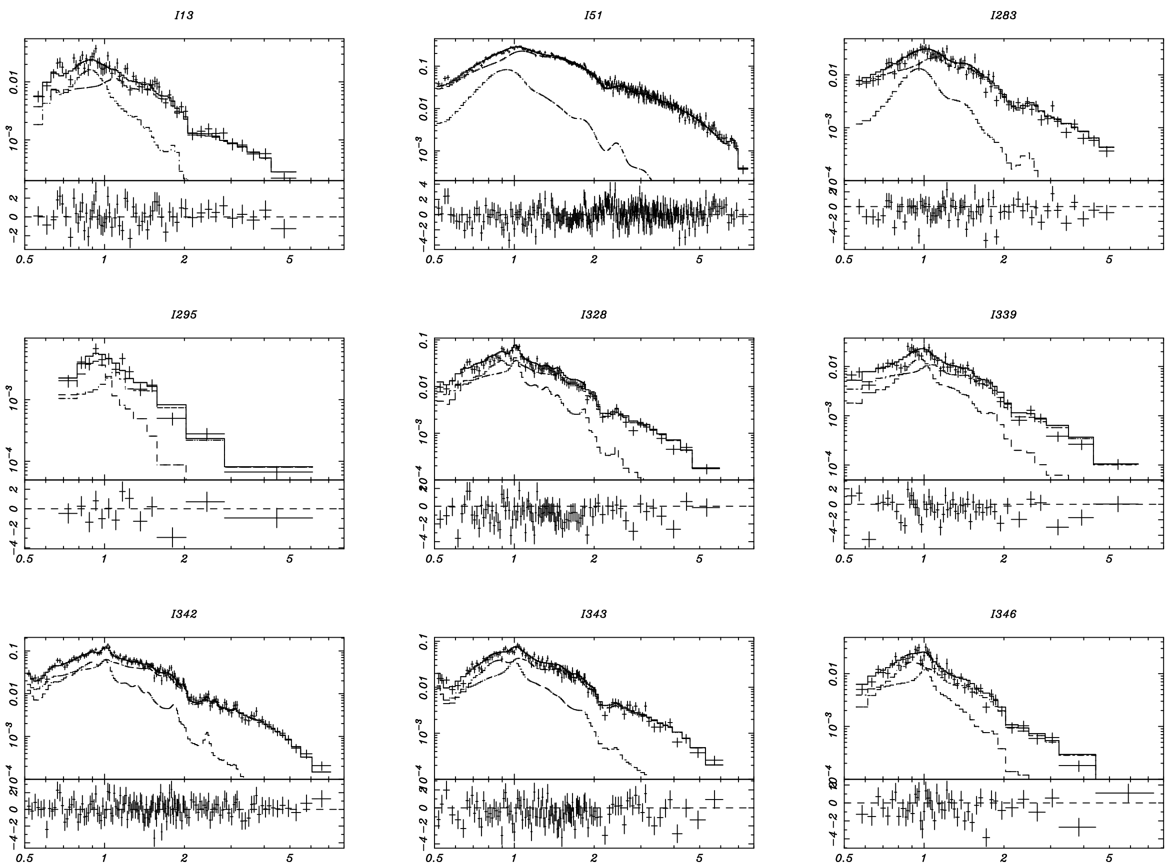


Figure 7.18: Spectra and the best-fit models of two-temperature thin-thermal plasma fittings of bright (S) NIR-IDed *Chandra* sources (I13–I346). The metallicity of all elements is fitted for elements with prominent line emissions. In the upper panels, the data (*pluses*) and the best-fit model (*solid steps*) are plotted over the energy (keV; *horizontal axis*) versus normalized spectral intensity (count rate keV^{-1}; *vertical axis*) plane. The dashed and dashed-and-dotted steps represent each spectral component. The response functions of the optics and the detector are convolved into the model. In the lower panels, the residuals between the background-subtracted data and the best-fit model are plotted over the energy (keV; *horizontal axis*) versus $\chi^{(i)} = (E^{(i)}_{data} - E^{(i)}_{model \otimes ARF \otimes RMF})/\Delta E^{(i)}_{data}$ (*vertical axis*) plane.

7.2.4 Time-sliced Spectroscopy

Finally, we conducted the time-sliced spectroscopy of the bright variable sources. We confined this analysis to those with flare-like variability and more counts than 500 in each slice. Six sources (I25, I51, I110, I200, I248, and I324) meet these criteria. We sliced their light curves (Figs. 7.19 and 7.20) into flare and quiescent phases, and fitted the spectrum of each slice with thin-thermal plasma models. Both slices of I110 were well fitted with one-temperature plasma model. Other slices were not fitted with one-temperature plasma with $\alpha < 0.05$ or a systematic residual, so two-temperature plasma model was used to fit these spectra. The best-fit values are summarized in Table 7.5.

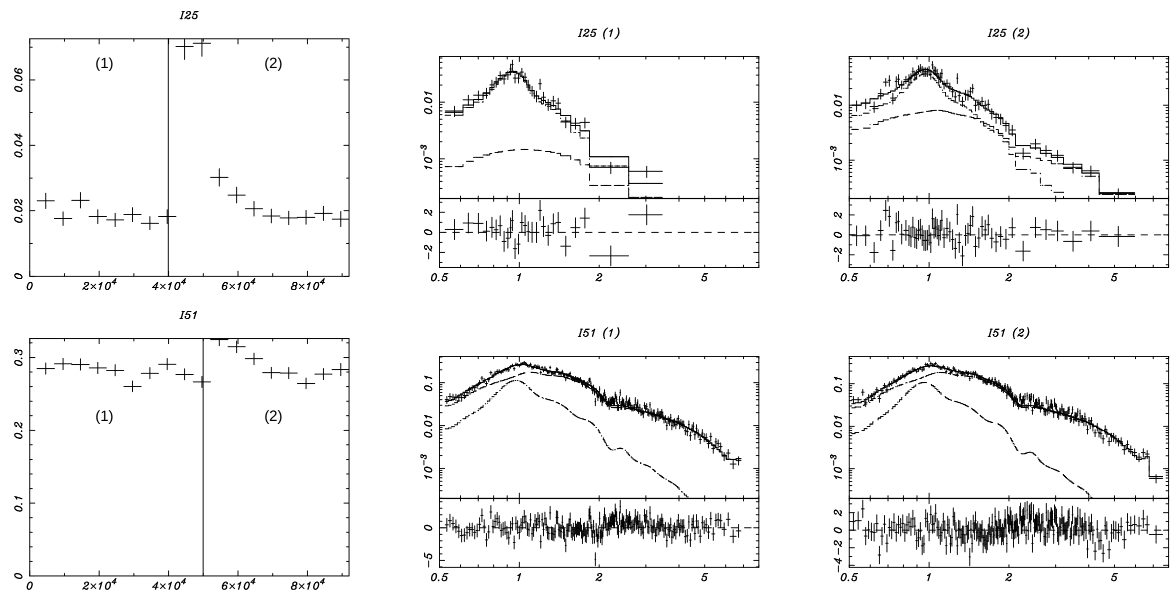


Figure 7.19: Time-sliced spectroscopy of some bright variable NIR-IDed *Chandra* sources (I25–I51). The light curves and the time slices are in the panels of the left column. The time-sliced spectra and the best-fit models of thin-thermal plasma fittings are in the middle and the right column. The metallicity of all elements is fixed to be 0.3 solar. In the upper panels, the data (*pluses*) and the best-fit model (*solid steps*) are plotted over the energy (keV; *horizontal axis*) versus normalized spectral intensity (count rate keV^{-1}; *vertical axis*) plane. The dashed and dashed-and-dotted steps represent each spectral component. The response functions of the optics and the detector are convolved into the model. In the lower panels, the residuals between the background-subtracted data and the best-fit model are plotted over the energy (keV; *horizontal axis*) versus $\chi^{(i)} = (E^{(i)}_{data} - E^{(i)}_{model \otimes ARF \otimes RMF})/\Delta E^{(i)}_{data}$ (*vertical axis*) plane.

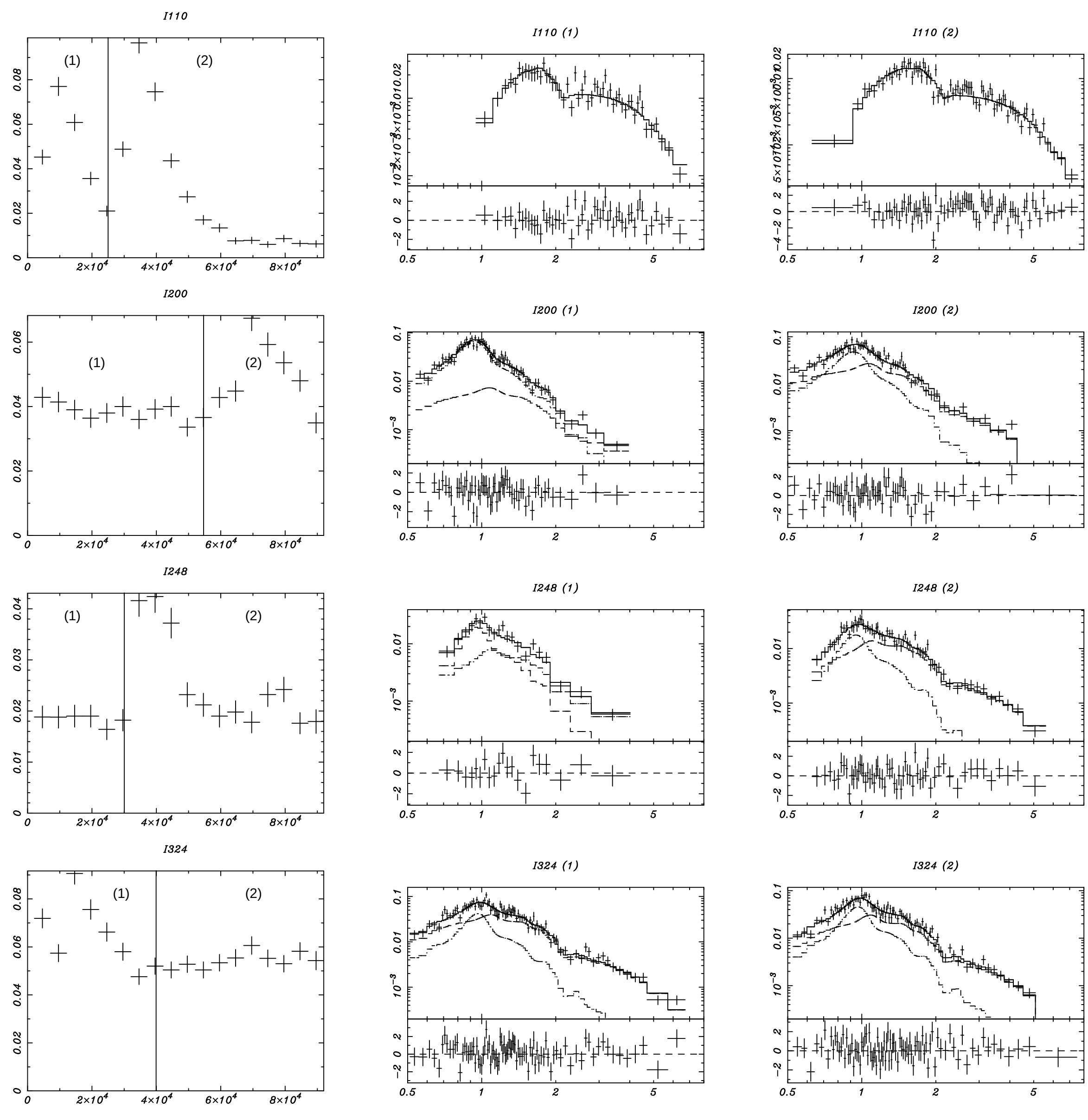


Figure 7.20: Time-sliced spectroscopy of some bright variable NIR-IDed *Chandra* sources (I110–I324).

Table 7.5: Time-sliced spectroscopy of bright NIR-IDed *Chandra* sources (1)

ID	phase[a]	counts[b]	$N_{\rm H}{}^c$ (10^{22} cm^{-2})	$k_{\rm B}T_1^{cd}$ (keV)	EM_1^d (cm^{-3})	$k_{\rm B}T_2^{ce}$ (keV)	EM_2^e (cm^{-3})
I25	1 (Q)	715	0.00 (0.00–0.06)	9.71 (0.21–79.9)	3.71 (1.04–9.14) e+52	0.92 (0.18–0.98)	2.15 (0.07–2.50) e+53
	2 (F)	1572	0.01 (0.00–0.04)	4.83 (2.61–8.77)	1.79 (1.12–2.50) e+53	1.04 (0.91–1.11)	2.78 (2.01–3.49) e+53
I51	1 (Q)	13470	0.19 (0.17–0.20)	3.28 (3.01–3.60)	4.66 (4.28–24.6) e+54	1.01 (0.91–1.10)	1.23 (0.92–1.52) e+54
	2 (F)	11930	0.20 (0.18–0.23)	3.23 (2.95–3.53)	5.04 (4.62–25.0) e+54	0.94 (0.84–1.01)	1.14 (0.88–1.46) e+54
I110	1 (F)	1154	1.52 (1.29–1.77)	6.82 (4.59–12.0)	2.55 (2.20–3.07) e+54	....	....
	2 (F)	1859	1.03 (0.91–1.15)	7.89 (5.85–11.6)	1.11 (1.02–1.23) e+54	....	....
I200	1 (Q)	2037	0.04 (0.00–0.08)	2.65 (1.73–9.42)	1.35 (0.70–2.03) e+53	0.87 (0.83–0.93)	4.19 (3.45–4.98) e+53
	2 (F)	1803	0.00 (0.00–0.03)	2.27 (1.86–2.83)	4.32 (3.59–5.29) e+53	0.84 (0.75–0.90)	2.73 (2.05–3.32) e+53
I248	1 (Q)	511	0.20 (0.04–0.78)	2.52 (0.50–27.4)	1.95 (0.77–4.08) e+53	0.92 (0.73–1.12)	1.91 (0.89–7.43) e+53
	2 (F)	1551	0.22 (0.14–0.32)	3.07 (2.53–3.79)	3.82 (3.08–4.42) e+53	0.85 (0.74–1.01)	1.84 (1.17–2.90) e+53
I324	1 (F)	2495	0.07 (0.05–0.10)	3.17 (2.61–3.99)	8.73 (7.21–10.2) e+53	1.00 (0.84–1.11)	3.27 (2.20–4.59) e+53
	2 (Q)	2775	0.11 (0.08–0.16)	2.63 (1.96–3.26)	6.70 (5.31–8.60) e+53	0.98 (0.80–1.08)	3.80 (2.63–5.06) e+53

[a] Q for the quiescent phases and F for the flare phases.

[b] Values in the 0.5–8.0 keV range.

[c] The lower and upper limit (1 σ) are given in parentheses.

[d] Best-fit parameters of the higher temperature component are given.

[e] Best-fit parameters of the lower temperature component are given.

7.3 Relations between Parameters

7.3.1 X-ray absorption versus NIR extinction

We examined the relation between A_V and $N_{\rm H}$. The former reflects the amount of ISM in the solid-state, while the latter reflects that in any states. The A_V values were determined from the NIR color-magnitude diagram (Fig. 6.7) and the $N_{\rm H}$ values were from the X-ray spectral fittings (Tables 7.1, 7.2, and 7.3). We have 118 sources with A_V and $N_{\rm H}$ values, for which we fitted with a linear relation in Figure 7.21 to obtain

$$\left(\frac{N_{\rm H}}{10^{21}\ {\rm cm}^{-2}}\right) = 1.60^{+0.18}_{-0.15} \times \left(\frac{A_V}{\rm mag}\right). \tag{7.6}$$

The slope of $1.60^{+0.18}_{-0.15}$ is smaller than that of the Galactic interstellar medium (1.79; Predehl & Schmitt 1995[154]) and is comparable to the values obtained in the ρ Ophiuchi dark cloud (1.59$\pm$0.40; Imanishi et al. 2001[91]).

7.3.2 X-ray counts versus X-ray flux

The relation between X-ray counts and the flux ($F_{\rm X}$) in the 0.5–8.0 keV range was examined using bright (S) samples. The linear relation of

$$\log\left(\frac{F_{\rm X}}{{\rm ergs\ s^{-1}\ cm^{-2}}}\right) = \log\left({\rm counts}\right) - 16.1 \tag{7.7}$$

(Figure 7.22) was used to estimate the detection limit of the *Chandra* observation (Sect. 5.1) and the expected number of background sources (Sect. 8.1).

7.3.3 X-ray counts versus X-ray luminosity

To estimate the $L_{\rm X}$ values of faint (S) sources, we examined the relation between the X-ray counts and luminosity ($L_{\rm X}$) in the 0.5–8.0 keV range of bright (S) sources (Fig. 7.23). The linear relation was derived to be

$$\log\left(\frac{L_{\rm X}}{{\rm ergs\ s^{-1}}}\right) = \log\left({\rm counts}\right) + 27.6. \tag{7.8}$$

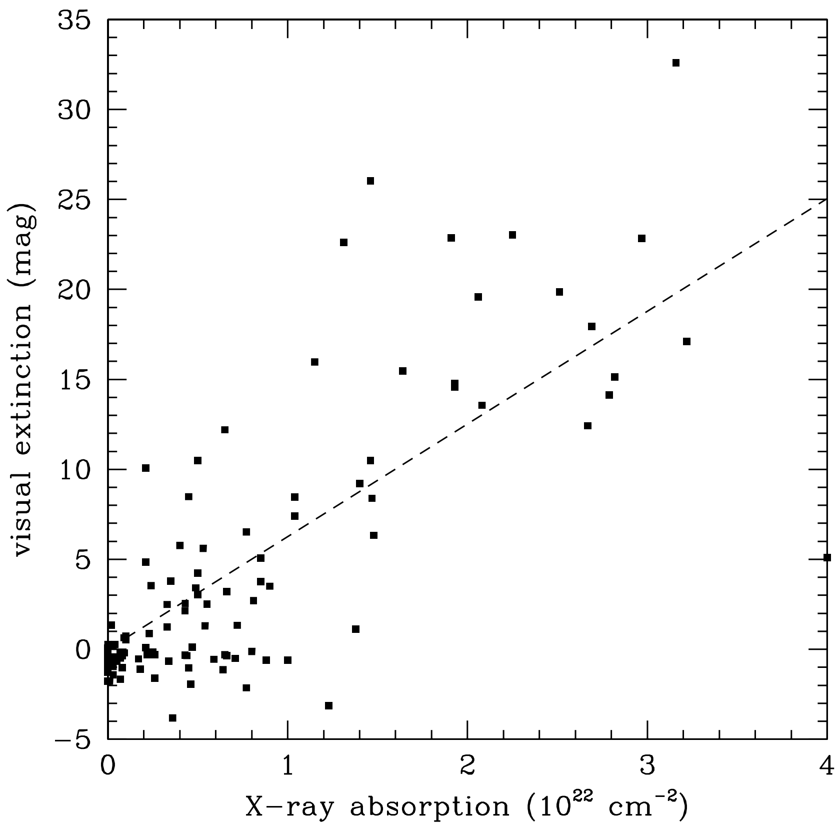


Figure 7.21: Relation between X-ray absorption ($N_{\rm H}$) and visual extinction (A_V) for NIR-IDed *Chandra* sources. The $N_{\rm H}$ values are derived from fitting the X-ray spectra, while A_V values are from the color-magnitude diagram (Fig. 6.7). The best-fit linear relation is shown with the dashed line.

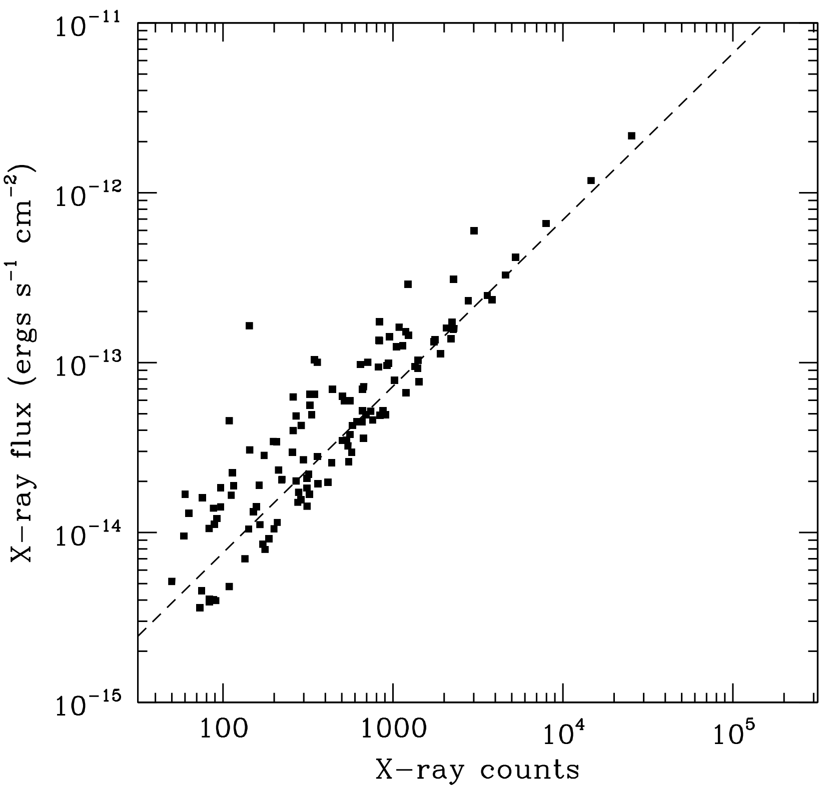


Figure 7.22: Relation between X-ray counts and X-ray flux (F_X) for NIR-IDed bright (S) sources. The best-fit linear relation is shown with the dashed line.

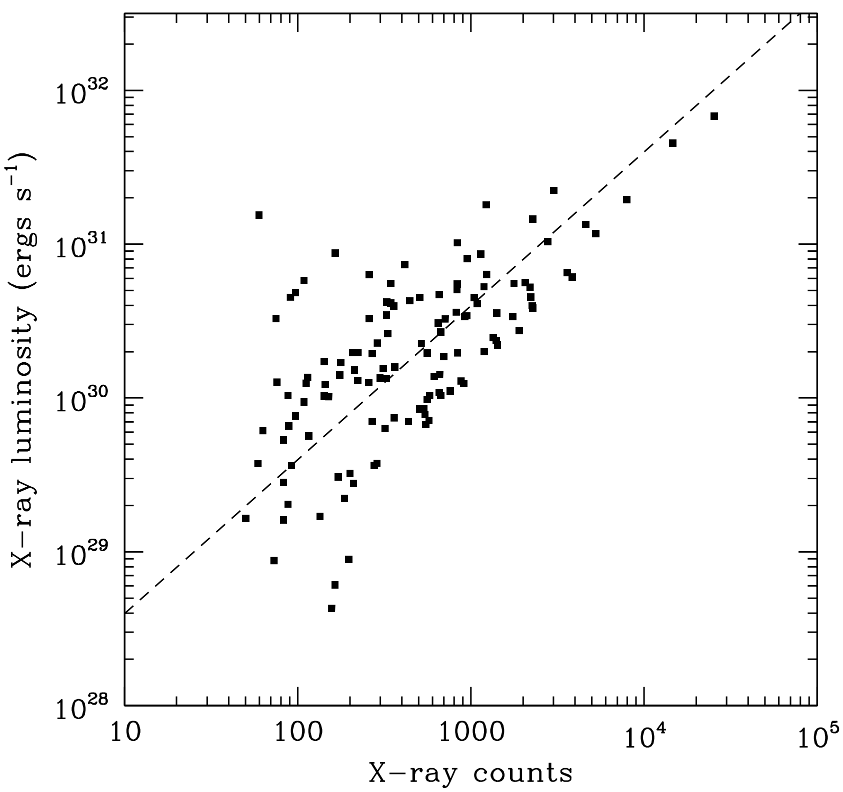


Figure 7.23: Relation between X-ray counts and X-ray luminosity (L_X) for NIR-IDed bright (S) sources. The best-fit linear relation is shown with the dashed line.

Chapter 8

NIR-unIDed X-ray Sources

Contents

In this chapter, we discuss the nature of NIR-unIDed *Chandra* sources. In Sect. 8.1, we conclude that most of the NIR-unIDed sources are background AGNs based on the spectral analysis of some bright sources, their HR, and spatial distribution. There remain a dozen of exceptions, however, that have YSO-like features among NIR-unIDed sources. Some have a softer HR than AGNs and an association with cloud cores. These X-ray emissions are compared with the H_2 image of OMC-3 that we obtained with QUIRC on UH88 (Sect. 8.2). We found that some NIR-unIDed sources are located at 1.3 mm cloud cores and are associated with H_2 outflows. As a test case, we studied the X-ray emissions at MMS 2 and MMS 3 with follow-up observations in the NIR (Sect. 8.3) and the centimeter (Sect. 8.4) band respectively using Subaru and VLA. The nature of NIR and centimeter sources is discussed in each section, which gives vital clues to discuss the origin of these X-ray emissions in Chap. 9.

8.1 The Nature of NIR-unIDed X-ray Sources

8.1.1 Background AGNs

We have 107 *Chandra* ACIS-I sources that were identified with neither 2MASS nor QUIRC sources. Most of these sources are too faint to conduct temporal and spectral analyses unlike NIR-IDed sources, which makes it difficult to discuss their nature on a source basis. However, on the following arguments, we consider that most, if not all, of these NIR-unIDed sources are background AGNs.

First, we fitted the spectra of NIR-unIDed bright (S) sources with a power-law model;

$$M(E) \propto E^{-\Gamma}, \tag{8.1}$$

convolved with the ISM absorption. NIR-unIDed bright (S) sources were defined in Figure 8.1 with the same criteria for NIR-IDed bright (S) sources (Fig. 7.7). Fifteen (14%) sources were found to be bright (S), while 92 (86%) were faint (S).

All of the bright (S) sources were well fitted with a power-law model, with the upper probability of $\alpha > 0.05$. Photons in the 1.0–8.0 keV range were used. Using the best-fit parameters, we derived $N_{\rm H}$, the photon index (Γ), and $F_{\rm X}$ in Table 8.1. The spectra and the best-fit models are shown in Figures 8.2 and 8.3.

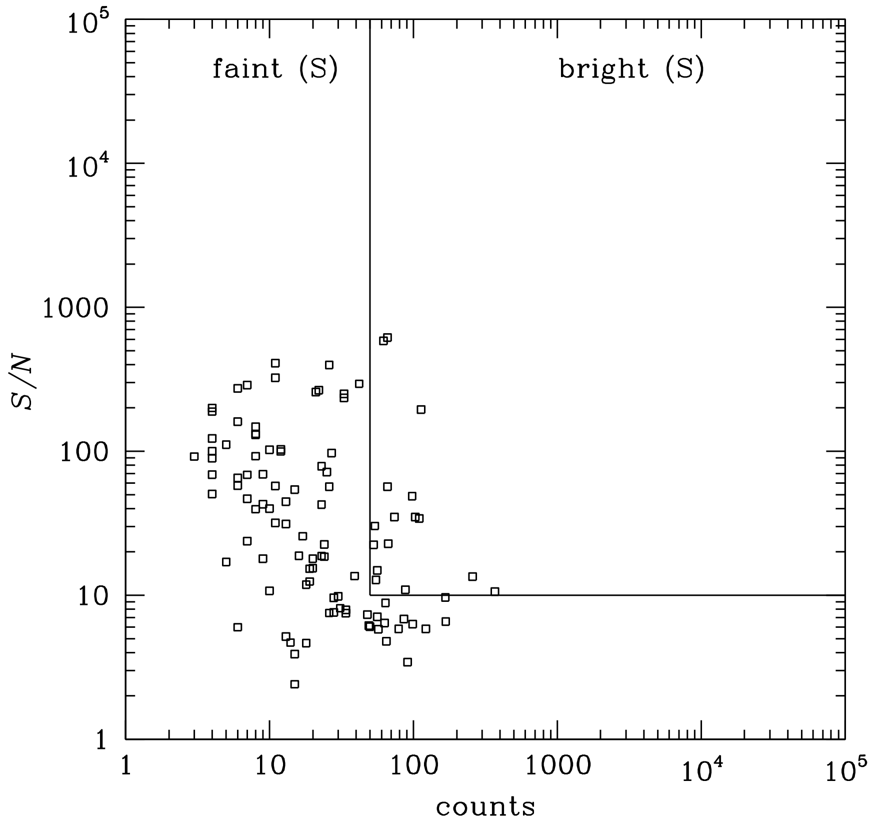


Figure 8.1: X-ray counts (0.5–8.0 keV) and S/N (2.0–8.0 keV) of NIR-unIDed ACIS-I sources. Sources with more counts than 50 and higher S/N than 10 (15 sources) are bright (S) sources, for which the spectral analysis was conducted.

The AGNs generally show the power-law index of $\Gamma \sim 1.7$ in the X-ray band regardless of their types and luminosities (Charles & Seward 1995[33]). Most spectra are consistent with that of a typical AGN. Moreover, when fitted with a thin-thermal plasma model, many spectra had the unphysical best-fit values of $k_B T > 15$ keV. These indicate that the X-ray emissions from these bright (S) sources are from mostly AGNs, not from YSOs.

Table 8.1: Power-law fittings of NIR-unIDed *Chandra* sources

ID^a	counts^b	S/N^c	N_H^{d} (10^{22} cm^{-2})	Γ^d	F_X^{b} (ergs s^{-1} cm^{-2})
I3†	368	10.6	1.43 (0.74–2.00)	0.51 (0.17–0.98)	9.74e−14
I11†	110	34.1	1.40 (0.05–3.48)	1.10 (0.14–2.11)	2.90e−14
I19†	103	35.0	1.21 (0.00–3.20)	1.58 (0.52–2.94)	1.87e−14
I22†	55	12.8	0.91 (0.00–5.27)	1.09 (−0.16–4.01)	1.13e−14
I23†	98	48.7	1.74 (0.00–4.63)	1.35 (0.04–2.78)	2.18e−14
I50	54	30.3	2.27 (0.00–7.29)	2.11 (0.07–6.18)	9.25e−15
I98	62	586.2	2.08 (....–....)	2.01 (.... – ....)	1.46e−14
I186	66	615.8	1.34 (0.00–45.8)	0.99 (−0.54–10.0)	2.64e−14
I266	53	22.4	2.91 (....–....)	0.92 (.... – ....)	1.98e−14
I275	66	56.7	16.2 (....–....)	1.30 (.... – ....)	2.92e−14
I301	113	194.8	3.68 (1.32–6.57)	2.40 (1.65–3.60)	2.41e−14
I325†	74	35.0	1.16 (0.00–5.61)	0.87 (−0.33–3.22)	1.97e−14
I330	67	22.8	7.15 (2.98–13.1)	4.36 (2.30–7.84)	1.03e−14
I336†	88	10.9	0.84 (0.00–2.85)	0.98 (−0.11–1.65)	1.59e−14
I345†	258	13.5	0.47 (0.00–1.10)	0.66 (0.09–0.88)	5.96e−14

[a] Sources that are rejected to have a thin-thermal plasma spectrum are marked with †.

[b] Values in the 0.5–8.0 keV range.

[c] Values in the 2.0–8.0 keV range.

[d] The lower and upper limit (1 σ) are given in parentheses. Three sources (I98, I266, and I275) have too few spectral bins to derive the uncertainty of their best-fit parameters.

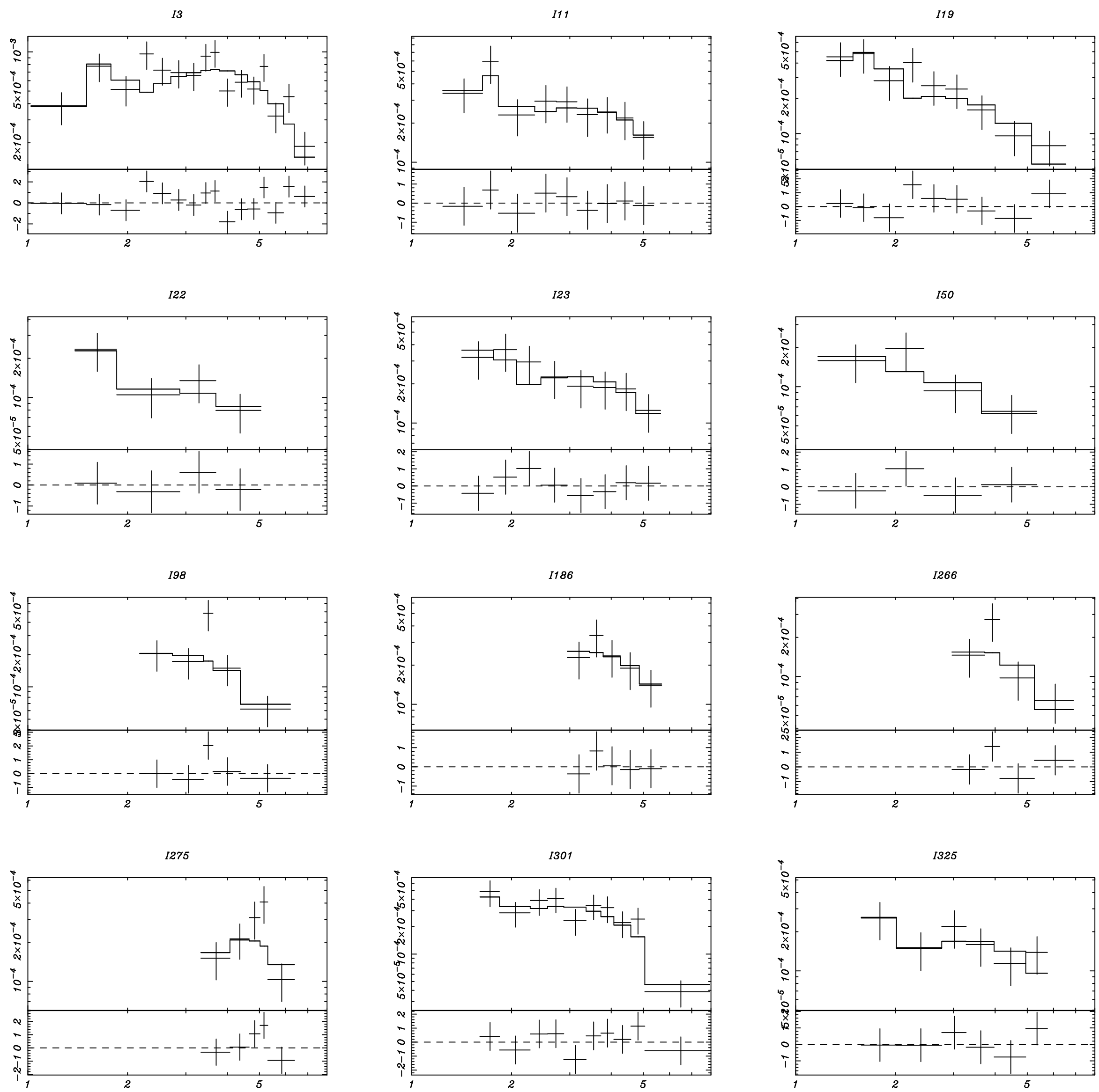


Figure 8.2: Spectra and the best-fit power-law models of NIR-unIDed bright (S) *Chandra* sources (I3–I325). In the upper panels, the data (*pluses*) and the best-fit model (*solid steps*) are plotted over the energy (keV; *horizontal axis*) versus normalized spectral intensity (count rate keV^{-1}; *vertical axis*) plane. The response functions of the optics and the detector are convolved into the model. In the lower panels, the residuals between the background-subtracted data and the best-fit model are plotted over the energy (keV; *horizontal axis*) versus $\chi^{(i)} = (E^{(i)}_{data} - E^{(i)}_{model \otimes ARF \otimes RMF})/\Delta E^{(i)}_{data}$ (*vertical axis*) plane.

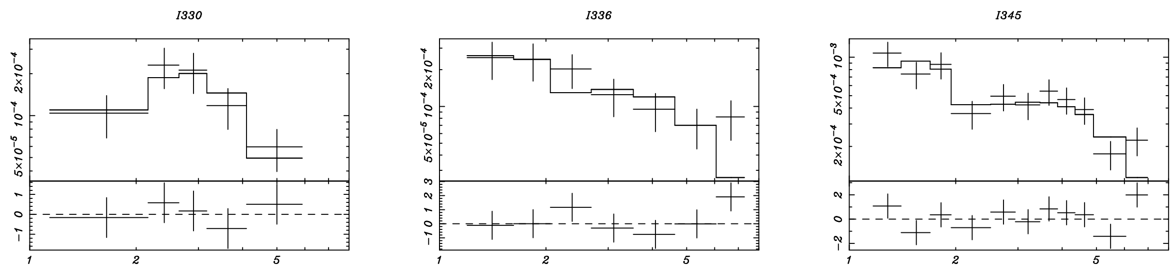


Figure 8.3: Spectra and the best-fit power-law models of NIR-unIDed bright (S) *Chandra* sources (I330–I345).

The second argument is the HR of NIR-unIDed sources. Figure 8.4 shows the histogram of the HR for NIR-IDed (*solid*) and NIR-unIDed (*dashed*) *Chandra* sources. These two histograms are in clear contrast with each other, where NIR-IDed source histogram peaks at $-0.8 < \mathrm{HR} < -0.6$ and NIR-unIDed source histogram at $0.4 < \mathrm{HR} < 0.6$. With completely different profiles, it is natural to consider that NIR-IDed and NIR-unIDed sources have different origins for their X-ray emissions.

The spectra of bright (S) sources show the N_{H} values of $\sim$1–3$\times10^{22}$ cm^{-2} (Table 8.1). When absorbed with this column density, a power-law spectrum of $\Gamma = 1.7$ has the HR of 0.40–0.81. This corresponds to the peak of the NIR-unIDed histogram. On the other hand, a thin-thermal plasma spectrum with $k_{\mathrm{B}}T = 0.8$ keV and $N_{\mathrm{H}} = 6.6\times10^{21}$ cm^{-2} (representative values for the lower-temperature component; see Sect. 9.2) has HR = –0.77, and that with $k_{\mathrm{B}}T = 3.0$ keV and $N_{\mathrm{H}} = 6.6\times10^{21}$ cm^{-2} (representative values for the higher-temperature component) has HR = 0.02. These account for the NIR-IDed histogram. All these infer that the histogram of the NIR-unIDed sources is mainly composed of AGNs, while that of NIR-IDed sources is of YSOs.

To confirm this further, we combined the spectra of the NIR-unIDed sources in the same HR bin in Figure 8.4 from HR = 0.0–0.2 to 0.8–1.0, and fitted them with a power-law model. The spectra and the best-fit models are shown in Figure 8.5, while the best-fit values are in Table 8.2. The indices of the power-law are consistent with $\Gamma = 1.7$ with increasing absorption from softer to harder HRs.

Third, the number of NIR-unIDed X-ray sources is roughly in the same order of the expected number of background AGNs. Giacconi et al. (2001)[68] derived the $\log N - \log S$

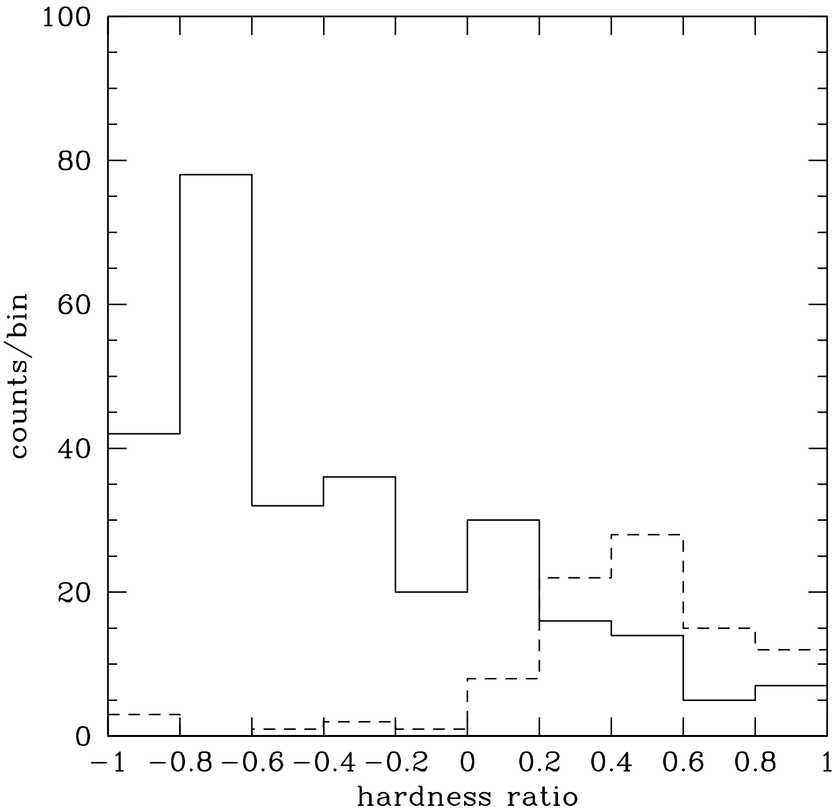


Figure 8.4: Histogram of *Chandra* hardness ratio separately for NIR-IDed (*solid*) and NIR-unIDed (*dashed*) sources.

Table 8.2: Power-law fittings of combined spectra of NIR-unIDed *Chandra* sources

HR range	N_{H} [a] (10^{22} cm^{-2})	Γ [a]
0.0–0.2	0.34 (0.14–0.75)	1.20 (0.81–1.81)
0.2–0.4	1.19 (0.80–1.72)	1.27 (1.00–1.72)
0.4–0.6	2.16 (1.55–3.08)	1.49 (1.12–1.81)
0.6–0.8	3.54 (2.07–5.33)	1.58 (1.11–2.55)
0.8–1.0	5.17 (0.00–6.57)	1.70 (0.15–2.45)

[a] The lower and upper limit (1 σ) are given in parentheses.

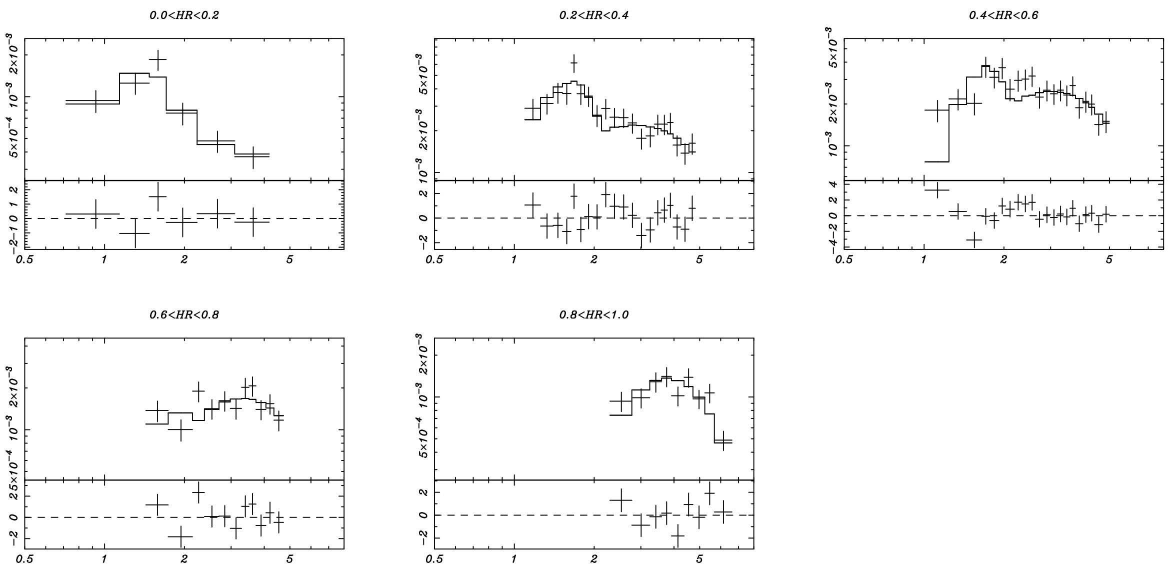


Figure 8.5: Spectra and the best-fit power-law models of the combined spectra of NIR-unIDed *Chandra* sources in the same HR bin from 0.0–0.2 to 0.2–1.0. In the upper panels, the data (*pluses*) and the best-fit model (*solid steps*) are plotted over the energy (keV; *horizontal axis*) versus normalized spectral intensity (count rate keV^{-1}; *vertical axis*) plane. The response functions of the optics and the detector are convolved into the model. In the lower panels, the residuals between the background-subtracted data and the best-fit model are plotted over the energy (keV; *horizontal axis*) versus $\chi^{(i)} = (E^{(i)}_{data} - E^{(i)}_{model \otimes ARF \otimes RMF})/\Delta E^{(i)}_{data}$ (*vertical axis*) plane.

relation of extragalactic X-ray sources as

$$N(>S) = 1200 \times \left(\frac{S}{2 \times 10^{-15} \text{ ergs s}^{-1} \text{ cm}^{-2}} \right)^{-1.0}, \qquad (8.2)$$

where $N(>S)$ is the number of X-ray sources per square degree brighter than S ergs s^{-1} cm^{-2} in the hard X-ray band (2.0–7.0 keV). In Sect. 5.1, we discussed the completeness limit of our *Chandra* observation to be $F_{\rm X} \sim 10^{-14.5}$ ergs s^{-1} cm^{-2} in 0.5–8.0 keV. Assuming a power-law spectrum of $\Gamma = 1.7$ and $N_{\rm H} = 1$–3×10^{22} cm^{-2}, the expected number of the extragalactic sources is 74–81.

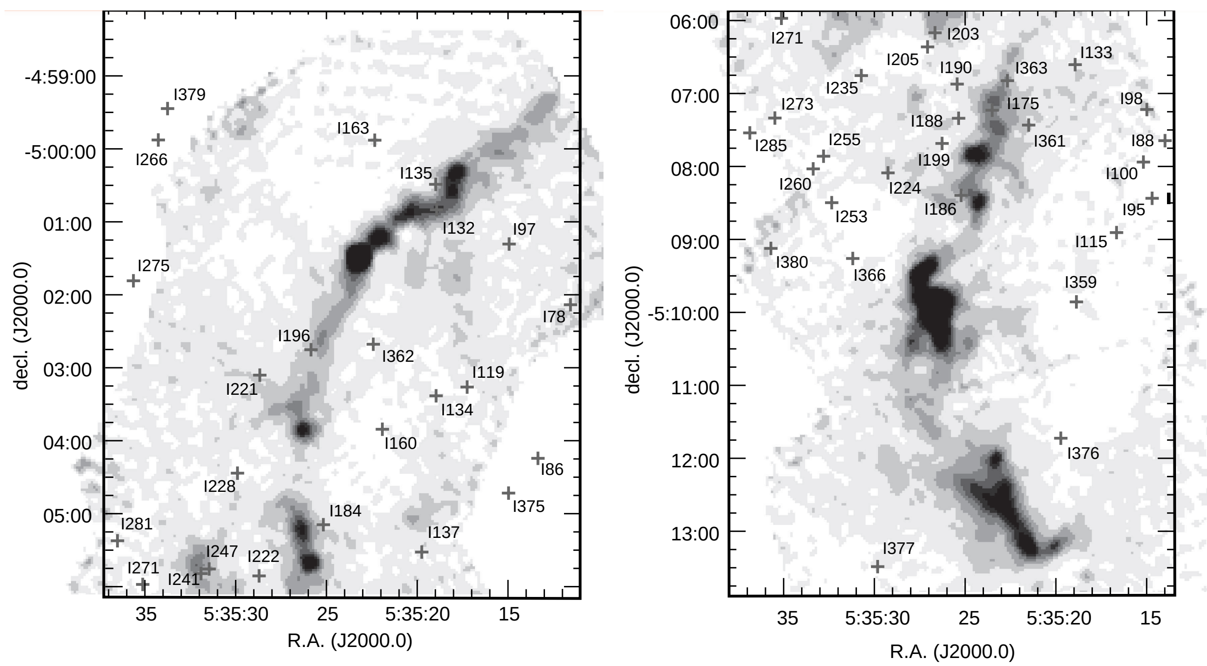
Finally, the spatial distribution of NIR-unIDed sources (Fig. 6.8 d) is not correlated with the 1.3 mm intensity map, but shows a rather uniform distribution. This is in contrast with protostar and cTTS X-ray sources.

8.1.2 Sources with YSO-like Features

There are, however, some exceptions for the NIR-unIDed X-ray sources to be of the extragalactic origin. Some sources show YSO-like features.

Figure 8.4 shows that there are some soft X-ray sources with HR < 0 among NIR-unIDed sources. The overall extinction of our study field is derived to be $A_K \sim 0.5$–1 mag (Fig. 6.4), which can be converted to $N_{\rm H} = 0.5$–1×10^{22} cm^{-2} using the equation (7.6). With this absorption, the power-law spectrum of $\Gamma = 1.7$ has the HR of 0.12–0.40. Even with no absorption, the spectrum has the HR of –0.45. Most of the NIR-unIDed X-ray sources with HR < 0 are unlikely to be AGNs.

Some sources appear to be spatially associated with the 1.3 mm cloud ridge (Fig. 8.6). If all the NIR-unIDed sources were distributed uniformly in the ACIS-I FOV, the number density would be $\sim$0.37 arcmin^2. The expected number of sources within the integral shape of the 1.3 mm intensity ($\sim$9.7 arcmin^2) is thus calculated to be $\sim$3.5. More than twice numbers of NIR-unIDed sources (I132, I135, I175, I186, I196, I203, I241, I247, and I363) are actually within the integral shape, indicating that some of them have physical associations with the cloud cores.



(a) OMC-3 (b) OMC-2

Figure 8.6: Positions of NIR-unIDed *Chandra* sources (*pluses*) with their name over the the 1.3 mm intensity in gray scale (Chini et al. 1997[35]). (*a*) OMC-3 and (*b*) OMC-2.

8.2 H_2 Imaging Observations on OMC-3

8.2.1 Observation

In order to identify the nature of the NIR-unIDed X-ray sources with YSO-like features, we further conducted H_2 mapping observations of OMC-3 with QUIRC (Fig. 8.7). A global map of molecular outflows in star-forming regions is quite useful to see the position and distribution of protostars. QUIRC provides a $3.2' \times 3.2'$ FOV with the pixel scale of $0.189''$ pixel^{-1}. With dithering observations, we took images of three $4.2' \times 4.2'$ regions in two narrow-bands (H_2 at 2.12 μm and *K-continuum* at 2.26 μm). Each FOV was exposed for 60 s on March 13, 2001. The seeing was $\sim 1.0''$.

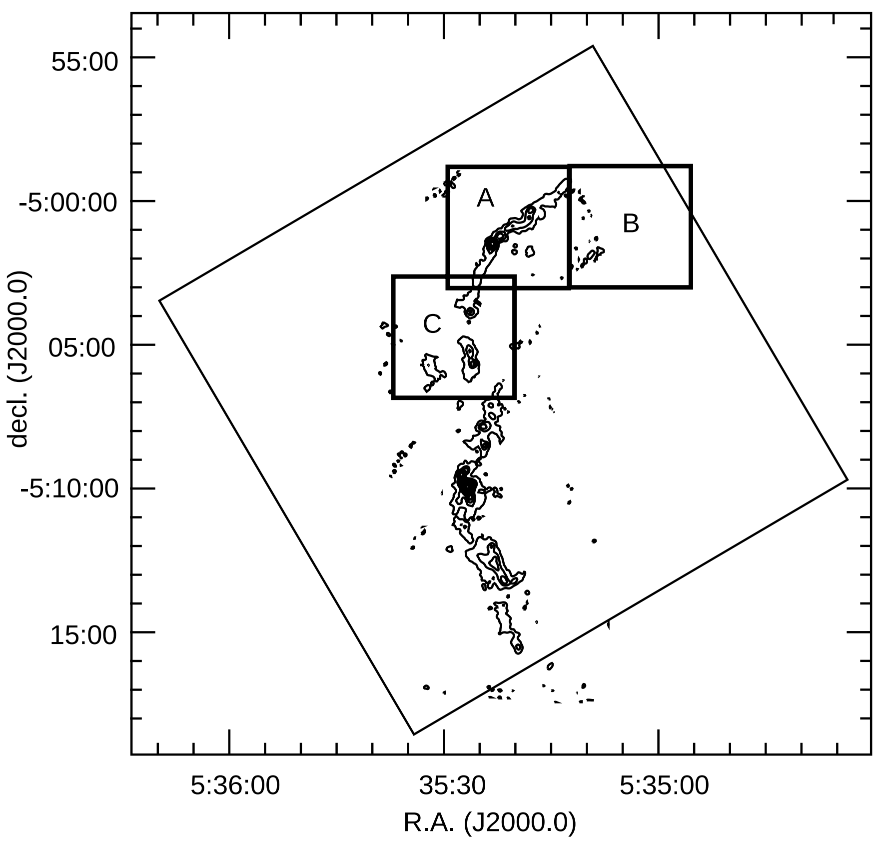


Figure 8.7: FOVs of the H_2 and *K-continuum* observations using QUIRC. They are respectively named region A, B and C. The *Chandra* ACIS-I FOV is also shown with the large oblique square. The contours are the 1.3 mm intensity (Chini et al. 1997[35]).

8.2.2 Analysis & Results

The images were reduced following the standard procedures using IRAF; dark subtraction, flat fielding, sky subtraction, bad pixel removal for each frame, and correction for dithering to construct the final images.

The H_2 and *K-continuum* images in three regions were binned with the neighboring 2×2 pixels and smoothed with a Gaussian function to attain better signal-to-noise ratio. The obtained H_2 maps are shown in Figure 8.8. Some *Chandra* sources (e.g.; I241 and I247) were found be associated both with 1.3 mm intensity and H_2 features.

8.3 NIR Observations on MMS 2 and MMS 3

8.3.1 Observation

As a test case to study the X-ray emission associated with the 1.3 mm cloud core and H_2 outflow, we focus on the X-ray sources at MMS 2 and MMS 3. Figure 8.9 gives a close-up view of MMS 2–MMS 3 region in the hard (3.0–6.0 keV) and soft (0.5–3.0 keV) X-ray band. These two images were binned with 0.1″ pixel^{-1} and smoothed with a Gaussian function of 6 bins to attain the sub-pixel resolution, utilizing the dithering observation of *Chandra*. We see that I128 is separated into four components; three (I128a, I128b and I128c) are in the hard and one (I128d) is in the soft band image. Based on the spatial coincidence with 1.3 mm cores and their hard X-ray spectra, Tsuboi et al. (2001)[188] proposed that I128a at MMS 2 and I132 and MMS 3 are the first candidates of the X-ray-emitting class 0 objects.

Previous observations in the radio and NIR band revealed that there is a molecular outflow (Aso et al. 2000[11]; Yu et al. 1997[199]), an optical jet (Reipurth et al. 1997[161]) and a radio jet (Reipurth et al. 1999[162]) originating from MMS 2, which strongly indicates that a protostar or protostars are embedded in this core. However, the spatial resolutions of these NIR and radio observations (including our QUIRC observations) are not high enough to match the *Chandra* resolution, which makes impossible to discuss the correlation of the X-ray emissions with the outflow and jet sources. Clearly, follow-up observations with much higher spatial resolution, comparable to that of *Chandra* are needed.

Using IRCS at the Cassegrain focus of the Subaru telescope, we took three broad-band

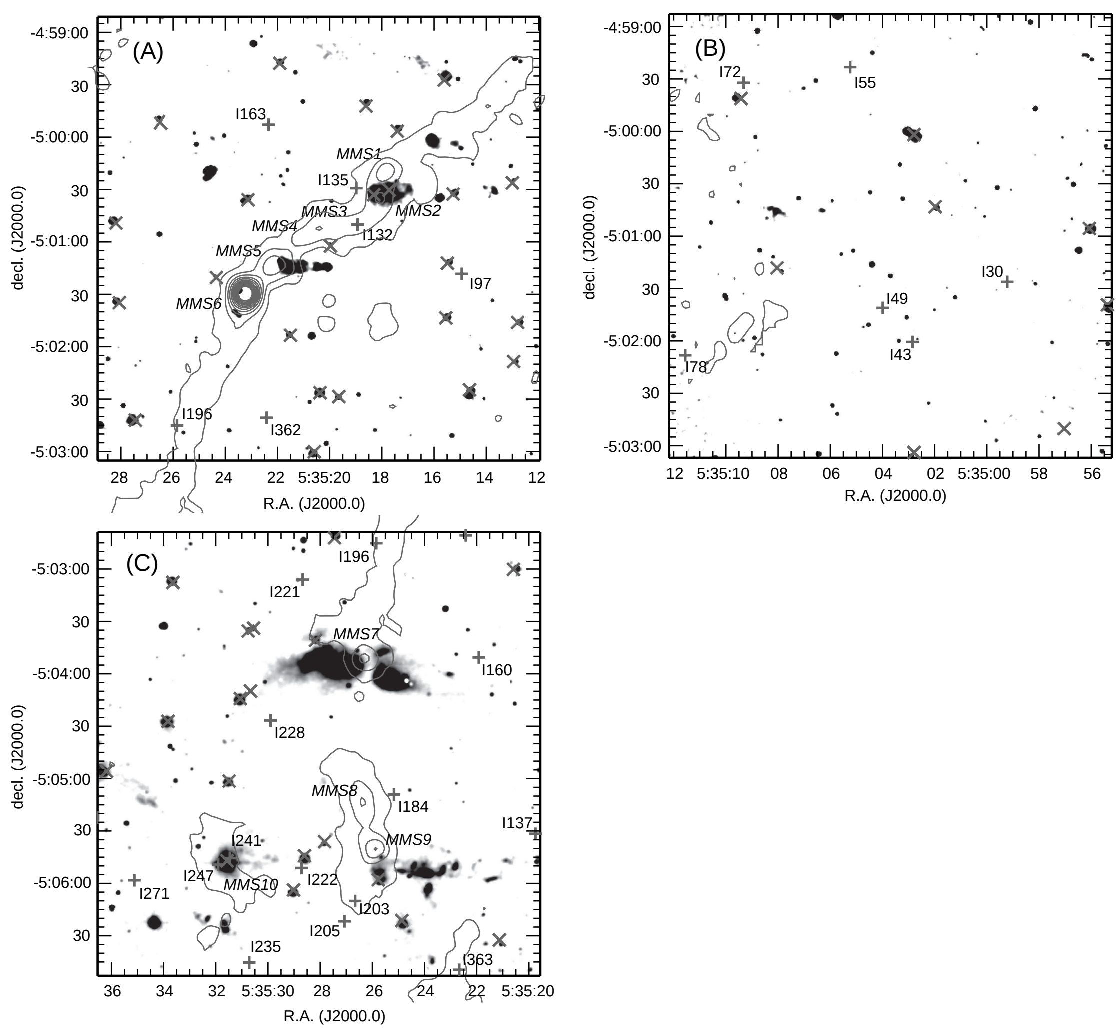
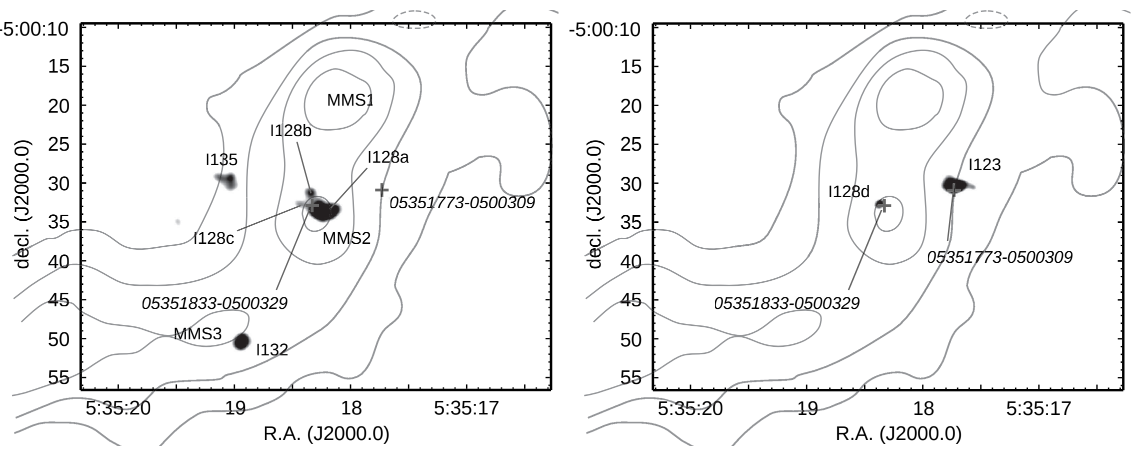


Figure 8.8: QUIRC H_2-band images of three regions in OMC-3. The gray scale gives the H_2 intensity with *K-continuum* not subtracted, while the contours give the 1.3 mm intensity (Chini et al. 1997[35]). The crosses and pluses show the position of NIR-IDed and NIR-unIDed *Chandra* sources. The names of NIR-unIDed X-ray sources are given in Roman, while those of 1.3 mm cores are in Italic.



(a) hard band image

(b) soft band image

Figure 8.9: Close-up images on I128 in MMS 2 for the (*a*) hard (3.0–6.0 keV) and (*b*) soft (0.5–3.0 keV) X-ray band. I128 can be separated into four components; I128a, I128b, and I128c in the hard, and I128d in the soft band image. The positions of QUIRC sources are given with pluses with the prefix in their names ("TKK J") omitted. The contours give the 1.3 mm intensity (Chini et al. 1997[35]).

(*J*, *H*, and *K*-band) and two narrow-band (H_2 and *K-continuum*) images on November 30 and December 4, 2000. The seeing was $\sim 0.5''$ on both nights. The *J*-, *K*-, H_2-, and *K-continuum*-band images were exposed for 600 s, while the *H*-band images were for 300 s.

IRCS provides a FOV of $60''\times60''$ with a pixel scale of $0.058''$ pixel^{-1}. With dithering of five FOVs (Fig. 8.10), we covered a $90''\times90''$ field encompassing both MMS 2 and MMS 3 in the central $30''\times30''$ region. Dithering compensates for the pixel-to-pixel variation in quantum efficiency of the detector and enables to construct the sky image by the median-sky technique without taking a sky frame. In a dithering observation of five FOVs, we have five ADU values at each pixel of the detector. By leaving the median values among the five, we can obtain the median-sky image.

As we had no detection in the *J* band from the two NIR sources at MMS 2, we obtained an additional L'-band image of MMS 2 with NSFCam at the Cassegrain focus of IRTF on December 23, 2000 with the integration time of 216 s. The seeing was $\sim 1.0''$. NSFCam provides a $38''\times38''$ FOV with the pixel scale of $0.148''$ pixel^{-1}. With dithering, we covered a $64''\times64''$ field.

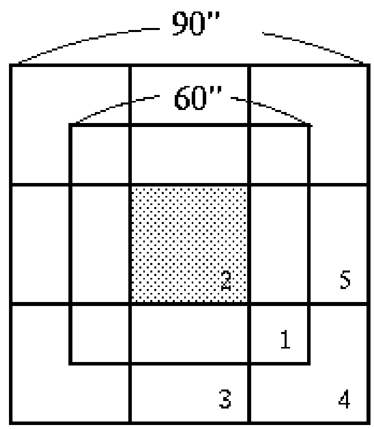


Figure 8.10: Configuration of IRCS dithering observations. Five frames (frame numbers are given at the bottom right of each frame) with $60'' \times 60''$ were combined to construct a $90'' \times 90''$ image. The hatched region at the center was covered by all frames. The same configuration with a different dithering amplitude was employed for QUIRC H_2 and NSFCam observations.

8.3.2 Analysis & Results

Source Extraction and Photometry of Broad-band Images

The images were reduced following the standard procedures using IRAF; dark subtraction, flat fielding, sky subtraction, bad pixel removal for each frame, and correction for dithering to construct a final image (Figs. 8.11 and 8.12).

SExtractor (Bertin & Arnouts 1996[21]) was used for source extraction and photometry. Nine sources (IRS 1–IRS 9) were extracted from the K-band image (Table 8.3). For each K-band detected source, we performed a $1.0''$-aperture photometry in the J, H, and K band. We transformed their magnitudes into the CIT color system in the following way. Seven sources have the counterpart in the Point Source Catalog of the 2MASS Second Incremental Data Release. Referring to their J-, H-, and K_s-band magnitudes, we derived a linear relation between IRCS and 2MASS magnitudes in each band. We first converted the IRCS magnitudes into 2MASS magnitudes using these relations and then into the CIT color system using the formulae given in Carpenter (2001)[27].

For the L'-band image with NSFCam, we performed a $2.0''$-aperture photometry of IRS 3, IRS 4, and IRS 5. We first calculated the magnitudes with the photometric zero-

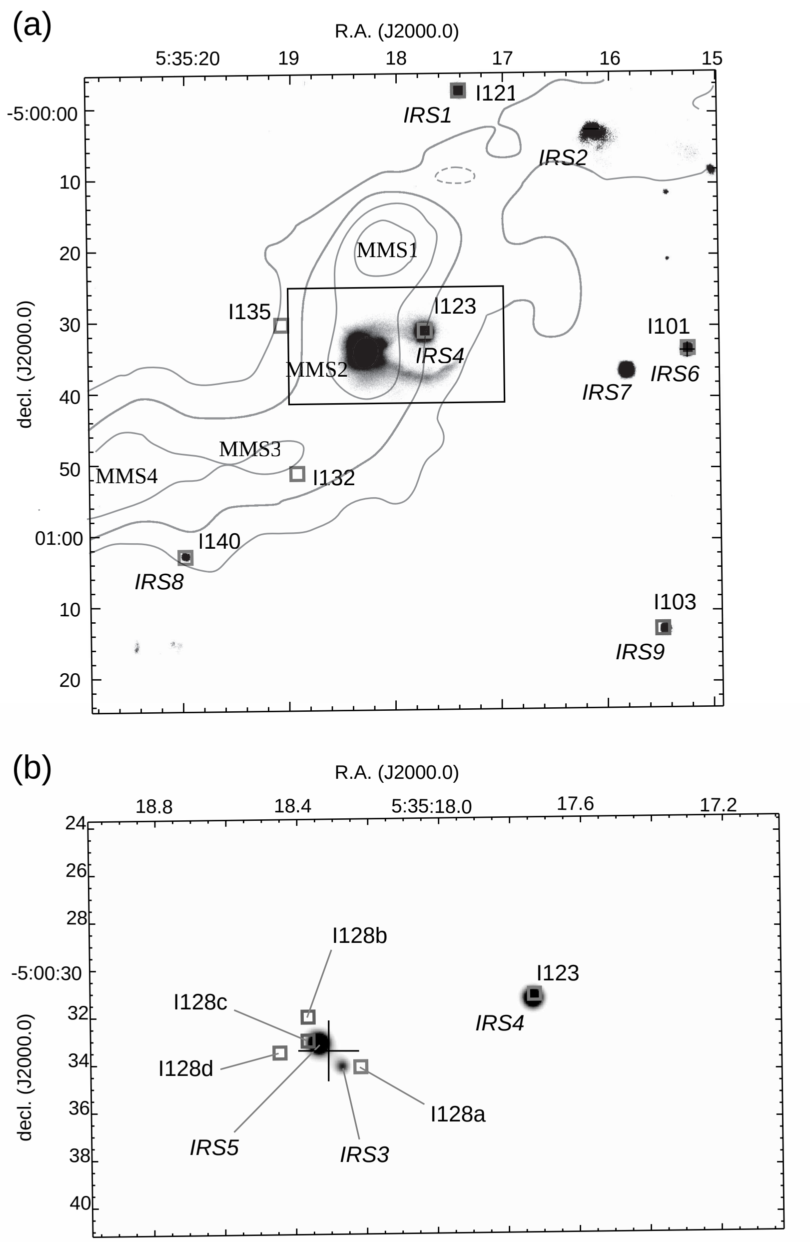


Figure 8.11: (*a*) IRCS *K*-band image with the logarithmic gray scale to stress diffuse features. (*b*) Close-up view of the MMS 2 region (shown in a rectangle in *a*) in the linear scale to show the accurate positions of point sources. The *K*-band sources (IRS 1–IRS 9) are labeled in Italic, while the positions of the X-ray sources are with squares with their names in Roman. The contours in (*a*) are the 1.3 mm intensity. Four 1.3 mm cores (MMS 1–MMS 4) are identified in this region (Chini et al. 1997[35]). The plus in (*b*) shows the position of the 3.6 cm source (Reipurth, Rodríguez, & Chini 1999[162]).

point of 20.3 mag[1], then converted them into the CIT L-band color using[2]

$$(K - L)_{\rm CIT} = 0.820 \times (K - L')_{\rm IRTF}, \tag{8.3}$$

where we assumed $K_{\rm CIT} = K_{\rm IRTF}$ as the first order approximation.

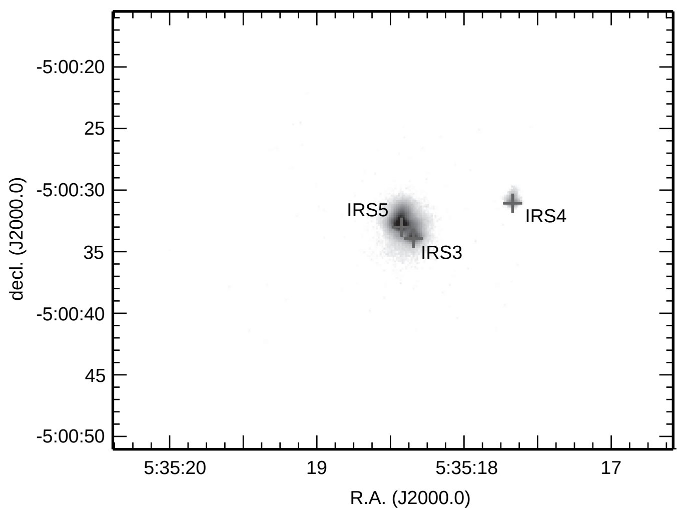


Figure 8.12: NSFCam L'-band image of MMS 2. The positions of the K-band sources (IRS 3–IRS 5) are shown with pluses.

Correlation with X-ray Sources

The X-ray counterpart was searched for each NIR source using Table A.1. From the visual inspection of the NIR and X-ray images, we identified the X-ray sources I121, I123, I101, I140, and I103 to be the counterpart of IRS 1, IRS 4, IRS 6, IRS 8, and IRS 9, respectively.

Two NIR sources (IRS 3 and IRS 5) and four X-ray sources (I128a, I128b, I128c, and I128d in Fig. 8.9. These respectively correspond to the source 8, 8a, 8b, and 8c in Tsuboi et al. 2001[188]) are crowded at MMS 2. In order to find the X-ray counterpart of the NIR sources, we adjusted the X-ray image by a shift and a rotation so that each X-ray source (I101, I103, I121, I123, and I140) comes closest to its NIR counterpart. After this procedure, the positional offset between the NIR sources and their X-ray counterparts is $\sim$0.25$''$ (1 σ).

[1]See http://irtf.ifa.hawaii.edu/Facility/nsfcam/hist/backgrounds.html.
[2]See http://irtf.ifa.hawaii.edu/Facility/nsfcam/hist/color.html.

Then, I128d is found to be the closest source to IRS 5 with the separation of $0.46''$, hence is the X-ray counterpart of IRS 5. On the other hand, IRS 3 is separated by $0.81''$ from the closest X-ray source; I128a. Assuming that the separation between a NIR and X-ray counterpart pair follows a Gaussian distribution of $\sigma = 0.25''$, the separation between IRS 3 and I128a is more than 3 σ. We therefore conclude that I128a is not the X-ray emission from IRS 3. In fact, no separation larger than $0.81''$ is found in any other NIR and X-ray counterpart pairs. I132 at MMS 3, as well as I128a at MMS 2, has no NIR counterpart.

We chose several source-free regions near the positions of I128a and I132 for a $1.0''$-aperture photometry in order to estimate the background level. We found the K-band upper limit of I128a and I132 to be ∼19.6 mag at the 3 σ level.

Table 8.3: IRCS & NSFCam sources

ID	R.A.[a] (J2000.0)	decl.[a] (J2000.0)	J^{b} (mag)	H^{b} (mag)	K^{b} (mag)	L'^{b} (mag)	2MASS[c] identification	X-ray[d] identification
IRS 1	05:35:17.415	−04:59:57.24	17.8	14.7	13.0	.....	0535174–045957	I121 (6)
IRS 2	05:35:16.168	−05:00:02.58	17.0	14.1	12.3	.....	0535161–050002	....
IRS 3[e]	05:35:18.275	−05:00:33.93	>19.6	18.0	13.2	9.17	0535183–050033	....
IRS 4	05:35:17.736	−05:00:31.07	15.4	13.2	11.7	10.5	0535177–050031	I123 (7)
IRS 5[e]	05:35:18.340	−05:00:33.01	>19.6	15.8	11.4	7.86	0535183–050033	I128d (8c)
IRS 6	05:35:15.265	−05:00:33.47	13.3	12.7	12.5	.....	0535152–050033	I101 (2)
IRS 7	05:35:15.837	−05:00:36.34	12.3	11.8	11.7	.....	0535158–050036	....
IRS 8	05:35:19.980	−05:01:02.64	>19.6	>18.8	14.4	.....	0535199–050102	I140 (12)
IRS 9	05:35:15.463	−05:01:12.59	14.2	13.6	13.4	.....	0535154–050112	I103 (3)

[a] The positions are determined from the IRCS K-band image in the equinox J2000.0.

[b] All magnitudes are in the CIT color system.

[c] 2MASS source names with "2MASSI J" omitted for the prefix. IRS 3 and IRS 5 are not resolved in the 2MASS data.

[d] Given in parentheses are the nomenclatures in Tsuboi et al. (2001)[188].

[e] Associated with MMS 2.

Narrow-band Images

The vibrational-rotational transition of $v = 1 - 0$ S(1) works as an effective coolant of the excited hydrogen molecules. Therefore, this emission line serves as a powerful tool to search for jets from a protostar and the position of its powering source (Bally et al. 1993[15]; Hodapp & Ladd 1995[88]). In the continuum-subtracted H_2-band image, we identified a bubble-like feature originating from MMS 2. A close-up view of this bubble-like emission is shown in Figure 8.13, where we see the origin of this feature spatially coincides with I128a. No similar feature was found for I132 at MMS 3.

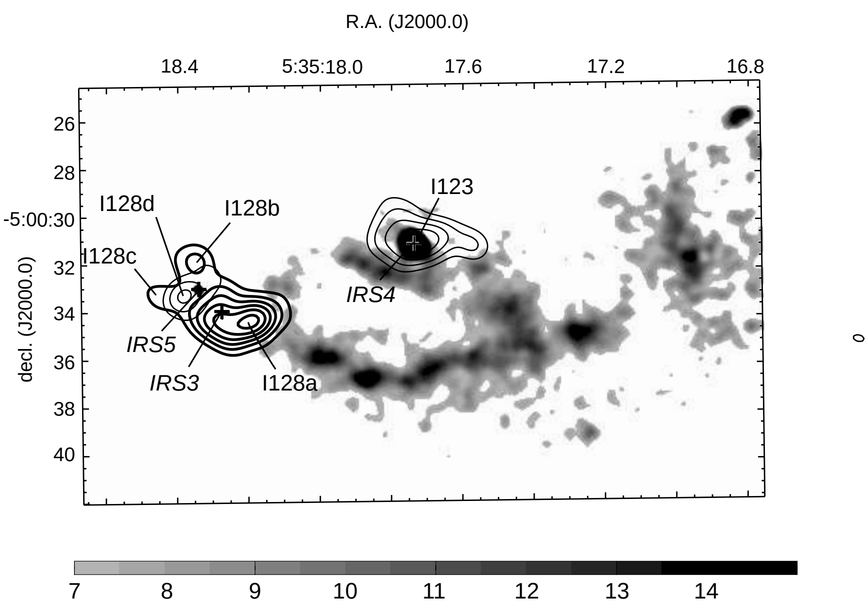


Figure 8.13: Continuum-subtracted H_2 intensity (gray scale). The hard (3.0–6.0 keV) and soft (0.5–3.0 keV) X-ray intensity are shown with thick and thin contours. The *K-continuum* image is multiplied by a factor and subtracted from the H_2 image, so that the emissions from IRS 3 and IRS 5 cancel out. Without *K-continuum* subtraction, however, we confirmed the same bubble-like feature in the H_2 image. The scale bar at the bottom is in the unit of intensity pixel^{-1}, where the background level (*white*) is $\sim$2.5. The positions of the *K*-band sources are shown with pluses. The X-ray and NIR sources are labeled in Roman and Italic, respectively.

8.3.3 Discussion

The Nature of NIR Sources

For the classification of IRCS sources, we used the color-color diagram (Lada & Adams 1992[114]). The $(J–H)/(H–K)$ diagram is given in Figure 8.14 (a). Since IRS 3 and IRS 5 have no detection in the J band, we also give the $(H–K)/(K–L)$ diagram in Figure 8.14 (b).

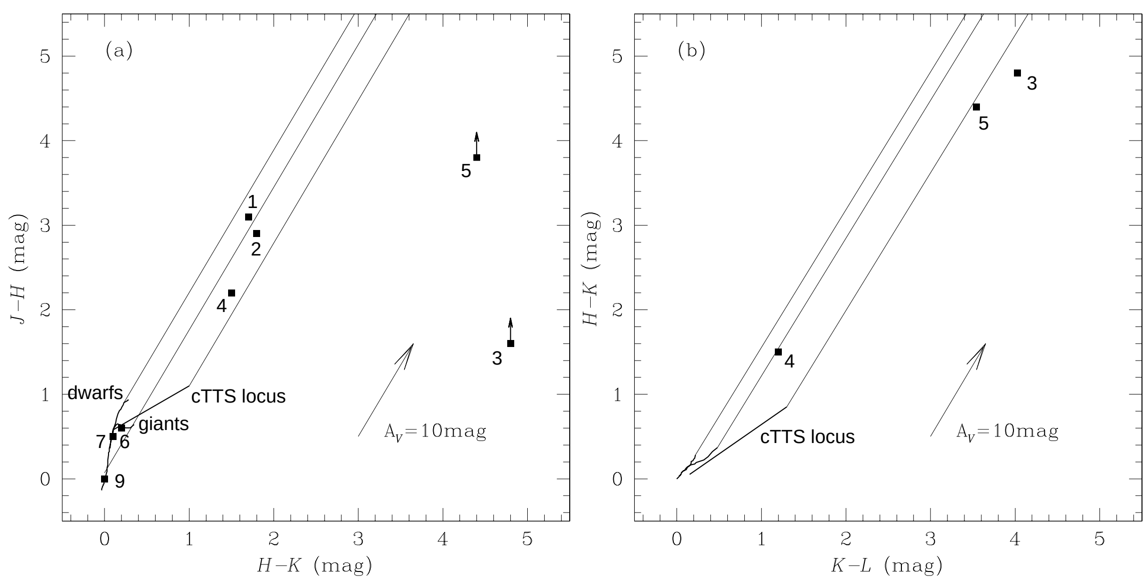


Figure 8.14: (a) $(J–H)/(H–K)$ color-color diagram and (b) $(H–K)/(K–L)$ color-color diagram. IRS 1–IRS 9 are plotted in the CIT color system with the label of their names ("IRS" is omitted). The uncertainty is less than roughly $\pm$ 0.1 mag for each color. The intrinsic colors of dwarfs and giants are given with thick solid curves (Tokunaga 2000[185]), and the cTTS locus is with the thick solid line (Meyer et al. 1997[130]). Their extinction vectors are given with the thin solid lines. We assumed the slope of the reddening lines to be $E(J–H)_{\rm reddening}/E(H-K)_{\rm reddening} = 1.69$ and $E(H-K)_{\rm reddening}/E(K-L)_{\rm reddening} = 1.63$ (Meyer et al. 1997[130]). The A_V of each source is estimated from $E(H-K)_{\rm reddening} = 0.065 \times A_V$ or $E(K-L)_{\rm reddening} = 0.04 \times A_V$ (Meyer et al. 1997[130]).

IRS 3 and IRS 5 are at the center of a millimeter core (MMS 2) and are located $\sim$1.34$''$ ($\sim$600 AU at the distance of 450 pc) apart from each other (Fig. 8.11). Together with their large extinction of more than $A_V > 50$ mag and large NIR excess seen in Figure 8.14 (b), they are class I protostars probably comprising a binary system.

IRS 1, IRS 2 and IRS 4 are at the reddening region of the cTTS locus with a moderate extinction of $A_V \sim 30$ mag (Fig. 8.14 a). They are located at the edge of 1.3 mm cores

(Fig. 8.11), thus are most likely to be cTTSs.

IRS 6, IRS 7, and IRS 9 are located away from the cloud cores (Fig. 8.11) and have less extinction (Fig. 8.14 a). Among them, IRS 6 and IRS 9 are considered to be wTTSs due to the association with X-ray emissions. IRS 7, which has no X-ray counterpart, may be a back- or foreground source.

It is hard to infer the nature of IRS 8 from NIR observation alone because it has only the K-band detection. However, its X-ray counterpart (I140) shows a thermal emission of $k_{\rm B}T = 3.12$ keV, $L_{\rm X} = 6.32\times10^{30}$ ergs s^{-1}, and $N_{\rm H} = 6.22\times10^{22}$cm^{-2} (Table 7.1). These are typical values for class I sources (Table 9.5), hence this source is most likely to be a class I protostar.

8.4 Centimeter Observations on MMS 2 and MMS 3

8.4.1 Observation

We further took a centimeter image on the MMS 2 region. The purpose of this observation is to determine the position of the protostellar core with an accuracy of 0.1″ and to compare it with the X-ray image. Protostars are often accompanied by free-free emissions within 100 AU (0.2″ at MMS 2), which is detectable mostly as point-like by the centimeter continuum imaging observations (Anglada et al. 1992[7]; Rodríguez, Anglada, & Raga 1995[164]). When observed with long-baseline interferometer observations, centimeter imaging is the most accurate method to determine the position of protostars, providing a vital clue to discuss the mechanism of the X-ray emissions from the youngest phase of protostars.

We conducted our centimeter observation with VLA on February 11, 2002. We used the A configuration to achieve the highest possible spatial resolution. A 3.6 cm map was obtained with the integration time of ∼3.5 hours, the band width of 50 MHz, and the phase center at R.A. = $05^{\rm h}35^{\rm m}18.3^{\rm s}$ and decl. = –05°00′33″. The map is sensitive to the structure smaller than ∼2″ with the angular resolution of ∼0.2″. 3C 48 (3.25 Jy) and 0541−056 (0.98 Jy) were used as the flux and phase calibrator, respectively.

8.4.2 Analysis & Results

Data reduction, calibration and analysis were performed using AIPS[3]. The natural weighted map is shown in Figure 8.15. We detected two sources (VLA 1a and VLA 1b) above the 3 σ level, for which we derived the position and the flux density (Table 8.4). VLA 1a is slightly extended, so we also determined the length of the major (minor) axis and position angle to be 0.40″ (0.09″) and 32.2 degree, respectively. These two sources were not resolved in the prior D configuration observation and were named altogether as VLA 1 (Reipurth et al. 1999[162]).

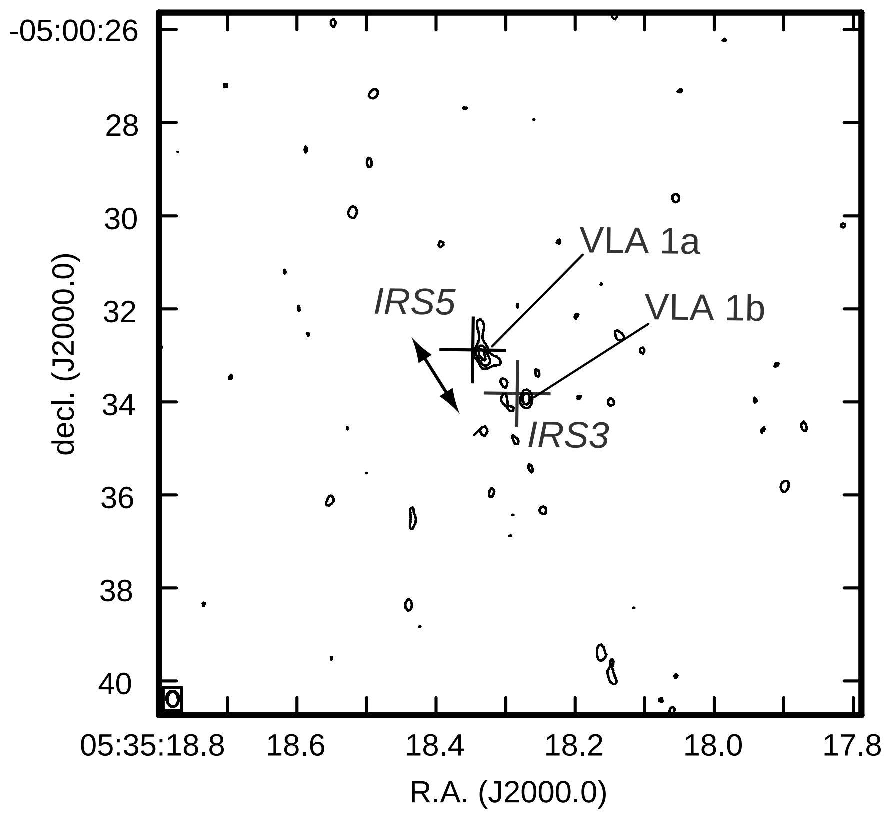


Figure 8.15: The 3.6 cm image of MMS 2 and MMS 3 with VLA. *Contours*: the 3.6 cm intensity. The contour levels are 3–9 σ with the step of 3 σ, where the background noise is $\sim$7.7 μJy beam^{-1}. The synthesized beam size is at the bottom left. *Pluses*: the positions of NIR sources. *Arrow*: the position angle of VLA 1a.

[3]See http://www.cv.nrao.edu/aips/.

Table 8.4: VLA sources

ID	R.A. (J2000.0)	decl. (J2000.0)	flux[a] (μJy)	major axis ($''$)	minor axis ($''$)	IRCS counterpart
VLA 1a	05:35:18.335	−05:00:32.97	128	0.40	0.09	IRS 5
VLA 1b	05:35:18.271	−05:00:33.93	54	....	....	IRS 3

[a] The flux density is corrected for the primary beam response.

8.4.3 Discussion

The Nature of Centimeter Sources

For the following reasons, we conclude that both VLA 1a and VLA 1b are free-free emissions from H_{II} regions ionized by the UV radiation from the shock front produced by the collision of a protostellar jet upon a dense ISM obstacle (Curiel et al. 1987[42]).

First, these centimeter sources are considered to be the counterpart of class I protostars (IRS 3 and IRS 5) from their proximity. The embedded protostars are frequently associated with centimeter emissions, and detailed studies indicate that most of them, if not all, are of free-free emission origin (Anglada 1996[8]).

Second, IRS 3 and IRS 5 have the NIR magnitudes and colors of $J_0 > 11.3$ mag and J–$H > 1.6$ mag, and $J_0 > 11.3$ mag and J–$H > 3.8$ mag, respectively (Table 8.3). This indicates that both sources have a mass less than 2 $M_\odot$, which rules out the possibility that the centimeter emissions are from H_{II} regions generated by stellar UV photons.

Third, the flux density multiplied by the square of the distance ($S_\nu D^2$) and the momentum rate in the outflow (dP/dt) of VLA 1a and VLA 1b fit well with a known empirical relation (Anglada et al. 1992[7]) and theoretical understandings (Curiel et al. 1987[42]) between these two parameters. Figure 8.16 shows the relation between $S_\nu D^2$ and dP/dt of 16 embedded objects, where we added VLA 1a/b with $S_\nu D^2 = 0.182 \times 0.45^2$ mJy kpc^2 (Table 8.4) and $dP/dt = 3 \times 10^{-5}$ $M_\odot$ yr^{-1} km s^{-1} (Table 3.2; Aso et al. 2000[11]). The limited spatial resolution of the HCO^+ and CO observations could not determine which of the three 1.3 mm clumps (MMS 2–MMS 4) is responsible for the molecular outflow. However, VLA 1a and VLA 1b in MMS 2 are the only 3.6 cm sources that are associated with MMS 2–MMS 4 (Reipurth et al. 1999[162]). We can therefore safely assume that the dP/dt

value determined by the molecular outflow represents the sum of the momentum rate from VLA 1a and VLA 1b. Similarly, we summed the flux density of VLA 1a and VLA 1b for the $S_\nu D^2$ value.

Fourth, in case of VLA 1a, the emission is elongated (the arrow in Fig. 8.15) along the direction of global outflow seen in the H_2 $v = 1 - 0$ S(1) band (Fig. 8.13), which is characteristic for the free-free centimeter emissions of shock induced plasma (Anglada 1996[8]).

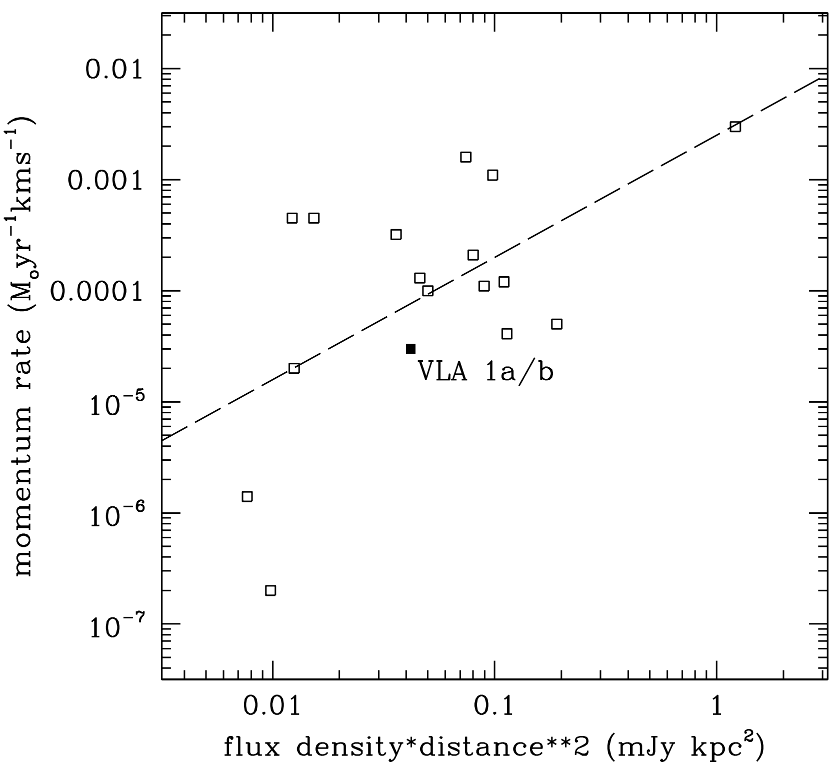


Figure 8.16: Relation between the flux density multiplied by the square of the distance ($S_\nu D^2$) and the outflow momentum rate (dP/dt). Open squares are from Anglada et al. (1992)[7], who derived an empirical relation among these sources (*dashed line*). The filled square (VLA 1a/b) is roughly consistent with this relation.

Chapter 9

Discussion

Contents

In this chapter, we discuss the X-ray emission mechanisms of X-ray sources based on their multi-wavelength features described in the previous chapters. The first two sections are for NIR-IDed X-ray sources. In Sect. 9.1, we compare the X-ray properties in different mass ranges (Sect. 9.1) and discuss that the IM, LM, and VLM sources have the same X-ray emission mechanisms in contrast to HM sources. In Sect. 9.2, we describe that their X-ray emissions consist of two components of different temperatures and propose that the hard component originates from flares while the soft component is from stellar coronae. In Sect. 9.3, we deal with NIR-unIDed X-ray sources, and interpret their X-ray emission in the context of jet-induced plasma and the magnetic activities of deeply embedded NIR invisible YSOs.

9.1 X-ray Emissions from NIR-IDed Sources: (1) Mass

9.1.1 X-ray Properties among Mass Ranges

Based on the mass estimates of NIR-IDed X-ray sources using the $J/(J\text{–}H)$ color-magnitude diagram, we separated them into four mass ranges (HM, IM, LM, and VLM) in Sect. 6.2. The NIR sources in the ACIS-I FOV are separated in the same manner in order to calculate the X-ray detection rate of each mass range. Readers should note that NIR sources classified as LM and VLM are contaminated by back- or foreground sources (Sect. 6.1), which gives the lower limit of the X-ray detection rate of these groups. The number of NIR and X-ray sources and the detection rate are summarized in Table 9.1 along with the results of the X-ray temporal and spectral analyses. The averaged X-ray properties are in Table 9.2.

9.1.2 High Mass Sources

We have one HM source in our sample. The source; ν Ori has the spectral type of B1 V (Kukarkin et al. 1971[111]), the mass of $10^{1.05} M_\odot$ and the bolometric luminosity of $10^{4.06} L_\odot$ (Greenstein 1998[71]). The X-ray counterpart (I242) shows a stable light curve with a soft spectrum. This source is faint (T) and faint (S), so no temporal and spectral analyses were performed (Sects. 7.1 and 7.2). In order to compare with sources in other mass ranges, we examined the temporal variation of I242 and found that the constant flux model was not rejected with $\alpha > 0.05$ (Fig. 9.1). We also fitted the spectrum in the range of 0.5–2.0 keV

Table 9.1: Number of sources among mass ranges

	HM	IM	LM	VLM	sum
NIR sources	1	26	210	462	699
X-ray sources	1	21	139	107	268
	(100%)	(81%)	(66%)	(23%)	(38%)
NIR sources with NIR excess . .	0	12	45	74	131
X-ray sources with NIR excess	0	11	31	11	53
	(N/A)	(92%)	(69%)	(15%)	(40%)
bright (T)	1	16	78	18	113
variable	0	12	42	7	66
	(0%)	(75%)	(54%)	(39%)	(58%)
bright (S)	0	18	94	22	134
one-temperature	0	10	57	13	80
	(N/A)	(56%)	(61%)	(59%)	(60%)
two-temperature	0	8	27	6	41
	(N/A)	(44%)	(29%)	(27%)	(31%)

Table 9.2: Comparison of X-ray properties among mass ranges

	HM[a]	IM	LM	VLM
$< \log N_{\rm H}$ (cm^{-2}) $>$	21.5	21.9±0.4	21.6±0.7	21.7±0.6
$< kT^{(1)}$ (keV) $>$	0.64	3.85±1.6	2.30±2.0	2.04±2.3
$< \log L_{\rm X}^{(1)}$ (ergs s^{-1}) $>$ . . .	30.3	30.5±0.4	30.1±0.5	29.8±0.4
$< kT_{high}^{(2)}$ (keV) $>$	N/A	4.00±2.1	2.75±1.2	2.41±1.1
$< \log L_{X\ high}^{(2)}$ (ergs s^{-1}) $>$	N/A	30.6±0.7	30.3±0.4	29.7±0.2
$< kT_{low}^{(2)}$ (keV) $>$	N/A	0.89±0.2	0.77±0.3	0.44±0.3
$< \log L_{X\ low}^{(2)}$ (ergs s^{-1}) $>$	N/A	30.5±0.5	30.1±0.5	29.9±0.7

[a] The result of the spectral fitting of I242 is shown.

to avoid background contamination in the hard X-ray band to have an acceptable fit with a thin-thermal plasma mode. The best-fit parameters are given in Table 9.3, while the spectrum and the best-fit model are in Figure 9.1.

The X-ray emissions from earlier-type main sequence stars than B2 are explained by the stellar wind model (Lucy & White 1980[118]; Lucy 1982[119]), in which the strong stellar wind propagating through the ambient matter at the speed of $\sim$1000 km s^{-1} ionizes the gas with shocks.

The past X-ray observations on this class of stars revealed that they show non-variable light curves, softer spectra than $k_{\rm B}T \sim 1$ keV, and $L_{\rm X}/L_{\rm bol} = 10^{-7.1}$–$10^{-7.6}$ (Berghöfer et al. 1997[20]; Feigelson et al. 2002[54]). The X-ray features of I242 follow very well with that of the typical early-type stars with $k_{\rm B}T = 0.64$, $L_{\rm X}/L_{\rm bol} = 10^{-7.3}$, and an non-variable light curve. We conclude that I242 is a high-mass main sequence source and the X-ray emission from it is of stellar wind origin.

Table 9.3: One-temperature plasma fittings of the high mass source

ID	counts[a]	S/N[b]	$N_{\rm H}$[c] (10^{22} cm^{-2})	$k_{\rm B}T$[c] (keV)	$L_{\rm X}$[a] (ergs s^{-1})
I242	660	8.54	0.34 (0.25–0.39)	0.64 (0.58–0.69)	2.14e+30

[a] Values in the 0.5–8.0 keV range.
[b] Values in the 2.0–8.0 keV range.
[c] The lower and upper limit (1 σ) are given in parentheses.

9.1.3 Intermediate Mass to Very Low Mass Sources

Previous X-ray observations mainly focused on the X-ray emissions from LM YSOs, revealing that they are of thin-thermal plasma origin generated and maintained by magnetic activities on the stellar surface. They show hard and strong X-ray emissions of $k_{\rm B}T = 0.5$–2.0 keV and $L_{\rm X}/L_{\rm bol} = 10^{-2}$–$10^{-5}$ with occasional flare-like variability (Feigelson & Montmerle 1999[53]). Less is known for higher mass YSOs because they have lower population and evolve more quickly, making these samples fewer and more distant. The X-ray emissions from VLM sources are also behind our understandings on LM YSOs because they are fainter. Using our VLM, LM, and IM samples, we discuss that YSOs in these mass ranges have the same X-ray emission mechanism.

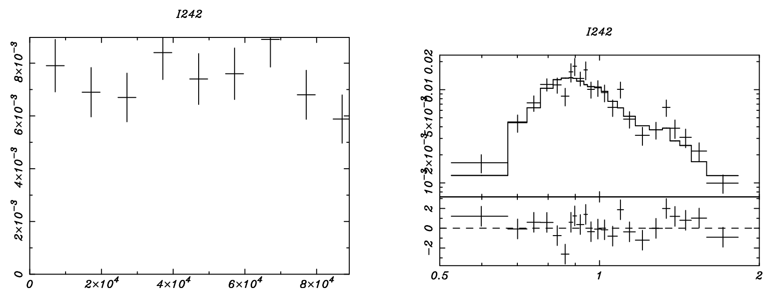


Figure 9.1: (*left*) Light curve of the high mass source (I242) over the count rate (s^{-1}; *vertical axis*) versus the time from the start of the observation (s; *horizontal axis*) plane. (*right*) Spectra and the best-fit models of one-temperature thin-thermal plasma fittings of the high mass source (I242). The metallicity of all elements is fixed to be 0.3 solar. In the upper panel, the data (*pluses*) and the best-fit model (*solid steps*) are plotted over the energy (keV; *horizontal axis*) versus normalized spectral intensity (count rate keV^{-1}; *vertical axis*) plane. The response functions of the optics and the detector are convolved into the model. In the lower panel, the residuals between the background-subtracted data and the best-fit model are plotted over the energy (keV; *horizontal axis*) versus $\chi^{(i)} = (E^{(i)}_{data} - E^{(i)}_{model\otimes ARF\otimes RMF})/\Delta E^{(i)}_{data}$ (vertical axis) plane.

First, we rule out the possibility that the X-ray emissions from IM YSOs are from their LM companion, which is the most favored scenario for the X-ray emissions from IM main sequence stars (Berghöfer & Schmitt 1994[20]). This is because IM main sequence stars have no mechanism to generate high temperature plasma. We compared the X-ray detection rate of our IM samples with the binary rate of the Orion and found that the former (∼80%) is much higher than the latter (∼15%). The binary rate is presented by Padgett et al (1997)[146], who derived the value based on the *Hubble Space Telescope* observations on three Orion stellar clusters (NGC 2024, NGC 2068, and NGC 2071). They examined 99 sources down to $I<$ 19 mag and found that 15 of them have a binary companion in the separation range of $0.3'' < \theta < 2.3''$ (∼100–1000 AU). As this rate gives the upper limit of IM sources to have a LM binary, all of the X-ray emissions from our IM sources can not be attributable to their companion sources.

We can confine our IM samples to those robustly in the pre-main-sequence stage by picking up sources with the NIR excess (the remainders are the mixture of main and pre-main sequence sources). All but one IM sources with NIR excess are found to have X-ray detections, indicating that virtually all IM YSOs emit X-rays.

We have three lines of evidence to conclude that the same X-ray emission mechanism functions for VLM, LM, and IM YSOs. First, all VLM–IM sources in our sample show similar X-ray features when averaged over the group (Table 9.2). About ~60% of them show one-temperature plasma spectra of 2–3 keV and ~30% show two-temperature plasma with the combination of ~1.0 keV and 2–3 keV. This is in contrast to the HM source (I242), which only shows low-temperature (~0.5 keV) plasma spectrum.

Second, the ratio of the X-ray and bolometric luminosity ($L_{\rm X}/L_{\rm bol}$) of most of these sources is in the range of 10^{-2}–10^{-5} (Fig. 9.2). The value is consistent with the previous observational results on LM YSOs (Feigelson & Montmerle 1999[53]) and higher than that of the typical HM main sequence stars with stellar wind origin ($\sim 10^{-7}$). In fact, the only HM (and some IM) source has $L_{\rm X}/L_{\rm bol} \sim 10^{-7}$. The trend of lower $L_{\rm X}/L_{\rm bol}$ values toward higher mass sources (Fig. 9.2) is due to the observational bias. We are dealing with the X-ray-selected samples, which causes lower mass (smaller bolometric luminosity) sources not to be detected even with the same $L_{\rm X}/L_{\rm bol}$ value. The bias is illustrated with the dashed curve representing the typical X-ray detection limit of 10^{29} ergs s^{-1} (Sect. 5.1).

Third, the $L_{\rm X}$ values of VLM–IM sources increase toward the higher mass sources (Fig. 9.3). An empirical relation between these two parameters was presented by Preibisch, & Zinnecker (2002)[156] for VLM–LM YSOs in a low-mass star forming region (IC 348). We see that the relation can be extrapolated to apply for IM and VLM sources.

All these arguments infer that VLM–IM sources (including young brown dwarfs) are emitting X-rays with the same mechanism and the level of activity can be scaled with their mass.

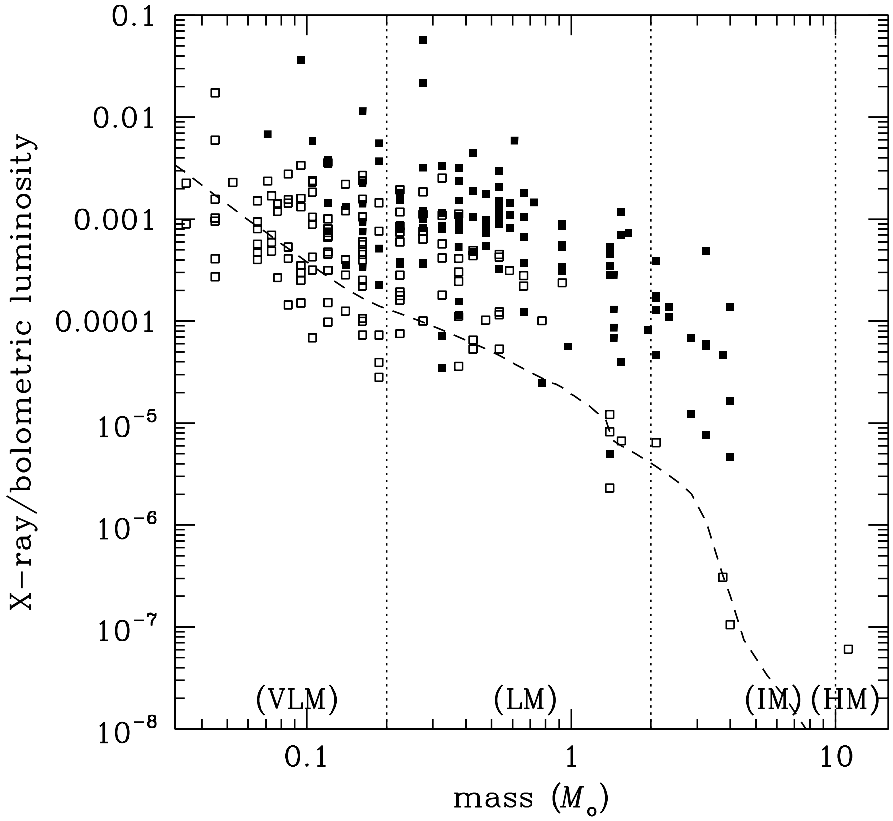


Figure 9.2: Relation between the mass (M) and the ratio of the X-ray (0.5–8.0 keV) and bolometric luminosity (L_X/L_{bol}) of NIR-IDed X-ray sources. Filled squares are with the X-ray luminosity derived from the spectral fittings, while open squares are with the L_X values estimated from their X-ray counts using the equation (7.8). The dashed curve shows the typical detection limit of $L_X = 10^{29}$ ergs s^{-1}.

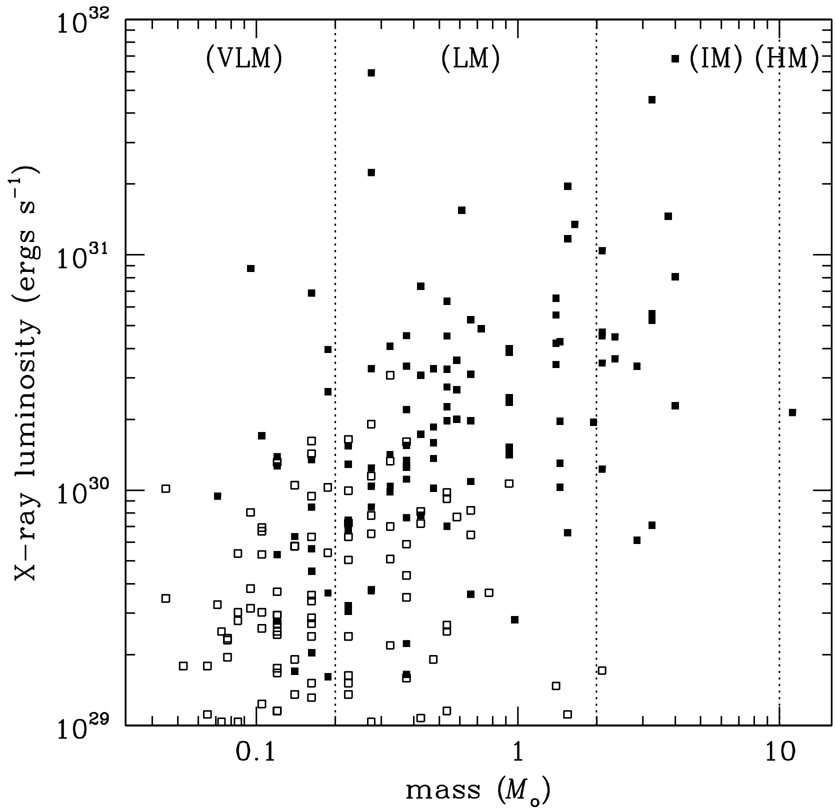


Figure 9.3: Relation between the mass (M) and the X-ray luminosity (L_X) of NIR-IDed X-ray sources in the 0.5–8.0 keV. Filled squares are with the luminosity derived from the spectral fittings, while open squares are with the L_X values estimated from their X-ray counts using the equation (7.8).

9.2 X-ray Emissions from NIR-IDed Sources: (2) Plasma Temperature

9.2.1 Two-temperature Plasma Emissions

The previous discussion that the same X-ray emission mechanism works for VLM–IM sources justifies us to deal with them collectively. We consider, for the following reasons, that these sources have the combination of two X-ray emission mechanisms of different temperature.

Figure 9.4 (a) shows the histogram of the plasma temperatures of VLM–IM sources, where we count the sources with two-temperature plasma twice at each temperature. We can see two peaks at $k_{\rm B}T \sim 1$ keV and 2–3 keV. According to the standard magnetic reconnection model of solar flares, the plasma temperature at a flare (T_f) is described as a function of the pre-flare density (n_0), the loop length (L), and the magnetic field strength (B); e.g.,

$$T_f = 2 \times 10^7 \left(\frac{B}{0.003\ \mathrm{T}}\right)^{\frac{6}{7}} \left(\frac{n_0}{10^{-15}\ \mathrm{m}^{-3}}\right)^{-\frac{1}{7}} \left(\frac{L}{10^7\ \mathrm{m}}\right)^{\frac{2}{7}}\ [\mathrm{K}] \tag{9.1}$$

(Yokoyama & Shibata 1998[197]). Taking into account that these parameters can change continuously, this histogram should appear with one peak with broad tails. It is more natural to understand, therefore, that these sources have two different X-ray emission mechanisms, each of which high and low temperature plasma are attributable to.

This is reinforced by examining sources with two-temperature plasma. Figure 9.4 (c) shows the histogram of plasma temperatures separately for the lower (*solid*) and higher (*dashed*) temperature component. The peak of the lower temperature component is at $k_{\rm B}T \sim 1$ keV, while that of the higher temperature component is at $k_{\rm B}T$ = 2–3 keV. These two components coexist at flare and quiescent phases, because for the reason that the time-sliced spectra of some brightest sources require a two-temperature plasma model in both phases (Sect. 7.2). It is not the case, therefore, that these X-ray emissions are attributable to only one component that shows a higher temperature at flare phases and a lower temperature at quiescent phases.

It is also unlikely that the sources with two-temperature plasma are binaries. First, we confirmed both in the X-ray and the K-band images that only four (I42, I211, I256, and I314) among the 41 two-temperature sources can be contaminated with their close companion. Second, among the 30 bright (S) sources with more counts than 1000 (sources

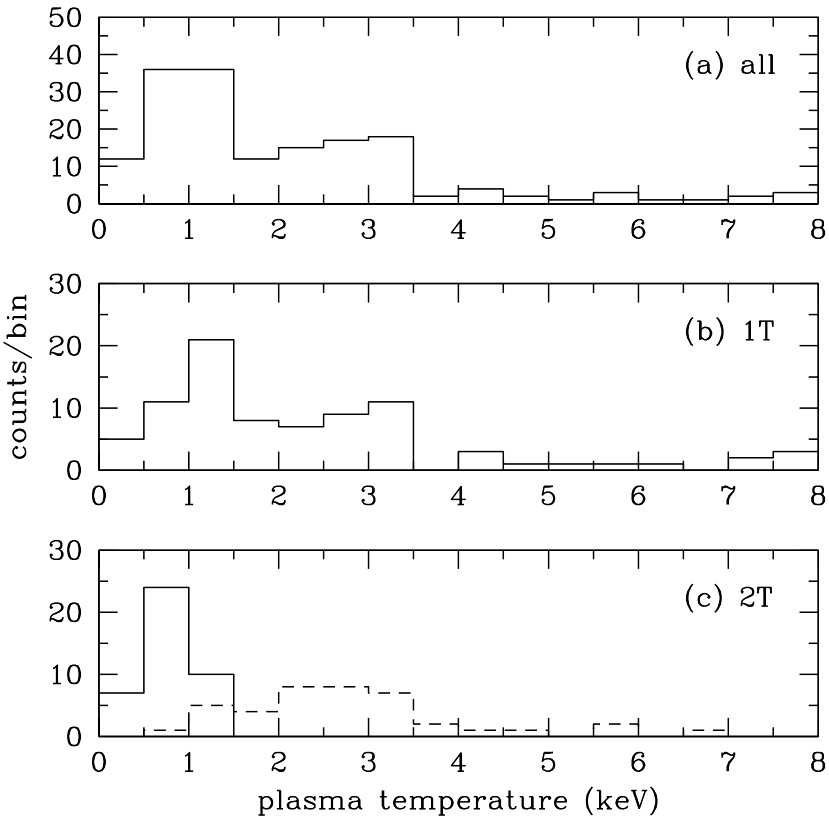


Figure 9.4: Histograms of plasma temperatures of bright (S) NIR-IDed X-ray sources; (*a*) total, (*b*) one-temperature plasma, and (*c*) two-temperature plasma sources. In (*c*), the histograms of the lower and higher temperature are given with solid and dashed, respectively.

with the spectrum of sufficient statistics to examine the two-temperature nature), 23 sources require two-temperature plasma models (Fig. 7.7). The rate of 77% is much higher than the binary rate (∼15%) of this region (Padgett, Strom, & Ghez 1997[146]). Finally, we have no reason to expect that the temperature of each binary component is always the combination of $k_{\rm B}T \sim 1$ keV and 2–3 keV.

One-temperature sources are considered to be with no or negligible contribution by either component. The two peak profile of their temperature histogram (Fig. 9.4 b) is consistent with this idea.

The two components show different temporal behavior. In Table 7.5, we see that the EM values of the high temperature component increase during flare phases, while those of the low temperature component stay constant or slightly decrease. This indicates that the flare-like variability is caused by the high temperature component. This is further strengthened by the fact that almost all the one-temperature plasma sources with flare-like variations (Figs. 7.2–7.6) have the plasma temperature of 2–3 keV or higher; e.g., I67 (2.7 keV), I131 (3.2 keV), I138 (7.5 keV), I165 (4.8 keV), I218 (3.1 keV), I264 (5.5 keV), and I280 (2.7 keV). In contrast, the light curves of the low temperature component show moderate variability of no flare-like episodes. Typical examples are the light curves of I54 (0.8 keV) and I282 (1.0 keV).

9.2.2 Origins of Two-temperature Plasma

A similar bimodal structure in the plasma temperature distribution is seen in the sun. Peres et al. (2000)[150] integrated all the X-ray emissions from the sun using the Soft X-ray Telescope (SXT) on *Yohkoh* and convolved the spectrum with the response function of *ASCA*/SIS in order to facilitate direct comparison with the X-ray emissions from other stars. The synthesized solar spectrum is well fitted with one- or two-temperature plasma model at the solar minimum and maximum. At the solar maximum, the plasma temperatures are $k_{\rm B}T \sim 0.2$ keV and ~ 0.5 keV. From the geometrically well-defined data on solar X-ray emissions, the higher and lower temperature components are found to originate from the solar coronae and flares. We interpret the two-temperature plasma of YSOs in the same analog; i.e., the higher temperature (2–3 keV) component is from flares and the lower temperature component is from the coronae.

Two pieces of evidence support our idea. First, similar bimodal temperature structures are also seen in other stars, with increasing plasma temperature toward younger samples (Figs. 2.7 and 9.5). Figure 9.6 shows the evolution of the representative plasma temperatures separately for the higher (*filled squares*) and lower (and middle) temperature component (*open squares*). Three G-type main sequence sources (EK Dra, HN Peg, and κ^1 Cet; Güdel, Guinan, & Skinner 1997[74]) and one G-type pre-main-sequence source (SU Aur; Skinner & Walter 1998[173]) are plotted. As the stellar rotation becomes slower as increasing ages, main sequence sources evolve magnetically inactive and the plasma temperature decreases (Güdel et al. 1997[74]). Our sample sources are at the age of ~1 Myr (or younger), which settles both of their typical high and low temperatures on this temperature-age relations.

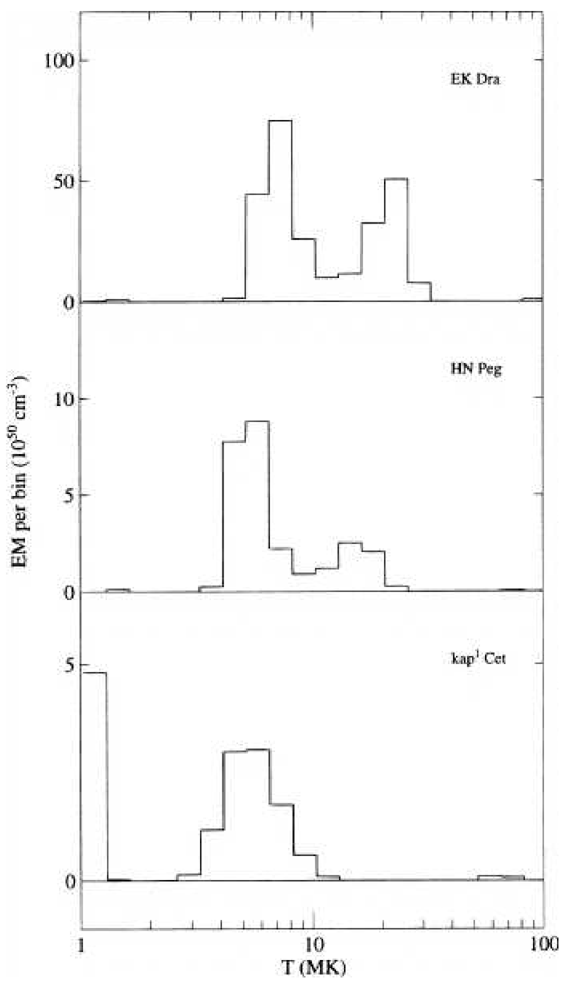


Figure 9.5: Differential emission measure distributions of coronally active main sequence stars (Güdel et al. 1997[74]); EK Dra (*top*), HN Peg (*middle*), and κ^1 Cet (*bottom*).

Second, the time-sliced spectroscopy on our bright X-ray sources (Sect. 7.2) indicates

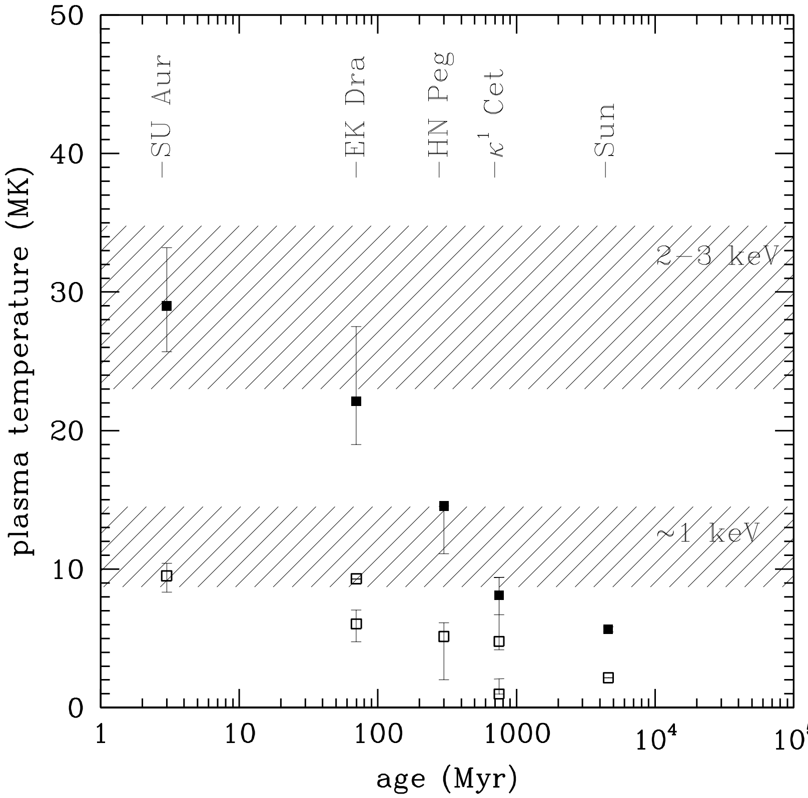


Figure 9.6: Evolution of the plasma temperatures of the soft (and middle) component (*open squares*) and the hard component (*filled squares*) of some main and pre-main-sequence stars (sun; Peres et al. 2000[150], EK Dra, HN Peg, and κ^1 Cet; Güdel, Guinan, & Skinner 1997[74], and SU Aur; Skinner & Walter 1998[173]).

that flaring activities are attributable to the high temperature component. This is the same with the solar plasma and main sequence stars (Güdel et al. 1997[74]; Güdel et al. 2002[75]).

The high-temperature plasma component seen in the spectra at quiescent phases may be the integration of small and temporally-unresolved flares. Güdel et al. (1997)[74] and Güdel et al. (2002)[75] further discussed, based on this interpretation, that the the EM ratio of both components can be a function of flare frequency. Future long-exposure observations on YSOs as well as main sequence stars will give an insight on this issue.

9.2.3 X-ray Properties among Evolutional Classes

Based on the evolutional class estimates of NIR-IDed X-ray sources using the $(J–H)/(H–K)$ color-color diagram as well as the UV excess and H_α emission data, we separated them into class I (protostars), class II (cTTSs), and class III (wTTSs) objects (Sect. 6.2). The statistics of each class are summarized in Table 9.4 with the results of X-ray temporal and spectral analyses, while the averaged X-ray properties are in Table 9.5.

Table 9.4: Number of sources among classes

	class I	class II	class III	others[a]	sum
X-ray sources...	13	59	170	36	278
bright (T)......	4	31	71	14	120
variable........	3	18	38	7	66
	(75%)	(58%)	(54%)	(50%)	(55%)
bright (S)......	7	37	81	17	142
one-temperature	6	25	43	13	87
	(86%)	(68%)	(53%)	(76%)	(61%)
two-temperature	0	10	29	2	41
	(0%)	(27%)	(36%)	(12%)	(29%)

[a] NIR-IDed but not classified either into class I, II, or III based on their NIR colors.

Despite that we do not have class I objects of two-temperature plasma, we can regard that the $< kT^{(1)} >$ value represents $< kT^{(2)}_{high} >$ in this class. This is because class I objects, even if they have two-temperature emissions, appear only with higher temperature plasma due to the heavier extinction at the soft X-ray band. A notable fact is that the

Table 9.5: Comparison of X-ray properties among classes

	class I	class II	class III
$< \log N_{\rm H}$ (cm^{-2}) $>$	22.6 (±0.3)	21.8 (±0.7)	21.5 (±0.6)
$< kT^{(1)}$ (keV) $>$	3.71 (±1.9)	2.82 (±1.9)	2.00 (±1.7)
$< kT^{(1)}$ (keV) $>_{\rm w}$[a]	3.30 (±1.1)	1.85 (±1.0)	1.32 (±0.4)
$< \log L_{\rm X}^{(1)}$ (ergs s^{-1}) $>$ . . .	30.6 (±0.5)	30.3 (±0.4)	30.0 (±0.5)
$< kT_{high}^{(2)}$ (keV) $>$	N/A	3.60 (±2.1)	2.64 (±1.2)
$< kT_{high}^{(2)}$ (keV) $>_{\rm w}$[a]	N/A	2.33 (±0.5)	2.30 (±0.8)
$< \log L_{X\ high}^{(2)}$ (ergs s^{-1}) $>$	N/A	30.4 (±0.3)	30.2 (±0.6)
$< kT_{low}^{(2)}$ (keV) $>$	N/A	0.83 (±0.1)	0.71 (±0.3)
$< kT_{low}^{(2)}$ (keV) $>_{\rm w}$[a]	N/A	0.85 (±0.1)	0.88 (±0.2)
$< \log L_{X\ low}^{(2)}$ (ergs s^{-1}) $>$	N/A	30.2 (±0.3)	30.1 (±0.6)

[a] Weighted means, which were calculated by weighting the values with the inverse square of their uncertainty.

$< kT_{high}^{(2)} >$ decreases along the evolution. A similar trend is seen in $< kT^{(1)} >$ and $< kT_{low}^{(2)} >$, although we have to pay attention to the fact that X-ray detections from younger and more obscured sources are biased for the harder emissions. This trend of decreasing higher plasma temperature can be understood by extrapolating the relation in Figure 9.6 toward younger ages than 1 Myr.

9.3 X-ray Emissions from NIR-unIDed Sources

Ten NIR-unIDed X-ray sources (I128a, I132, I135, I175, I186, I196, I203, I241, I247, and I363) in the 1.3 mm integral-shaped ridge appear to be separated into two groups. The first group consists of four sources (I128a, I175, I241, and I247), which are associated with jet and outflow systems. The remaining sources comprise the second group, which has no such features. We try to interpret the X-ray emissions from these two groups in the context of jet-induced plasma and magnetic activities of deeply embedded NIR invisible YSOs.

9.3.1 Jet-induced Plasma Emissions

Sources in the first group share many characteristics in common (Table 9.6). They are (1) located at the 1.3 mm cloud cores, and (2) are associated with H_2 outflows seen in the QUIRC H_2-band image (Fig. 9.7), (3) CO and HCO^+ outflows (Aso et al. 2000[11]), and (4) the centimeter emissions (Reipurth et al. 1999[162]). VLA 7 is extended in the direction of molecular outflows, indicating that this is a free-free emission from the H_{II} region ionized by protostar jets. We resolved VLA 1 into two sources (VLA 1a and VLA 1b) and discussed that these sources have free-free emissions in origin as well (Sect. 8.4).

Their close-up view in the QUIRC K-band image (Fig. 9.7) reveals more of their common features. (5) The 1.3 mm cores that contain these X-ray sources also possess a NIR source that is classified either into class I protostars or cTTSs. For TKK J05351833−0500329 at MMS 2, we resolved it into two NIR sources (IRS 3 and IRS 5) with our Subaru and IRTF observations, and classified them into class I protostars using their J-, H-, K-, and L'-band colors (Sect. 8.3). TKK J05352333−0507096 at FIR 1c has the color $(J–H)_{\rm CIT}$ = 2.83 and $(H–K)_{\rm CIT} = 2.19$, showing its class I nature (Fig. 6.2). 2MASSI J0535315−050547 at MMS 10, although it shows no NIR excess in the $(J–H)/(H–K)$ diagram (Fig. 6.2), have a large extinction with $(J–H)_{\rm CIT} = 2.00$ and $(H–K)_{\rm CIT} = 1.42$. Considering its large extinction and the insensitivity of the K band to the NIR excess emissions, this source can also be a class I protostar. (6) These NIR sources show apparent reflection features in the K-band image. Finally, (7) the X-ray sources are close to but significantly offset from these NIR protostars in the direction of jet and outflow system.

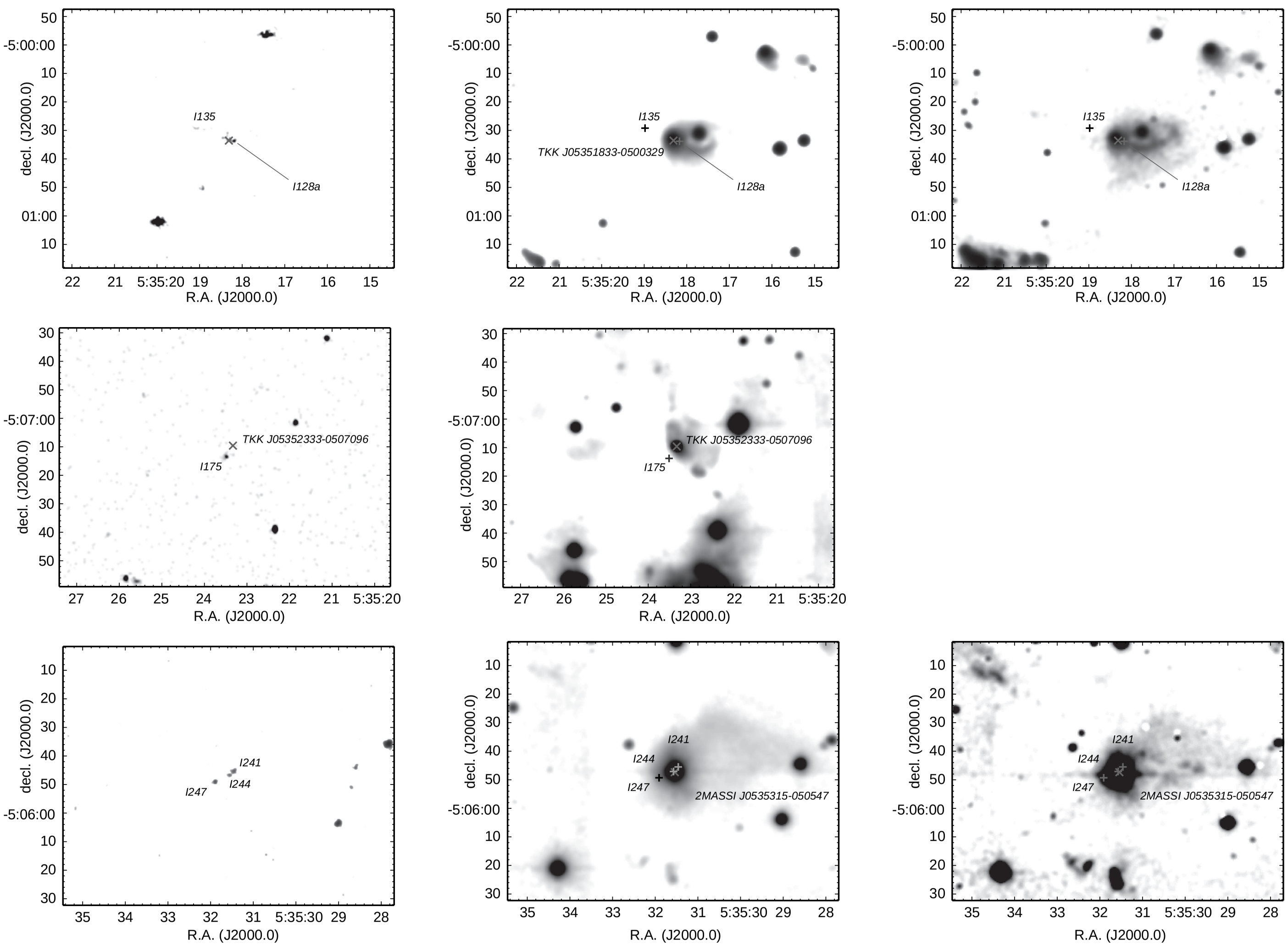


Figure 9.7: X-ray and NIR images of I128a at MMS 2 (*top*), I175 at FIR 1c (*middle*), and I241 and I247 at MMS 10 (*bottom*). The panels in the left column are the X-ray images in the hard band, while those in the middle column are the *K*-band images. The H_2 intensity images are in the right column. I175 was out of the FOV of our H_2 observations. The positions of X-ray and NIR sources are marked with pluses and crosses, respectively.

Table 9.6: NIR-unIDed X-ray sources associated with jet and outflow systems

ID	X-ray[a] counts	1.3 mm core[b]	jet/outflow associations H_2[c]	CO[d]	$H^{13}CO^+$[d]	3.6 cm[e]	class I protostar[f]
I128a	66	MMS 2	yes	yes	yes	VLA 1	TKK J05351833−0500329
I175	33	FIR 1c	N/A	yes	yes	VLA 7	TKK J05352333−0507096
I241	22	MMS 10	yes	yes	yes	......	2MASSI J0535315−050547
I247	33	MMS 10	yes	yes	yes	......	2MASSI J0535315−050547

[a] Values in the 0.5–8.0 keV range.

[b] Chini et al. (1997)[35].

[c] The H_2 emissions found in QUIRC images (Fig. 8.8). I175 was out of the FOV of our observation.

[d] Aso et al. (2000)[11].

[e] Reipurth et al. (1999)[162]. VLA 7 has an extended structure. VLA 1 was resolved into VLA 1a and VLA 1b in our higher resolution VLA observation (Sect. 8.4), in which we found VLA 1a is also extended in the direction of the molecular outflow as well as VLA 7.

[f] The class I protostars located in the same 1.3 mm cloud.TKK J05351833−0500329 was resolved into two class I protostars (IRS 3 and IRS 5) in our higher resolution Subaru observation (Sect. 8.3). I244 is the X-ray counterpart of 2MASSI J0535315−050547.

We pick up I128a as a test case. We revealed the vicinity of this source with our follow-up studies with high resolution NIR and centimeter imaging observations as follows (Figs. 8.13 and 9.8). Two NIR sources of class I protostar nature (IRS 3 and IRS 5) are located at the center of 1.3 mm cloud core (MMS 2). These two protostars accompany centimeter emissions (VLA 1a and VLA 1b) that originate from the the H_{II} region ionized by the protostellar jets from the two NIR sources. A global outflow is seen in the H_2 band, the direction of which aligns with the extended structure of VLA 1a (Sect. 8.3 and 8.4). A hard X-ray source (I128a) is located at the origin of this outflow and has a significant offset from these two protostars. The X-ray characteristics that we obtained on I128 can be applied to I128a, because I128a occupies most (~70%) of the X-ray photons of this complex (Fig. 8.9). It has a thin-thermal plasma spectrum of $k_BT = 3.05$ keV and $EM = 2.1\times10^{53}$ cm^{-3} with the absorption of $N_H = 1.27\times10^{23}$ cm^{-2} (Sect. 7.2).

We propose an interpretation to explain these hard X-ray emissions together with the centimeter emissions based on the shock-induced plasma scenario. Figure 9.9 shows the schematic view, where the jet from a protostar collides into a dense obstacle and produces the shock front. The hard X-ray is emitted from the post shock (PS) region, while the centimeter emission is from the recombination zone (RZ) behind the shock. The RZ is maintained by the continuous ionization by UV photons from the PS region.

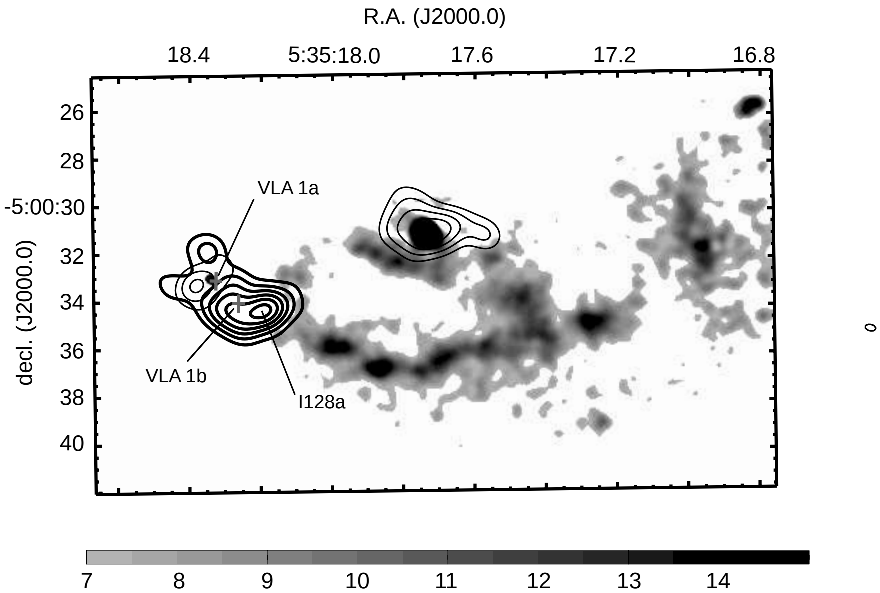


Figure 9.8: Multi-wavelength view on MMS 2. The H_2 intensity is in gray scale, the hard and soft X-rays are with the thin and the thick contours, and the position of centimeter sources (VLA 1a and VLA 1b) are marked with pluses. I128a, the hard X-ray peak, is significantly offset from two centimeter emissions in the direction of global H_2 outflow. See also Figure 8.13 for the positions of NIR sources.

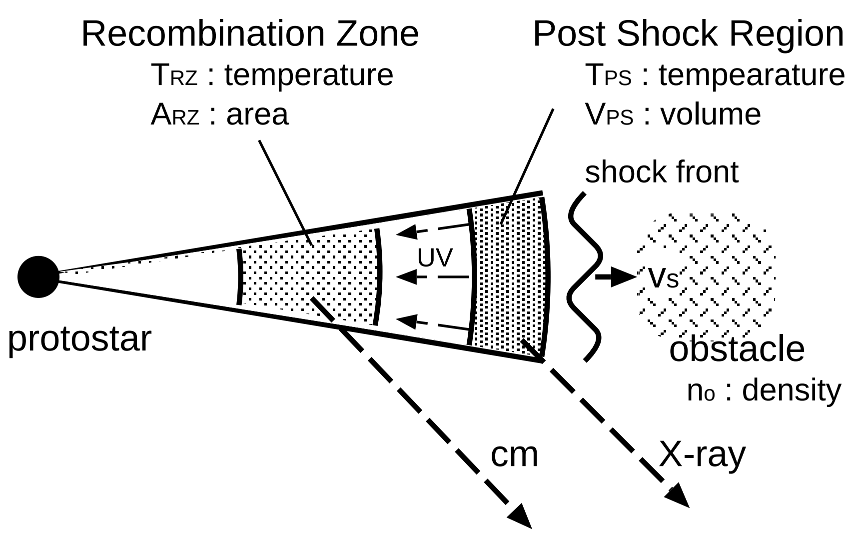


Figure 9.9: Schematic view of the X-ray and centimeter emissions of protostellar jet origin.

In PS, the temperature ($T_{\rm PS}$) and the density ($n_{\rm PS}$) are expressed as (Raga et al. 2002[158])

$$T_{\rm PS} = 1.5 \times 10^5 \left(\frac{v_{\rm s}}{100\ {\rm km\ s^{-1}}}\right)^2\ [{\rm K}], \tag{9.2}$$

$$n_{\rm PS} = 4n_0\ [{\rm cm^{-3}}], \tag{9.3}$$

where $v_{\rm s}$ and n_0 are the velocity and density of the shock produced by the collision of the jet upon the obstacle. Assuming that light elements are fully ionized in PS, the emission measure (EM) is given with the electron density ($n_{\rm PS}$) and the volume ($V_{\rm PS}$) as

$$EM = n_{\rm PS}^2 V_{\rm PS}\ [{\rm cm^{-3}}]. \tag{9.4}$$

By substituting the observed values ($T_{\rm PS}$ = 35 MK and $EM = 2.1\times10^{53}$ cm^{-3}), we obtained $v_{\rm s} = 1.5\times10^3$ km s^{-1} and $n_0 = 5.3\times10^2$ cm^{-3}. Here, we assumed that PS is a cube with the length of $0.5''$ (=the scale of an ACIS-I pixel).

The values of $v_{\rm s}$ and n_0 are consistent with what can be independently derived from the centimeter observations. In RZ, the centimeter intensity (S_ν) is given by

$$S_\nu = \frac{A_{\rm RZ}}{D^2} 2k_B T_{\rm RZ} \left(\frac{\nu}{c}\right)^2 \tau_\nu \tag{9.5}$$

at the optically-thin limit and with the Rayleigh-Jeans approximation. Here, D is the distance to the source, and $T_{\rm RZ}$ and $A_{\rm RZ}$ are the temperature and the surface area of RZ. Curiel et al. (1989)[43] showed that the optical depth (τ_ν) is expressed in terms of the shock parameters by

$$\tau_\nu = 1.55 \times 10^{-7} \left(\frac{n_0}{1\ {\rm cm^{-3}}}\right) \left(\frac{v_{\rm s}}{100\ {\rm km\ s^{-1}}}\right)^{1.68} \left(\frac{T_{\rm RZ}}{10^4\ {\rm K}}\right)^{-0.55} \left(\frac{\nu}{5\ {\rm GHz}}\right)^{-2.1}. \tag{9.6}$$

By substituting the observed values of VLA 1a (ν = 8.3 GHz, S_ν = 0.128 mJy, and D = 450 pc) as a typical value, we obtained

$$\left(\frac{n_0}{1\ {\rm cm^{-3}}}\right) \left(\frac{v_{\rm s}}{100\ {\rm km\ s^{-1}}}\right)^{1.68} = 1.7 \times 10^5. \tag{9.7}$$

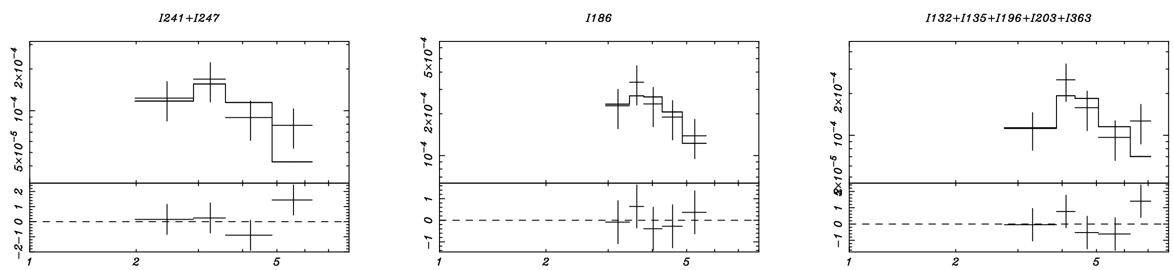
Here, we assumed that $T_{\rm RZ} = 10^4$ K and $A_{\rm RZ}$ is an ellipse of the observed major and minor axis lengths of VLA 1a.

From the X-ray observation, we can independently derive that

$$\left(\frac{n_0}{1\ {\rm cm^{-3}}}\right) \left(\frac{v_{\rm s}}{100\ {\rm km\ s^{-1}}}\right)^{1.68} = 5.2 \times 10^4. \tag{9.8}$$

This is in good agreement with the value obtained with the centimeter data in a factor of a few, supporting the protostellar jet scenario for the origin of the X-ray and centimeter emissions.

Other three X-ray emissions with the same multi-wavelength features can be understood with the same scheme. The positions of X-ray and NIR sources are similar to the case of I128a, where X-ray emissions are $1''$–$6''$ offset from the protostar in the direction of jet and outflow. More interestingly, I241 and I247 are positioned at the opposite side of the central NIR source (Fig. 9.7). We may be seeing a jet and counter-jet pair of this system. The combined spectrum of I241 and I247 (Fig. 9.10 a; Table 9.7) supports this idea, which are well fitted by a thin-thermal plasma model with similar best-fit values ($k_{\rm B}T$ = 2.8 keV, EM = 9.2$\times 10^{22}$, and $N_{\rm H}$ = 5.8$\times 10^{22}$ cm^{-2}) with I128a. Likewise, I135 at MMS 2 may be the counter-jet of I128a (Fig. 9.7).



(a) (b) (c)

Figure 9.10: Spectra and the best-fit models of one-temperature thin-thermal plasma fittings of (a) I241+I247, (b) I186, and (c) I132+I135+I196+I203+I363. The metallicity of all elements is fixed to be 0.3 solar. In the upper panels, the data (*pluses*) and the best-fit model (*solid steps*) are plotted over the energy (keV; *horizontal axis*) versus normalized spectral intensity (count rate keV^{-1}; *vertical axis*) plane. The response functions of the optics and the detector are convolved into the model. In the lower panels, the residuals between the background-subtracted data and the best-fit model are plotted over the energy (keV; *horizontal axis*) versus $\chi^{(i)} = (E^{(i)}_{data} - E^{(i)}_{model\otimes ARF\otimes RMF})/\Delta E^{(i)}_{data}$ (*vertical axis*) plane.

9.3.2 Magnetic Activities of Deeply Embedded YSOs

Other X-ray sources (I132, I135, I186, I196, I203, and I363) that are associated with 1.3 mm emissions but not with jet and outflow systems can be deeply embedded YSOs such as class 0 objects for the following reasons.

Table 9.7: One-temperature plasma fittings of NIR-unIDed *Chandra* sources

ID	counts[a]	$N_{\rm H}$[b] (10^{22} cm^{-2})	$k_{\rm B}T$[b] (keV)	$L_{\rm X}$[a] (ergs s^{-1})
I241+I247	55	5.83 (0.0–6.0)	2.80 (1.9–...)	1.00e+30
I186	66	14.2 (0.0–41)	1.91 (0.5–...)	2.07e+30
I132+I135+I196+I203+I363	66	20.2 (8.0–42)	3.71 (1.0–...)	3.47e+30

[a] Values in the 0.5–8.0 keV range. For combined fittings, the sum of all sources are shown for the count and luminosity.

[b] The lower and upper limit (1 σ) are given in parentheses. Some have too few spectral bins to derive the uncertainty of their best-fit parameters.

First, it has no NIR counterpart brighter than the QUIRC limiting magnitude of $K \sim 16.5$ mag. I132 at MMS 3, which was observed with the Subaru deep imaging observation, has much tighter upper limit of $K \sim 19.6$ mag. Low-mass class I protostars are generally detected at 10–14 mag and class II at 8–12 mag in the K band at the distance of 450 pc (for example, see Fig. 24. in Aspin, Sandell, & Russel 1994[12]), hence are easily detected with our K-band sensitivity. In fact, all class I and class II sources with X-ray emissions (Table A.1) have the K-band magnitude of 6.5–15.1 mag and 9.3–15.3 mag, respectively. These six X-ray sources are fainter than typical class I protostars by more than 10–100 times in the K band, which indicates that these sources are either much more obscured than typical class I protostars, or are intrinsically faint in NIR because of their much lower mass than low-mass sources.

Our results of the X-ray spectral analysis on these sources favor the high obscuration interpretation. The spectrum of the brightest source (I186) and the combined spectrum of the rest of faint others (I132, I135, I196, I203, and I363) are well fitted with a thin-thermal plasma model (Fig. 9.10) with the best-fit value of $L_{\rm X} = 0.5$–2.0×10^{30} ergs s^{-1} (Table 9.7). Here, we assumed that the $L_{\rm X}$ values are the same for I132, I135, I196, I203, and I363, because they have the same order of X-ray counts (Table A.1). If these sources follow the relation between the mass and $L_{\rm X}$ (Fig. 9.3), this X-ray luminosity corresponds to the source of 0.1–0.5 $M_\odot$. If these X-ray sources are in this mass range with moderate extinction, they should be easily detected at the detection limit of our deep NIR observations. This rules out the possibility of these sources just to have a very low mass.

Second, they have much larger $N_{\rm H}$ values of 1–2×10^{23} cm^{-2} than typical class I sources

(Table 9.5). This is converted to $A_V \approx 100$ mag using the equation (7.6).

Third, some of the sources spatially coincide with the 350 μm cores (Lis et al. 1998[117]), which provides the position of protostellar cores with a better spatial resolution than 1.3 mm mapping observations by Chini et al. (1997)[35]. I132 and I186 are located within the positional uncertainty of $\sim 5''$ respectively from CSO 7 and CSO 19, and I363 is within $9''$ from CSO 15. I132 is also associated with MMS 3, a class 0 candidate source, detected by 1.3 mm mapping observations (Chini et al. 1997[35]).

These sources, if they are class 0 sources, would have magnetic activities in the same way as class I protostars. The best-fit values of their $k_{\rm B}T$ (2–3 keV) and $L_{\rm X}$ (0.5–2.0×10^{30} ergs s^{-1}) is in the typical range of that of class I sources, which supports this idea. We may see the higher temperature component of these sources.

We note that no X-ray emissions from bona-fide class 0 sources were found in other star forming regions (e.g., NGC 2068) that were observed by *Chandra* with better sensitivity (E. D. Feigelson 2002, private communication). However, our X-ray observations are as deep as, or deeper than any other wavelength observations with unprecedented spatial resolution in this field. Just as an infrared source with no optical counterpart found by Becklin & Neugebauer (1967)[19] turned out to be a new class of YSOs in the Orion nebula, our X-ray sources can be of this kind. Further studies are mandatory to identify the nature of these sources, such as millimeter–sub-millimeter interferometer observations to detect circumstellar clumps of these prospective YSOs.

Chapter 10

Conclusions

Using OMC-2 and OMC-3 as our study field and taking a multi-wavelength observational approach, we discussed the origins and mechanisms of ∼400 X-ray sources to understand the wide variety of X-ray-emitting phenomena seen in star-forming regions.

1. We conducted deep X-ray and NIR observations respectively using ACIS on *Chandra* and QUIRC on the University of Hawaii 88 inch (2.2 m) telescope. The X-ray observation is complete down to $F_{\mathrm{X}} \sim 10^{-14.5}$ ergs s^{-1} cm^{-2} with the faintest detected source of $F_{\mathrm{X}} \sim 10^{-15.5}$ ergs s^{-1} cm^{-2} in the 0.5–8.0 keV energy range. The NIR image has the 90% completeness limit of $J \sim 17.5$, $H \sim 16.5$, and $K \sim 16.0$ mag, matching well with the *Chandra* limit (Sects. 5.1 and 5.2).

2. We extracted 385 X-ray sources in the 17×17 arcmin^2 ACIS-I FOV and 1448 NIR sources in the 512 arcmin^2 QUIRC FOV. Combining the 2MASS catalog with our QUIRC source list and correlating them with our *Chandra* data, we identified the NIR counterpart for 278 (∼72%) X-ray sources (Sects. 5.1 and 5.2). Most of NIR-IDed X-ray sources are YSOs that belong to OMC-2 and OMC-3, considering their K-band flux and luminosity function (Sect. 6.2).

3. The NIR-IDed X-ray sources were examined for their J-, H-, and K-band colors to estimate their mass, bolometric luminosity, and evolutional class (Sect. 6.2). Their X-ray temporal and spectral features were also analyzed to derive their flux variability, plasma temperature, emission measure, and X-ray luminosity (Sects. 7.1 and 7.2).

4. The averaged X-ray properties among different mass ranges were compared. We found

that IM ($2.0\ M_\odot \leq M < 10.0\ M_\odot$), LM ($0.2\ M_\odot \leq M < 2.0\ M_\odot$), and VLM ($M < 0.2\ M_\odot$) YSOs have the same X-ray emission mechanisms based on their $L_{\rm bol}/L_{\rm X}$ values, $L_{\rm X}$-mass relation, and averaged X-ray features. This is in contrast with the HM ($M \geq 10.0\ M_\odot$) sources that emit X-rays of the stellar wind origin (Sect. 9.1).

5. We revealed that the X-ray emission from IM, LM, and VLM sources consists of two components of different temperatures with $k_{\rm B}T \sim 1.0$ keV and 2.0–3.0 keV. Based on the time-sliced X-ray spectroscopy as well as comparison with the sun and other sources in the literature, we proposed that the soft component is from coronal and the hard component is from flare activities (Sect. 9.2).

6. Most of the NIR-unIDed X-ray sources are AGNs based on their X-ray spectra, HR distribution, and the number counts. However, the spatial distribution of these sources has an excess along the ridge of 1.3 mm cloud cores, indicating that ~10 of them are related to star formation (Sect. 8.1).

7. We conducted follow-up observations to identify the nature of the NIR-unIDed X-ray sources associated with the 1.3 mm ridge using QUIRC on the University of Hawaii 88 inch (2.2 m) telescope in the H_2 $v = 1-0$ S(1) band (Sect. 8.2), IRCS on the Subaru telescope and NSFCam on IRTF in the J, H, K, L', and H_2 $v = 1-0$ S(1) bands (Sect. 8.3) in addition to the centimeter interferometer imaging observation using VLA (Sect. 8.4).

8. Four NIR-unIDed X-ray sources along the 1.3 mm ridge have many multi-wavelength features in common; association with the 1.3 mm cores, H_2, CO, HCO^+ outflows, and centimeter emissions from radio jets. These X-ray sources are very close to but significantly offset from the class I NIR sources in the direction of the jet and outflow system. We proposed a picture of the jet-induced plasma to account for these X-ray emissions (Sect. 9.3).

9. We have six NIR-unIDed X-ray sources that are located along the 1.3 mm ridge but not with the jet and outflow associations. Theses sources can be deeply embedded X-ray-emitting YSOs such as class 0 objects, based on their X-ray luminosity, temperature, NIR flux upper limit, and association of some with 350 μm and 1.3 mm cores (Sect. 9.3).

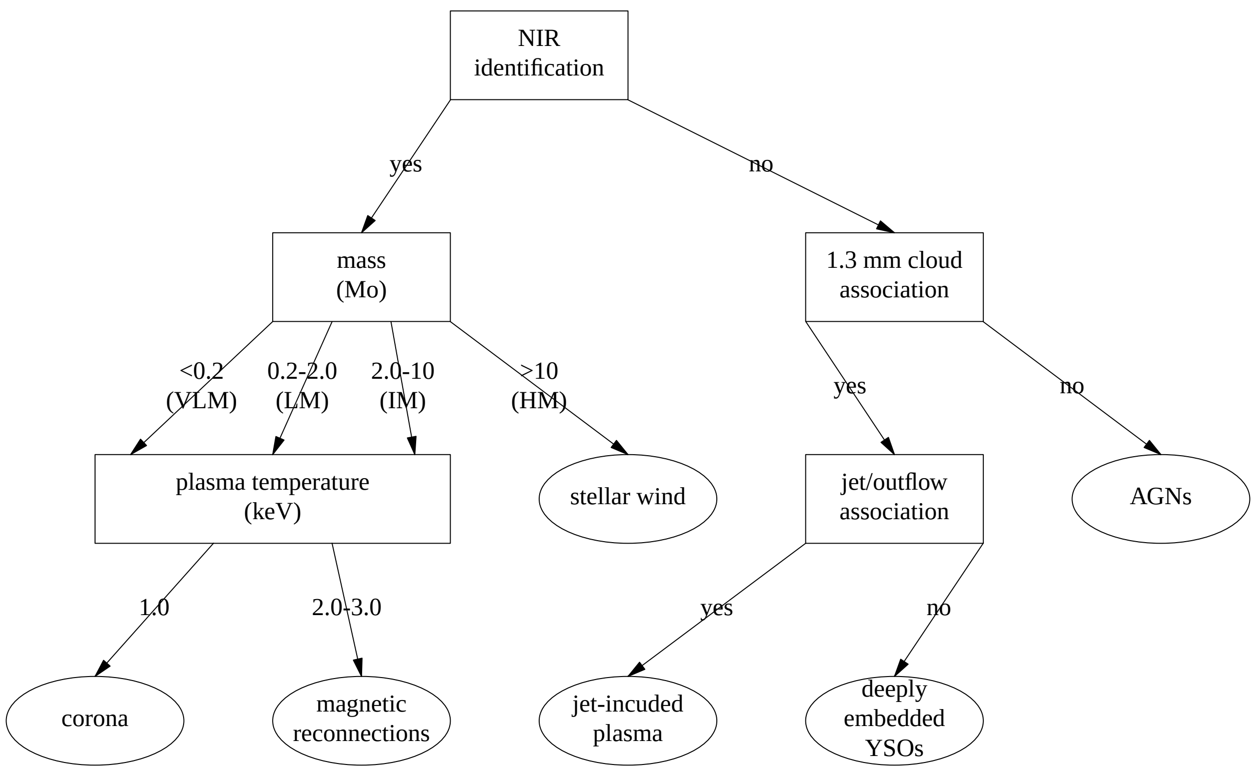


Figure 10.1: Summary

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Appendix A

Chandra Sources and NIR Counterpart

Table A.1: *Chandra* source list

ID[a]	R.A.[b] (hh:mm:ss.ss)	decl.[b] (dd:mm:ss.s)	Photon Counts	HR[c]	NIR counterpart[d]	J^e (mag)	$J-H^e$ (mag)	$H-K^e$ (mag)	class[f]	mass[gh] ($M_\odot$)
I1	05:34:38.60	−05:08:42.8	255	0.31	TKK J05343791−0508480	16.22	0.94	0.25	others	0.040–0.050
I2	05:34:40.86	−05:06:58.2	63	−0.02	TKK J05344064−0506586	15.50	2.33	1.28	III	0.500–0.570
I3	05:34:43.10	−05:09:18.0	368	0.62						
I4	05:34:44.89	−05:06:49.1	403	−0.54	2MASSI J0534449−050649	11.88	0.72	0.38	III	0.350–0.400
I5	05:34:45.19	−05:10:45.1	176	−0.25	2MASSI J0534451−051047	12.17	0.75	0.52	others	0.300–0.350
I6	05:34:48.29	−05:07:13.1	535	−0.66	2MASSI J0534482−050713	12.99	0.66	0.22	III	0.150–0.175
I7	05:34:48.37	−05:05:01.1	549	−0.77	2MASSI J0534484−050501	12.48	0.68	0.19	III	0.200–0.250
I8	05:34:49.18	−05:04:38.2	127	−0.97	2MASSI J0534492−050438	12.55	0.69	0.23	III	0.200–0.250
I9	05:34:50.33	−05:06:38.1	70	−0.29	TKK J05345056−0506383	13.96	0.51	0.25	III	0.080–0.090
I10	05:34:50.43	−05:11:11.2	773	−0.75	2MASSI J0534505−051110	11.96	0.66	0.17	III	0.300–0.350
I11	05:34:50.59	−05:06:49.8	110	0.40						
I12	05:34:50.64	−05:04:07.6	361	−0.70	2MASSI J0534506−050407	12.65	0.72	0.29	III	0.200–0.250
I13	05:34:53.03	−05:03:26.9	1747	−0.73	2MASSI J0534530−050327	11.95	0.75	0.45	II	0.350–0.400
I14	05:34:53.42	−05:10:28.2	579	−0.70	2MASSI J0534534−051027	12.22	0.63	0.23	III	0.250–0.300
I15	05:34:53.45	−05:01:30.2	49	0.10	2MASSI J0534536−050129	14.34	0.62	0.21	III	0.075–0.080
I16	05:34:53.71	−05:05:48.9	145	−0.75	2MASSI J0534537−050548	13.33	0.64	0.25	III	0.130–0.150
I17	05:34:55.38	−05:01:39.3	534	−0.83	2MASSI J0534554−050139	12.51	0.73	0.15	III	0.250–0.300
I18	05:34:55.66	−05:08:44.3	39	0.38						
I19	05:34:55.76	−05:07:42.5	103	0.24						
I20	05:34:56.07	−05:00:55.6	359	−0.69	2MASSI J0534561−050055	12.80	0.64	0.21	III	0.150–0.175
I21	05:34:56.13	−05:06:01.5	276	−0.82	2MASSI J0534561−050601	12.56	0.61	0.24	III	0.175–0.200
I22	05:34:56.64	−05:06:26.1	55	0.20						
I23	05:34:56.67	−05:04:37.9	98	0.47						
I24	05:34:56.71	−05:10:43.6	48	−0.33	2MASSI J0534567−051043	13.46	0.73	0.30	III	0.130–0.150
I25	05:34:56.82	−05:11:33.5	2286	−0.82	2MASSI J0534568−051132	10.89	0.71	0.47	others	0.900–0.950
I26	05:34:57.35	−05:06:04.0	19	0.58						
I27	05:34:58.19	−05:09:27.9	558	−0.71	2MASSI J0534581−050927	12.29	0.76	0.28	III	0.300–0.350
I28	05:34:58.21	−05:11:54.6	258	−0.64	2MASSI J0534581−051153	12.65	0.60	0.22	III	0.175–0.200
I29	05:34:58.52	−05:12:27.0	135	0.04	2MASSI J0534586−051226	14.03	0.55	0.24	III	0.080–0.090
I30	05:34:59.22	−05:01:26.2	28	0.43						
I31	05:34:59.32	−05:05:30.0	740	−0.75	2MASSI J0534593−050530	12.41	0.77	0.31	III	0.250–0.300
I32	05:34:59.47	−05:08:25.6	30	0.27						
I33	05:35:00.40	−05:09:44.8	213	0.13	2MASSI J0535003−050944	13.81	0.57	0.31	III	0.090–0.100
I34	05:35:00.87	−05:09:39.4	135	−0.72	2MASSI J0535008−050938	13.01	0.56	0.32	III	0.130–0.150
I35	05:35:01.37	−05:11:06.1	60	0.33	2MASSI J0535014−051104	13.86	0.96	0.57	II	0.150–0.175

(*cont.*)

ID[a]	R.A.[b]	decl.[b]	Counts	HR[c]	NIR counterpart[d]	J^e	$J-H^e$	$H-K^e$	class[f]	mass[gh]
I36	05:35:01.43	−05:09:32.8	670	−0.85	2MASSI J0535014−050932	12.13	0.67	0.18	III	0.300–0.350
I37	05:35:01.98	−05:00:43.2	63	−0.46	2MASSI J0535019−050043	14.51	0.63	0.22	III	0.072–0.075
I38	05:35:02.16	−05:04:38.5	83	−0.76	2MASSI J0535021−050438	13.88	0.76	0.37	III	0.110–0.130
I39	05:35:02.52	−05:04:10.0	11	0.82						
I40	05:35:02.56	−05:04:58.5	9	1.00						
I41	05:35:02.79	−05:03:03.9	20	−0.40	2MASSI J0535028−050303	15.83	0.58	0.00	others	0.030–0.040
I42	05:35:02.79	−05:00:01.9	1023	−0.57	2MASSI J0535027−050002	12.29	1.17	0.45	III	0.620–0.700
I43	05:35:02.85	−05:02:00.6	20	0.50						
I44	05:35:03.12	−05:09:17.3	90	−0.76	2MASSI J0535031−050917	13.15	0.72	0.34	III	0.150–0.175
I45	05:35:03.30	−05:04:48.1	26	0.69						
I46	05:35:03.33	−05:13:07.2	159	0.17	TKK J05350365−0513104	15.07	1.69	0.74	III	0.200–0.250
I47	05:35:03.41	−05:05:40.4	1771	−0.62	2MASSI J0535034−050540	9.17	0.22	0.03	III	1.400–1.400
I48	05:35:03.81	−05:05:19.1	17	0.53						
I49	05:35:03.99	−05:01:41.1	20	0.70						
I50	05:35:04.23	−05:08:42.5	54	0.26						
I51	05:35:04.30	−05:08:12.9	25402	−0.55	2MASSI J0535042−050812	8.18	0.66	0.15	III	>4.000
I52	05:35:04.45	−05:07:36.5	14	−0.43	TKK J05350440−0507356	14.75	1.10	0.54	III	0.110–0.130
I53	05:35:04.55	−04:58:29.9	109	−0.21	2MASSI J0535046−045829	12.83	1.05	0.74	II	0.350–0.400
I54	05:35:04.64	−05:09:56.1	437	−0.78	2MASSI J0535046−050955	11.38	0.65	0.26	III	0.500–0.570
I55	05:35:05.24	−04:59:23.2	26	0.54						
I56	05:35:05.65	−04:58:53.6	334	−0.83	2MASSI J0535056−045853	11.94	0.65	0.24	III	0.300–0.350
I57	05:35:05.66	−05:04:53.4	7	0.43						
I58	05:35:05.76	−05:11:34.5	92	−0.28	2MASSI J0535057−051135	12.91	1.39	0.66	III	0.620–0.700
I59	05:35:05.95	−05:08:37.5	18	−0.67	2MASSI J0535058−050838	15.17	1.08	0.50	III	0.090–0.100
I60	05:35:06.09	−05:12:17.0	43	−0.30	2MASSI J0535061−051216	8.13	0.00	0.00	III	2.000–2.200
I61	05:35:06.36	−05:09:00.0	10	0.20						
I62	05:35:06.49	−04:58:40.4	58	−0.45	2MASSI J0535063−045841	14.21	0.56	0.30	III	0.075–0.080
I63	05:35:06.58	−05:07:16.0	15	0.60						
I64	05:35:06.61	−05:04:51.2	42	0.43						
I65	05:35:06.71	−05:11:45.4	257	−0.19	2MASSI J0535067−051145	13.85	1.44	0.62	III	0.350–0.400
I66	05:35:06.83	−05:10:39.1	200	−0.76	2MASSI J0535068−051038	12.59	0.67	0.24	III	0.200–0.250
I67	05:35:07.54	−05:11:14.7	1232	−0.21	2MASSI J0535075−051114	12.38	1.07	0.58	III	0.500–0.570
I68	05:35:08.04	−05:01:18.2	14	0.14	TKK J05350779−0501193	16.19	0.93	0.43	III	0.040–0.050
I69	05:35:08.75	−05:04:40.5	26	−0.85	2MASSI J0535087−050440	12.37	0.74	0.36	III	0.250–0.300
I70	05:35:08.96	−05:10:25.6	12	−0.67	2MASSI J0535089−051025	14.56	0.61	0.10	III	0.060–0.070
I71	05:35:09.15	−05:06:47.1	1192	−0.83	2MASSI J0535091−050647	11.24	0.67	0.21	III	0.570–0.600
I72	05:35:09.32	−04:59:32.2	56	0.39						
I73	05:35:09.42	−04:59:41.2	72	−0.56	2MASSI J0535094−045941	12.89	0.68	0.35	III	0.150–0.175
I74	05:35:09.85	−04:58:49.1	442	0.24	TKK J05350988−0458495	16.22	3.04	1.64	III	1.400–1.500

(*cont.*)

ID[a]	R.A.[b]	decl.[b]	Counts	HR[c]	NIR counterpart[d]	J[e]	$J-H$[e]	$H-K$[e]	class[f]	mass[gh]
I75	05:35:10.14	−05:13:40.6	86	0.00						
I76	05:35:10.32	−05:11:11.1	13	0.54						
I77	05:35:10.51	−04:58:45.5	835	0.52	TKK J05351053−0458460	>20.55	>6.77	>3.28	others	. . . .
I78	05:35:11.56	−05:02:08.2	16	0.75						
I79	05:35:11.81	−05:14:00.0	237	−0.09	TKK J05351134−0514016	13.76	0.97	0.39	III	0.150–0.175
I80	05:35:12.72	−05:12:00.8	59	−0.46	2MASSI J0535127−051200	14.11	0.52	0.25	III	0.075–0.080
I81	05:35:12.80	−05:01:46.1	79	−0.95	2MASSI J0535128−050146	13.93	0.59	0.28	III	0.090–0.100
I82	05:35:12.92	−04:55:55.7	325	0.48	TKK J05351310−0455523	12.82	1.88	1.64	I	2.000–2.200
I83	05:35:12.95	−05:02:08.5	109	−0.83	2MASSI J0535129−050208	14.38	0.56	0.33	III	0.070–0.072
I84	05:35:13.00	−05:00:26.2	29	0.38	TKK J05351303−0500261	>20.68	>3.77	2.45	II	
I85	05:35:13.34	−05:09:19.9	1088	−0.74	2MASSI J0535133−050919	12.14	0.68	0.19	III	0.300–0.350
I86	05:35:13.35	−05:04:14.5	7	0.71						
I87	05:35:13.88	−04:58:03.2	207	0.11	2MASSI J0535138−045803	15.63	2.36	1.16	III	0.500–0.570
I88	05:35:13.99	−05:07:38.8	21	0.81						
I89	05:35:14.28	−05:14:26.3	268	0.01	2MASSI J0535142−051427	15.19	2.40	1.50	II	0.900–0.950
I90	05:35:14.29	−05:13:40.3	116	0.00	2MASSI J0535143−051340	16.30	1.94	1.26	II	0.150–0.175
I91	05:35:14.63	−05:06:25.3	34	−0.65	2MASSI J0535146−050625	14.39	1.16	0.65	II	0.130–0.150
I92	05:35:14.64	−05:02:24.7	312	−0.82	TKK J05351464−0502250	12.77	1.02	0.46	III	0.350–0.400
I93	05:35:14.66	−05:03:12.4	574	−0.85	2MASSI J0535146−050312	12.45	0.68	0.18	III	0.200–0.250
I94	05:35:14.67	−05:08:52.3	33	−0.76	2MASSI J0535146−050852	12.83	0.57	0.28	III	0.150–0.175
I95	05:35:14.70	−05:08:26.2	4	1.00						
I96	05:35:14.87	−05:06:48.8	333	0.16	TKK J05351486−0506489	16.75	2.23	1.09	III	0.175–0.200
I97	05:35:14.94	−05:01:18.3	24	0.58						
I98	05:35:14.99	−05:07:13.2	62	0.71						
I99	05:35:15.08	−05:06:53.6	112	0.21	2MASSI J0535150−050653	15.89	2.29	1.20	III	0.350–0.400
I100	05:35:15.18	−05:07:56.4	4	0.50						
I101	05:35:15.27	−05:00:32.6	29	−0.86	2MASSI J0535152−050033	13.42	0.63	0.26	III	0.110–0.130
I102	05:35:15.36	−05:13:40.2	68	−0.53	TKK J05351533−0513382	13.47	0.86	0.12	others	0.150–0.175
I103	05:35:15.49	−05:01:12.2	26	−0.31	2MASSI J0535154−050112	14.34	0.68	0.34	III	0.080–0.090
I104	05:35:15.55	−05:01:43.5	34	−0.53	2MASSI J0535155−050143	13.66	1.09	0.53	III	0.200–0.250
I105	05:35:15.59	−05:09:32.2	21	−0.33	2MASSI J0535156−050931	13.93	0.59	0.28	III	0.090–0.100
I106	05:35:15.61	−04:59:27.3	151	−0.30	2MASSI J0535156−045927	14.63	1.90	1.09	II	0.450–0.500
I107	05:35:15.64	−04:57:12.7	159	−0.47	2MASSI J0535156−045713	13.30	0.77	0.32	III	0.150–0.175
I108	05:35:15.76	−05:12:27.4	72	0.06	TKK J05351580−0512264	18.18	2.99	1.26	others	0.250–0.300
I109	05:35:15.80	−05:03:26.0	177	−0.92	2MASSI J0535158−050326	13.69	0.58	0.34	III	0.100–0.110
I110	05:35:15.93	−05:14:59.5	3013	0.20	TKK J05351596−0514591	11.57	0.44	>0.39	others	0.250–0.300
I111	05:35:16.14	−05:09:19.3	88	−0.91	2MASSI J0535161−050919	12.91	0.62	0.23	II	0.150–0.175
I112	05:35:16.35	−05:04:36.3	17	−0.76	2MASSI J0535163−050436	14.57	0.56	0.33	III	0.060–0.070
I113	05:35:16.40	−04:58:01.9	264	−0.79	2MASSI J0535164−045802	13.22	0.60	0.16	III	0.130–0.150

(*cont.*)

ID[a]	R.A.[b]	decl.[b]	Counts	HR[c]	NIR counterpart[d]	J^e	$J-H^e$	$H-K^e$	class[f]	mass[gh]
I114	05:35:16.49	−05:03:30.2	294	−0.80	2MASSI J0535164−050330	12.23	0.66	0.20	III	0.250–0.300
I115	05:35:16.66	−05:08:54.4	4	0.50						
I116	05:35:16.80	−05:07:27.4	45	−0.91	2MASSI J0535167−050727	14.92	0.49	0.34	others	0.050–0.055
I117	05:35:16.86	−05:07:47.8	16	−0.62	2MASSI J0535168−050747	13.81	0.99	0.54	III	0.150–0.175
I118	05:35:17.12	−05:12:39.1	74	−0.54	TKK J05351715−0512394	13.51	0.62	0.08	III	0.110–0.130
I119	05:35:17.25	−05:03:15.9	10	0.80						
I120	05:35:17.36	−05:12:29.3	65	−0.66	TKK J05351738−0512296	13.70	0.62	0.12	III	0.100–0.110
I121	05:35:17.41	−04:59:56.6	259	0.20	TKK J05351742−0459570	17.26	2.67	1.68	II	0.250–0.300
I122	05:35:17.46	−05:09:49.2	44	−1.00	2MASSI J0535174−050949	13.31	0.58	0.30	II	0.110–0.130
I123	05:35:17.72	−05:00:30.3	29	−0.17	2MASSI J0535177−050031	15.29	2.22	1.37	II	0.500–0.570
I124	05:35:17.92	−05:15:33.2	830	−0.14	TKK J05351792−0515329	11.47	>0.73	....	......	....
I125	05:35:18.24	−05:13:07.1	561	−0.40	TKK J05351824−0513069	11.31	>0.99	....	......	....
I126	05:35:18.24	−05:03:54.1	92	−1.00	2MASSI J0535182−050354	9.29	0.04	0.01	III	0.750–0.800
I127	05:35:18.26	−05:08:05.1	7	−0.14	TKK J05351830−0508048	16.86	2.31	1.28	III	0.175–0.200
I128	05:35:18.31	−05:00:33.0	60	0.93	TKK J05351833−0500329	17.67	3.25	>3.73	I	0.600–0.620
I129	05:35:18.45	−05:08:30.9	15	−0.73	2MASSI J0535184−050830	13.75	0.59	0.23	III	0.090–0.100
I130	05:35:18.61	−04:59:42.2	165	−0.71	2MASSI J0535186−045942	13.78	0.58	0.20	III	0.090–0.100
I131	05:35:18.84	−05:14:46.3	661	−0.48	TKK J05351886−0514456	12.45	0.81	0.34	III	0.300–0.350
I132	05:35:18.93	−05:00:50.1	23	0.65						
I133	05:35:18.94	−05:06:36.3	26	0.85						
I134	05:35:18.96	−05:03:23.1	7	1.00						
I135	05:35:18.98	−05:00:29.2	24	0.83						
I136	05:35:19.65	−05:02:28.6	11	−0.82	2MASSI J0535196−050228	14.47	0.65	0.31	III	0.075–0.080
I137	05:35:19.75	−05:05:31.6	6	0.67						
I138	05:35:19.75	−05:15:34.9	359	0.67	TKK J05351982−0515354	17.41	2.55	2.51	I	0.175–0.200
I139	05:35:19.86	−05:15:08.4	309	0.25	2MASSI J0535198−051508	12.93	1.61	1.03	II	
I140	05:35:19.97	−05:01:02.3	259	0.83	TKK J05351998−0501024	>20.72	>2.61	3.95	I	
I141	05:35:19.98	−05:14:03.8	202	0.03	TKK J05351980−0514054	13.51	0.43	0.38	others	0.090–0.100
I142	05:35:19.99	−05:12:51.1	97	0.86	2MASSI J0535200−051250	14.49	2.06	1.31	II	0.700–0.750
I143	05:35:19.65	−05:13:27.0	143	0.08	TKK J05351966−0513265	13.47	1.95	1.25	II	1.400–1.500
I144	05:35:20.14	−05:13:15.5	953	0.05	2MASSI J0535201−051315	9.74	1.92	1.31	II	>4.000
I145	05:35:20.37	−05:02:26.4	73	−0.86	2MASSI J0535203−050226	12.24	0.77	0.20	III	0.300–0.350
I146	05:35:20.59	−05:03:00.2	7	0.43	TKK J05352062−0503007	18.13	3.73	2.44	II	1.400–1.400
I147	05:35:20.74	−05:15:49.5	2276	−0.12	2MASSI J0535207−051549	9.84	1.18	0.56	III	3.500–4.000
I148	05:35:20.76	−04:58:33.8	673	−0.37	2MASSI J0535207−045834	12.33	1.12	0.46	III	0.570–0.600
I149	05:35:21.13	−05:06:32.4	109	0.63	TKK J05352115−0506324	>20.95	>2.52	3.14	I	
I150	05:35:21.26	−05:09:16.4	2202	−0.81	2MASSI J0535212−050916	8.36	0.55	0.49	II	3.000–3.500
I151	05:35:21.31	−05:12:13.1	14667	−0.63	2MASSI J0535213−051212	8.71	0.60	0.18	III	3.000–3.500
I152	05:35:21.39	−05:09:42.4	29	−0.79	2MASSI J0535213−050942	13.71	0.69	0.49	others	0.110–0.130

(*cont.*)

ID[a]	R.A.[b]	decl.[b]	Counts	HR[c]	NIR counterpart[d]	J^e	$J-H^e$	$H-K^e$	class[f]	mass[gh]
I153	05:35:21.44	−05:09:03.8	142	−0.62	TKK J05352147−0509037	12.23	0.92	>0.37	III	0.400–0.450
I154	05:35:21.50	−05:01:53.5	97	0.13	TKK J05352152−0501539	18.24	3.23	1.92	II	0.350–0.400
I155	05:35:21.56	−05:09:39.2	76	0.84	2MASSI J0535215−050938	13.98	0.84	0.71	others	0.110–0.130
I156	05:35:21.57	−05:09:49.9	260	−0.64	2MASSI J0535215−050949	12.88	0.78	0.53	others	0.200–0.250
I157	05:35:21.86	−04:56:48.7	93	0.08	TKK J05352169−0456487	13.64	0.70	0.11	III	0.110–0.130
I158	05:35:21.87	−05:07:01.9	838	−0.71	2MASSI J0535218−050701	10.93	0.95	0.56	II	1.400–1.500
I159	05:35:21.91	−04:59:17.7	15	−0.20	TKK J05352188−0459197	15.26	1.42	0.80	II	0.130–0.150
I160	05:35:21.92	−05:03:50.6	8	0.25						
I161	05:35:21.94	−05:14:28.3	345	−0.69	2MASSI J0535219−051427	11.75	0.75	0.23	III	0.450–0.500
I162	05:35:22.11	−05:15:05.9	479	−0.32	TKK J05352191−0515012	12.75	0.90	0.32	III	0.250–0.300
I163	05:35:22.34	−04:59:52.9	23	0.65						
I164	05:35:22.35	−05:07:39.2	507	0.02	2MASSI J0535223−050739	14.76	2.62	1.65	II	2.000–2.200
I165	05:35:22.40	−05:08:05.2	833	0.21	2MASSI J0535224−050805	10.88	1.40	0.87	II	3.000–3.500
I166	05:35:22.45	−05:09:11.3	415	−0.92	2MASSI J0535224−050911	11.35	0.58	0.12	III	0.400–0.450
I167	05:35:22.55	−05:08:00.8	921	−0.34	2MASSI J0535225−050800	11.19	1.41	0.53	III	2.700–3.000
I168	05:35:22.67	−05:14:12.0	206	−0.69	2MASSI J0535226−051411	11.69	0.90	0.52	II	0.620–0.700
I169	05:35:23.13	−05:00:35.9	13	−0.08	TKK J05352308−0500364	16.30	2.45	1.38	III	0.350–0.400
I170	05:35:23.20	−05:13:43.9	164	0.04	2MASSI J0535231−051343	13.02	1.76	0.74	III	1.400–1.400
I171	05:35:23.22	−05:08:44.0	13	−0.38	TKK J05352322−0508436	>20.72	>5.77	1.77	others	
I172	05:35:23.33	−04:57:20.5	182	−0.63	TKK J05352330−0457207	12.88	0.83	0.10	others	0.200–0.250
I173	05:35:23.33	−05:08:21.8	59	0.19	TKK J05352335−0508216	16.60	2.40	1.78	I	0.250–0.300
I174	05:35:23.44	−05:10:52.1	1909	−0.82	2MASSI J0535234−051051	11.31	0.63	0.25	III	0.500–0.570
I175	05:35:23.51	−05:07:13.8	33	0.76						
I176	05:35:23.53	−05:15:24.1	96	−0.25	TKK J05352348−0515234	14.03	0.66	0.41	III	0.090–0.100
I177	05:35:24.34	−05:01:20.4	277	0.58	TKK J05352434−0501205	>20.80	>1.51	4.69	I	
I178	05:35:24.58	−05:11:29.7	83	−0.25	2MASSI J0535246−051129	12.11	1.23	0.67	III	0.950–1.000
I179	05:35:24.61	−05:11:58.9	363	−0.81	2MASSI J0535246−051158	11.71	0.77	0.51	others	0.450–0.500
I180	05:35:24.66	−05:09:27.6	9	0.56	TKK J05352469−0509264	16.86	1.81	1.49	I	0.090–0.100
I181	05:35:24.87	−05:06:21.3	144	0.68	2MASSI J0535248−050621	15.63	2.99	1.85	II	2.000–2.200
I182	05:35:25.03	−05:09:09.6	26	−0.85	2MASSI J0535250−050909	14.54	0.63	0.32	III	0.072–0.075
I183	05:35:25.07	−05:10:23.5	22	0.82	2MASSI J0535251−051023	14.86	1.96	1.87	I	0.400–0.450
I184	05:35:25.17	−05:05:09.1	12	0.00						
I185	05:35:25.17	−05:15:38.8	162	−0.06	TKK J05352524−0515357	11.67	0.89	0.68	others	0.620–0.700
I186	05:35:25.22	−05:08:23.8	66	0.94						
I187	05:35:25.24	−05:09:27.6	544	−0.86	2MASSI J0535252−050927	11.77	0.69	0.23	III	0.400–0.450
I188	05:35:25.36	−05:07:20.4	4	0.50						
I189	05:35:25.40	−05:10:48.3	762	−0.83	2MASSI J0535254−051048	11.60	0.60	0.13	III	0.350–0.400
I190	05:35:25.44	−05:06:52.3	6	1.00						
I191	05:35:25.53	−04:57:22.8	145	0.10	TKK J05352567−0457183	13.32	0.73	0.16	III	0.130–0.150

(*cont.*)

ID[a]	R.A.[b]	decl.[b]	Counts	HR[c]	NIR counterpart[d]	J[e]	J − H[e]	H − K[e]	class[f]	mass[gh]
I192	05:35:25.64	−05:07:57.4	324	−0.64	2MASSI J0535256−050757	11.77	0.66	0.89	others	0.350–0.400
I193	05:35:25.72	−05:07:03.4	5	−1.00	2MASSI J0535257−050703	13.72	0.65	0.28	III	0.100–0.110
I194	05:35:25.73	−05:09:49.7	1411	−0.68	2MASSI J0535257−050949	11.09	0.61	0.34	III	0.570–0.600
I195	05:35:25.77	−05:05:57.9	9	0.11	TKK J05352576−0505579	16.21	1.75	1.80	I	0.110–0.130
I196	05:35:25.85	−05:02:45.2	5	1.00						
I197	05:35:25.86	−05:07:56.5	878	−0.79	TKK J05352587−0507564	11.12	0.17	0.86	others	0.200–0.250
I198	05:35:26.05	−05:08:37.9	83	−0.76	TKK J05352606−0508377	12.51	0.59	1.91	others	0.175–0.200
I199	05:35:26.27	−05:07:41.1	8	1.00						
I200	05:35:26.29	−05:08:40.1	3840	−0.84	2MASSI J0535262−050840	10.10	0.50	0.22	II	
I201	05:35:26.34	−05:16:12.0	231	−0.06	2MASSI J0535264−051612	13.28	1.46	0.76	III	0.500–0.570
I202	05:35:26.47	−04:59:52.0	88	0.20	TKK J05352646−0459519	16.52	2.34	1.26	III	0.250–0.300
I203	05:35:26.66	−05:06:10.1	11	1.00						
I204	05:35:26.85	−05:11:07.9	289	0.13	2MASSI J0535268−051107	8.85	1.04	0.77	II	>4.000
I205	05:35:27.07	−05:06:21.7	12	0.67						
I206	05:35:27.15	−05:15:45.3	196	−0.08	2MASSI J0535270−051544	12.75	0.89	0.51	II	0.250–0.300
I207	05:35:27.41	−05:09:04.0	18	−0.89	2MASSI J0535274−050903	14.84	0.77	0.45	III	0.072–0.075
I208	05:35:27.45	−05:02:42.1	270	0.16	TKK J05352744−0502424	16.60	3.32	2.22	II	1.900–2.000
I209	05:35:27.63	−05:09:36.8	24	0.42	2MASSI J0535276−050937	12.90	2.44	1.50	II	3.500–4.000
I210	05:35:27.73	−05:13:55.6	55	0.60	2MASSI J0535275−051356	13.62	1.28	0.66	III	0.300–0.350
I211	05:35:27.78	−05:17:03.0	711	−0.46	TKK J05352793−0516573	11.58	0.73	0.27	III	0.500–0.570
I212	05:35:27.83	−05:05:36.1	345	0.57	TKK J05352786−0505363	>20.62	>4.33	2.79	II	
I213	05:35:28.06	−05:01:34.8	326	0.42	TKK J05352806−0501351	18.99	4.09	2.69	II	1.400–1.400
I214	05:35:28.14	−05:10:13.5	37	−0.41	2MASSI J0535281−051013	11.16	1.00	0.66	II	1.400–1.400
I215	05:35:28.18	−05:03:40.9	41	0.41	TKK J05352819−0503413	20.03	3.61	4.62	I	0.200–0.250
I216	05:35:28.19	−05:00:49.2	172	−0.79	2MASSI J0535281−050049	12.49	0.65	0.22	III	0.200–0.250
I217	05:35:28.21	−05:11:37.2	50	−0.28	2MASSI J0535281−051137	12.76	1.06	0.52	III	0.350–0.400
I218	05:35:28.27	−04:58:38.0	1228	0.61	TKK J05352826−0458384	>20.58	>6.47	2.65	others	
I219	05:35:28.50	−05:07:46.8	89	0.03	TKK J05352852−0507469	15.99	2.99	1.88	II	1.500–1.600
I220	05:35:28.60	−05:05:44.2	28	0.29	2MASSI J0535285−050544	14.92	2.58	1.50	II	1.500–1.600
I221	05:35:28.67	−05:03:06.1	8	0.50						
I222	05:35:28.71	−05:05:51.2	11	0.82						
I223	05:35:29.02	−05:06:03.8	1424	−0.82	2MASSI J0535290−050604	11.64	0.58	0.17	III	0.350–0.400
I224	05:35:29.25	−05:08:05.3	6	0.67						
I225	05:35:29.45	−05:16:33.5	1189	−0.53	2MASSI J0535294−051633	11.59	0.86	0.26	III	0.620–0.700
I226	05:35:29.81	−05:16:07.6	476	−0.63	2MASSI J0535298−051606	12.06	0.78	0.31	III	0.350–0.400
I227	05:35:29.89	−05:12:10.6	136	−0.78	2MASSI J0535299−051210	12.68	0.65	0.20	III	0.175–0.200
I228	05:35:29.90	−05:04:26.7	8	1.00						
I229	05:35:30.00	−05:12:27.7	157	−0.61	2MASSI J0535299−051227	12.08	0.70	0.49	others	0.300–0.350
I230	05:35:30.11	−05:09:09.5	23	−0.74	2MASSI J0535301−050909	13.77	0.60	0.27	III	0.100–0.110

(*cont.*)

ID[a]	R.A.[b]	decl.[b]	Counts	HR[c]	NIR counterpart[d]	J^e	$J-H^e$	$H-K^e$	class[f]	mass[gh]
I231	05:35:30.13	−05:14:19.0	148	−0.27	2MASSI J0535301−051418	13.16	1.18	0.61	III	0.350–0.400
I232	05:35:30.32	−05:13:52.0	213	−0.10	2MASSI J0535302−051352	13.19	1.61	0.85	III	0.900–0.950
I233	05:35:30.55	−05:03:33.9	13	−0.08	TKK J05353054−0503345	15.31	1.67	0.87	III	0.175–0.200
I234	05:35:30.64	−04:59:35.6	222	−0.34	2MASSI J0535306−045936	14.31	2.29	1.52	II	1.400–1.500
I235	05:35:30.72	−05:06:45.3	6	0.33						
I236	05:35:30.86	−04:58:12.4	67	0.25	TKK J05353080−0458136	17.60	3.13	1.65	III	0.500–0.570
I237	05:35:31.06	−05:04:14.2	91	−0.93	2MASSI J0535310−050415	11.13	0.57	0.12	II	0.500–0.570
I238	05:35:31.22	−05:12:28.3	311	−0.80	2MASSI J0535312−051228	13.03	0.66	0.23	III	0.150–0.175
I239	05:35:31.27	−05:12:01.8	82	−0.41	2MASSI J0535313−051201	14.57	0.62	0.30	III	0.070–0.072
I240	05:35:31.31	−05:15:33.3	2773	−0.64	2MASSI J0535312−051533	10.05	0.78	0.32	III	2.000–2.200
I241	05:35:31.46	−05:05:45.5	22	0.91						
I242	05:35:31.48	−05:16:03.2	660	−0.78	2MASSI J0535313−051602	5.86	0.23	0.05	III	>4.000
I243	05:35:31.49	−05:05:01.4	187	−0.80	2MASSI J0535315−050501	11.96	0.69	0.48	others	0.350–0.400
I244	05:35:31.57	−05:05:47.1	13	0.85	2MASSI J0535315−050547	11.72	2.00	1.42	II	>4.000
I245	05:35:31.58	−05:16:56.4	181	−0.28	2MASSI J0535316−051658	12.95	1.17	0.50	III	0.400–0.450
I246	05:35:31.61	−05:00:14.0	63	0.49	2MASSI J0535316−050014	14.25	2.63	1.64	II	2.700–3.000
I247	05:35:31.91	−05:05:49.3	33	0.39						
I248	05:35:31.96	−05:09:27.9	2062	−0.65	2MASSI J0535319−050927	9.34	0.81	0.41	II	3.000–3.500
I249	05:35:32.02	−05:08:05.2	7	−0.14	2MASSI J0535320−050805	15.16	1.05	0.62	II	0.080–0.090
I250	05:35:32.24	−05:11:58.3	38	−0.53	2MASSI J0535321−051157	12.61	0.75	0.48	III	0.200–0.250
I251	05:35:32.31	−05:11:44.0	175	0.29	2MASSI J0535323−051143	14.58	2.18	1.14	III	0.900–0.950
I252	05:35:32.35	−05:14:26.3	150	−0.45	2MASSI J0535324−051424	12.62	0.98	0.36	III	0.350–0.400
I253	05:35:32.35	−05:08:29.8	25	0.76						
I254	05:35:32.51	−05:02:09.9	14	−0.71	2MASSI J0535325−050209	14.71	0.60	0.19	III	0.060–0.070
I255	05:35:32.80	−05:07:51.6	23	0.65						
I256	05:35:32.86	−05:16:04.6	1143	−0.39	2MASSI J0535329−051605	11.59	1.09	0.30	others	
I257	05:35:32.97	−05:12:05.9	25	0.12	2MASSI J0535329−051204	13.77	2.04	1.15	III	1.400–1.400
I258	05:35:33.00	−05:17:32.3	406	−0.63	2MASSI J0535331−051733	12.81	0.62	0.27	III	0.150–0.175
I259	05:35:33.16	−05:14:11.1	974	−0.14	2MASSI J0535331−051410	12.36	0.78	0.40	III	0.300–0.350
I260	05:35:33.37	−05:08:02.0	11	0.82						
I261	05:35:33.52	−05:15:19.3	646	−0.11	TKK J05353359−0515232	13.08	1.24	0.38	others	0.400–0.450
I262	05:35:33.61	−05:00:41.7	128	−0.83	2MASSI J0535336−050042	11.93	0.65	0.20	III	0.300–0.350
I263	05:35:33.64	−05:03:07.7	223	−0.35	2MASSI J0535336−050308	13.06	1.45	0.69	III	0.620–0.700
I264	05:35:33.83	−05:04:27.2	346	0.88	2MASSI J0535338−050427	13.75	1.93	1.26	II	
I265	05:35:33.86	−05:09:05.6	20	−0.80	2MASSI J0535338−050905	13.99	0.59	0.26	III	0.080–0.090
I266	05:35:34.26	−04:59:52.7	53	0.89						
I267	05:35:34.48	−05:13:07.3	76	−0.61	2MASSI J0535345−051307	14.14	0.58	0.16	III	0.080–0.090
I268	05:35:34.52	−05:00:51.8	27	0.48	2MASSI J0535345−050052	15.79	2.31	2.17	I	0.400–0.450
I269	05:35:34.68	−05:15:53.3	330	−0.67	2MASSI J0535346−051552	13.33	0.56	0.55	others	0.110–0.130

(*cont.*)

ID[a]	R.A.[b]	decl.[b]	Counts	HR[c]	NIR counterpart[d]	J[e]	J − H[e]	H − K[e]	class[f]	mass[gh]
I270	05:35:35.11	−05:14:47.2	99	0.33						
I271	05:35:35.12	−05:05:58.2	13	1.00						
I272	05:35:35.36	−05:11:11.8	506	−0.72	2MASSI J0535353−051111	12.66	0.58	0.22	III	0.150–0.175
I273	05:35:35.48	−05:07:20.2	27	0.93						
I274	05:35:35.55	−05:06:58.5	210	−0.90	2MASSI J0535355−050658	13.35	0.63	0.21	III	0.110–0.130
I275	05:35:35.63	−05:01:48.5	66	0.91						
I276	05:35:35.64	−05:10:50.0	22	−0.55	2MASSI J0535356−051050	14.24	0.52	0.35	III	0.072–0.075
I277	05:35:35.71	−05:15:43.3	204	−0.26	2MASSI J0535356−051543	11.63	0.68	0.28	III	0.400–0.450
I278	05:35:36.02	−05:12:25.1	270	−0.61	2MASSI J0535360−051225	11.09	1.55	1.04	II	3.000–3.500
I279	05:35:36.19	−05:04:55.8	114	0.60	2MASSI J0535362−050455	13.17	1.35	0.72	III	0.450–0.500
I280	05:35:36.40	−05:01:15.2	943	−0.46	2MASSI J0535364−050115	10.89	0.84	0.47	II	1.400–1.400
I281	05:35:36.51	−05:05:22.3	9	0.33						
I282	05:35:36.56	−05:04:39.1	909	−0.92	2MASSI J0535365−050439	12.24	0.65	0.19	III	0.250–0.300
I283	05:35:36.69	−05:04:14.0	2221	−0.67	2MASSI J0535366−050414	11.96	0.75	0.32	III	0.350–0.400
I284	05:35:36.77	−05:09:59.1	11	−0.45	2MASSI J0535367−051000	13.30	0.80	0.56	others	0.150–0.175
I285	05:35:36.86	−05:07:32.4	10	0.60						
I286	05:35:36.95	−05:05:25.8	24	−0.33	2MASSI J0535369−050526	14.59	0.51	0.43	others	0.060–0.070
I287	05:35:37.17	−05:10:29.6	312	−0.88	2MASSI J0535371−051029	12.53	0.68	0.27	III	0.200–0.250
I288	05:35:37.57	−05:04:47.5	16	0.12	TKK J05353759−0504471	17.46	2.62	1.59	II	0.200–0.250
I289	05:35:37.70	−05:06:31.8	298	−0.64	2MASSI J0535376−050631	12.95	0.67	0.28	III	0.150–0.175
I290	05:35:38.04	−05:15:08.3	122	0.48						
I291	05:35:38.07	−05:03:17.1	13	0.38						
I292	05:35:38.30	−05:14:19.1	246	−0.42	2MASSI J0535382−051418	11.56	0.75	0.22	III	0.500–0.570
I293	05:35:38.52	−04:59:40.2	518	−0.31	TKK J05353853−0459410	12.00	0.92	>0.62	II	0.500–0.570
I294	05:35:38.57	−05:08:03.4	75	0.12	2MASSI J0535385−050803	15.96	2.53	1.36	III	0.500–0.570
I295	05:35:38.65	−05:09:56.9	320	−0.72	2MASSI J0535386−050956	13.06	0.58	0.30	III	0.130–0.150
I296	05:35:38.75	−05:04:55.3	40	−0.55	2MASSI J0535387−050455	14.96	1.94	1.03	III	0.350–0.400
I297	05:35:38.88	−05:12:42.2	1047	−0.26	2MASSI J0535388−051241	10.81	1.11	0.66	II	2.200–2.500
I298	05:35:39.03	−05:07:04.1	279	−0.84	2MASSI J0535390−050704	12.60	0.56	0.29	III	0.150–0.175
I299	05:35:39.08	−05:08:56.4	198	−0.80	2MASSI J0535390−050856	10.86	0.66	0.16	III	0.750–0.800
I300	05:35:39.17	−05:05:40.0	7	0.14						
I301	05:35:39.39	−05:05:07.7	113	0.59						
I302	05:35:39.72	−04:59:18.4	65	0.35						
I303	05:35:39.97	−05:06:36.5	75	−0.68	2MASSI J0535399−050636	12.67	1.13	0.61	III	0.450–0.500
I304	05:35:40.19	−05:04:27.3	9	0.33						
I305	05:35:40.47	−05:04:18.7	19	0.58						
I306	05:35:40.59	−05:12:19.8	85	−0.13	2MASSI J0535406−051219	12.96	0.69	0.27	III	0.150–0.175
I307	05:35:40.66	−05:11:10.0	76	0.11	2MASSI J0535407−051111	13.77	0.65	0.29	III	0.100–0.110
I308	05:35:40.80	−05:09:01.1	825	−0.61	2MASSI J0535407−050901	10.33	0.90	0.30	II	2.200–2.500

(*cont.*)

ID[a]	R.A.[b]	decl.[b]	Counts	HR[c]	NIR counterpart[d]	J^e	$J-H^e$	$H-K^e$	class[f]	mass[gh]
I309	05:35:41.04	−05:06:24.9	119	−0.76	2MASSI J0535410−050625	12.66	0.62	0.30	III	0.175–0.200
I310	05:35:41.38	−05:03:52.6	38	−0.58	2MASSI J0535413−050352	14.21	1.17	0.49	III	0.150–0.175
I311	05:35:41.72	−05:03:28.3	61	−0.74	2MASSI J0535417−050329	13.45	0.61	0.20	III	0.110–0.130
I312	05:35:41.81	−05:05:19.1	15	−0.33	2MASSI J0535417−050519	15.52	1.70	0.83	III	0.150–0.175
I313	05:35:42.04	−05:10:12.0	660	−0.32	2MASSI J0535420−051011	11.47	1.33	0.62	III	2.000–2.200
I314	05:35:42.04	−05:13:00.0	617	−0.73	2MASSI J0535420−051259	13.45	0.56	0.29	III	0.110–0.130
I315	05:35:42.40	−05:11:02.5	56	0.32						
I316	05:35:42.64	−04:59:55.4	79	0.34						
I317	05:35:42.78	−05:11:55.1	658	−0.79	2MASSI J0535427−051154	11.18	0.70	0.16	III	0.620–0.700
I318	05:35:42.98	−05:03:05.3	42	−0.33	2MASSI J0535430−050307	13.34	0.64	0.27	III	0.110–0.130
I319	05:35:43.27	−05:09:16.8	2262	−0.79	2MASSI J0535432−050917	10.76	0.68	0.17	II	0.900–0.950
I320	05:35:43.54	−05:05:40.2	48	−0.08	2MASSI J0535435−050541	11.96	0.87	0.69	others	0.450–0.500
I321	05:35:43.82	−05:09:59.8	174	−0.67	2MASSI J0535438−050958	13.69	0.60	0.27	III	0.100–0.110
I322	05:35:44.52	−05:07:31.4	288	−0.87	2MASSI J0535445−050731	12.25	0.70	0.20	II	0.250–0.300
I323	05:35:44.54	−05:00:04.9	45	0.24	TKK J05354461−0459576	14.91	0.65	0.22	III	0.060–0.070
I324	05:35:44.88	−05:07:16.7	5273	−0.70	2MASSI J0535448−050716	10.07	0.63	0.25	II	1.500–1.600
I325	05:35:45.70	−05:06:44.7	74	0.51						
I326	05:35:46.10	−05:10:52.4	695	−0.76	2MASSI J0535461−051051	11.83	0.78	0.41	III	0.450–0.500
I327	05:35:47.25	−05:02:50.1	64	0.41						
I328	05:35:47.77	−05:10:31.2	3604	−0.81	2MASSI J0535477−051030	10.18	0.54	0.21	II	1.400–1.400
I329	05:35:48.26	−05:11:15.5	60	0.07	TKK J05354826−0511103	12.77	0.82	0.26	III	0.200–0.250
I330	05:35:48.39	−05:05:20.6	67	0.55						
I331	05:35:48.40	−05:01:27.9	193	−0.60	2MASSI J0535483−050128	11.84	0.93	0.68	II	0.570–0.600
I332	05:35:48.88	−05:00:31.6	88	−0.45	TKK J05354884−0500285	12.63	0.96	0.29	II	0.350–0.400
I333	05:35:49.00	−05:05:42.2	31	0.10						
I334	05:35:49.04	−05:01:40.6	288	−0.49	2MASSI J0535489−050139	12.59	0.77	0.35	II	0.250–0.300
I335	05:35:49.21	−05:02:23.2	49	0.55						
I336	05:35:49.96	−05:04:24.4	88	0.39						
I337	05:35:50.07	−05:10:22.2	63	−0.33	TKK J05355012−0510294	13.21	0.58	0.24	III	0.110–0.130
I338	05:35:50.51	−05:02:24.7	63	0.27						
I339	05:35:51.12	−05:07:08.7	1348	−0.78	2MASSI J0535510−050708	11.00	0.74	0.37	III	0.900–0.950
I340	05:35:51.46	−05:02:03.7	167	0.46						
I341	05:35:51.61	−05:05:58.7	48	0.25						
I342	05:35:51.66	−05:08:09.0	7972	−0.73	2MASSI J0535516−050809	10.04	0.62	0.14	III	1.500–1.600
I343	05:35:52.66	−05:05:05.2	4600	−0.75	2MASSI J0535526−050505	9.82	0.58	0.37	III	1.600–1.700
I344	05:35:53.71	−05:02:33.8	164	−0.07	2MASSI J0535535−050234	13.12	1.04	0.35	III	0.250–0.300
I345	05:35:53.86	−05:06:54.3	258	0.34						
I346	05:35:54.05	−05:04:14.5	1398	−0.81	2MASSI J0535540−050414	11.29	0.87	0.49	II	0.900–0.950
I347	05:35:54.22	−05:05:43.2	74	−0.03	2MASSI J0535542−050545	13.55	0.65	0.20	III	0.110–0.130

(*cont.*)

ID[a]	R.A.[b]	decl.[b]	Counts	HR[c]	NIR counterpart[d]	J^e	$J-H^e$	$H-K^e$	class[f]	mass[gh]
I348	05:35:54.67	−05:06:28.6	168	−0.56	2MASSI J0535546−050627	13.48	0.55	0.30	III	0.100–0.110
I349	05:35:55.86	−05:04:42.5	87	0.24	TKK J05355574−0504377	16.88	1.26	1.44	I	0.040–0.050
I350	05:35:57.79	−05:05:02.8	169	0.14						
I351	05:36:00.05	−05:04:32.2	134	−0.27	2MASSI J0535599−050430	13.54	0.57	0.41	others	0.100–0.110
I352	05:36:00.36	−05:05:50.5	68	0.06	TKK J05360045−0505540	13.38	0.55	0.19	III	0.110–0.130
I353	05:36:00.55	−05:05:03.5	413	−0.27	TKK J05360033−0505000	12.22	0.57	0.19	III	0.200–0.250
I354	05:36:04.60	−05:04:00.5	250	0.10	TKK J05360415−0504088	12.75	0.74	0.31	III	0.200–0.250
I355	05:34:54.47	−05:04:12.5	18	0.67						
I356	05:34:57.03	−05:02:50.2	6	1.00	TKK J05345731−0502508	18.31	1.82	0.77	others	0.040–0.050
I357	05:35:08.43	−05:12:24.5	34	0.47						
I358	05:35:13.13	−05:07:51.5	4	1.00						
I359	05:35:18.87	−05:09:51.4	5	1.00						
I360	05:35:21.33	−05:12:46.3	15	0.47	TKK J05352138−0512444	15.88	0.70	0.53	others	0.040–0.050
I361	05:35:21.50	−05:07:26.0	4	1.00						
I362	05:35:22.42	−05:02:40.7	4	1.00						
I363	05:35:22.67	−05:06:49.5	3	1.00						
I364	05:35:29.78	−04:59:53.9	27	0.63	TKK J05352956−0459567	17.93	0.77	1.08	others	<0.002
I365	05:35:30.67	−05:04:09.9	4	1.00	TKK J05353065−0504111	18.18	1.76	1.39	II	0.040–0.050
I366	05:35:31.19	−05:09:15.7	8	1.00						
I367	05:35:33.14	−05:15:08.7	268	0.26	TKK J05353244−0515067	13.53	1.30	0.75	II	0.300–0.350
I368	05:35:47.84	−05:08:21.5	28	0.57						
I369	05:35:50.91	−05:01:04.9	91	0.45						
I370	05:34:44.08	−05:05:58.3	57	0.30						
I371	05:34:52.91	−05:08:03.2	18	0.00						
I372	05:35:00.97	−05:05:39.6	8	−0.75	2MASSI J0535009−050540	15.81	0.41	0.35	others	0.030–0.040
I373	05:35:07.89	−05:09:03.5	3	−1.00						
I374	05:35:10.79	−05:10:43.5	3	−1.00						
I375	05:35:14.97	−05:04:43.1	3	−1.00						
I376	05:35:19.73	−05:11:43.5	34	0.00						
I377	05:35:29.82	−05:13:29.1	49	0.18						
I378	05:35:30.76	−05:03:35.5	5	−1.00	TKK J05353072−0503355	15.25	1.66	0.81	III	0.175–0.200
I379	05:35:33.75	−04:59:26.9	15	−0.07						
I380	05:35:35.70	−05:09:07.4	6	−0.33						
I381	05:35:38.44	−05:10:08.7	28	−0.50	2MASSI J0535384−051009	14.85	0.61	0.31	III	0.060–0.070
I382	05:35:42.37	−05:02:50.3	15	−0.60						
I383	05:35:44.67	−05:00:38.6	31	−0.55	2MASSI J0535447−050039	14.04	0.75	0.35	III	0.100–0.110
I384	05:35:46.94	−05:06:58.9	14	−0.29						
I385	05:35:54.84	−05:05:25.8	23	−0.30	TKK J05355468−0505213	16.87	1.10	0.67	II	0.040–0.050
S1	05:35:54.73	−04:58:07.1	2987	−0.87	2MASSI J0535546−045808	10.40	0.48	0.17	III	0.750–0.800

(*cont.*)

ID[a]	R.A.[b]	decl.[b]	Counts	HR[c]	NIR counterpart[d]	J^e	$J-H^e$	$H-K^e$	class[f]	mass[gh]
S2	05:35:56.14	−04:56:55.9	1681	−0.75	2MASSI J0535560−045655	11.74	0.62	0.19	III	0.350–0.400
S3	05:35:56.28	−04:55:04.1	510	0.13						
S4	05:35:58.84	−04:55:35.2	799	0.03	2MASSI J0535588−045537	>15.42	>0.36	0.81	II	
S5	05:35:59.31	−04:58:44.3	743	−0.64	2MASSI J0535592−045846	12.55	0.70	0.16	III	0.200–0.250
S6	05:36:02.99	−04:57:34.4	304	−0.30						
S7	05:36:04.24	−04:57:42.5	285	−0.25						
S8	05:36:05.75	−04:51:08.5	832	−0.39						
S9	05:36:11.84	−05:00:30.4	567	−0.47	2MASSI J0536118−050032	12.23	0.69	0.17	III	0.250–0.300
S10	05:36:12.81	−04:55:13.7	462	−0.29						
S11	05:36:28.11	−04:57:07.5	185	0.43						
S12	05:36:14.94	−04:57:05.9	72	0.47						
S13	05:36:20.26	−04:54:22.2	164	0.34						

[a] I1–I354 and S1–S11 are detected in the total band (0.5–8.0 keV) image of ACIS-I and ACIS-S2. I355-I369 and S12-S13 are detected only in the hard band (2.0–8.0 keV) image of ACIS-I and ACIS-S2. I370–I385 are detected only in the soft band (0.5–2.0 keV) image of ACIS-I. No new source was detected in the soft band image of ACIS-S2.

[b] The equinox in J2000.0.

[c] Defined as $(H-S)/(H+S)$ where H and S are the X-ray photon counts in the hard (2.0–8.0 keV) and soft (0.5–2.0 keV) band, respectively.

[d] The QUIRC and 2MASS counterpart has the prefix "TKK J" and "2MASSI J", respectively. The nomenclatures follow the IAU convention.

[e] NIR colors in the CIT color system.

[f] The evolutional classes of NIR-IDed *Chandra* sources estimated from the NIR excess, UV excess, and H_α emission data (see Sect. 6.2).

[g] The mass of NIR-IDed *Chandra* sources estimated from the $J/(J–H)$ color-magnitude diagram (see Sect. 6.2).

[h] Four sources have two estimates of their mass on Figure 6.7; i.e., 1.050–1.000 $M_\odot$ or 1.300–1.400 $M_\odot$ for I139 and I264, and 1.100–1.150 $M_\odot$ or 1.400–1.400 $M_\odot$ for I200 and I256.

Appendix B

QUIRC Sources and 2MASS Counterpart

Table B.1: QUIRC source list

ID	R.A.[a]	decl.[a]	J^{bcd} (mag)	H^{bcd} (mag)	K^{cd} (mag)	2MASS counterpart
1	34:28.94	−5:08:38.7	15.52	14.13	13.74	0534289−050838
2	34:29.36	−5:11:14.4	17.45	16.75	16.51	
3	34:29.48	−5:05:24.7	17.64	16.82	16.44	
4	34:29.60	−5:04:29.0	14.05	13.47	13.15	0534295−050428
5	34:29.86	−5:04:05.6	12.83	12.16	11.95	0534298−050405
6	34:30.15	−5:11:56.7	15.70	14.35	13.99	0534301−051156
7	34:30.27	−5:11:48.2	11.53	10.64†	9.95†	0534302−051148
8	34:30.61	−5:11:05.0	13.90	13.30	13.08	0534305−051104
9	34:30.97	−5:03:56.9	16.94	16.14	15.80	
10	34:31.07	−5:06:17.1	17.42	16.64	16.23	
11	34:31.47	−5:08:54.1	16.45	15.84	15.56	
12	34:31.52	−5:12:30.0	9.49†	8.40†	7.95†	0534315−051230
13	34:31.58	−5:11:50.3	17.78	16.59	16.21	
14	34:31.85	−5:04:09.3	13.31	12.66	12.50	0534318−050408
15	34:32.03	−5:10:23.2	15.21	14.63	14.51	0534320−051023
16	34:32.05	−5:11:24.8	13.01	12.19	11.83	0534320−051124
17	34:32.36	−5:12:59.6	18.64	15.67	15.12	
18	34:32.55	−5:04:21.3	17.65	16.83	16.63	
19	34:32.88	−5:10:15.5	16.25	15.33	15.06	0534328−051015
20	34:32.89	−5:09:30.9	20.34	>20.72	17.05	
21	34:33.26	−5:07:26.8	14.51	14.22	13.85	0534332−050726
22	34:33.27	−5:11:44.1	17.50	16.44	16.00	
23	34:34.06	−5:12:27.0	15.37	14.43	14.14	0534340−051227
24	34:34.16	−5:05:57.0	17.18	16.44	16.12	
25	34:34.17	−5:05:17.1	13.30	12.71	12.40	0534341−050516
26	34:34.62	−5:04:49.7	14.02	13.25	13.03	0534346−050449
27	34:34.66	−5:05:08.4	16.22	15.37	15.18	0534346−050508
28	34:35.55	−5:08:45.0	17.42	16.49	16.19	
29	34:35.55	−5:09:04.8	17.43	16.77	16.46	
30	34:35.95	−5:09:03.2	16.10	15.27	14.94	0534359−050903
31	34:36.14	−5:05:08.0	16.34	15.48	15.46	0534361−050507
32	34:36.26	−5:00:48.2	19.98	18.96	16.79	
33	34:36.35	−5:02:54.7	16.61	15.67	15.40	
34	34:36.39	−5:11:40.6	14.82	14.19	14.12	0534363−051140
35	34:36.69	−5:05:03.4	16.05	15.23	15.05	0534366−050503
36	34:36.69	−5:08:17.5	18.13	17.49	16.55	

(cont.)

ID	R.A.[a]	decl.[a]	J^{bcd} (mag)	H^{bcd} (mag)	K^{cd} (mag)	2MASS counterpart
37	34:36.72	−5:09:37.3	17.71	16.98	16.44	
38	34:36.76	−5:01:55.3	20.87	19.90	17.27	
39	34:36.83	−5:06:09.9	15.74	14.96	14.73	0534368−050609
40	34:36.84	−5:05:19.8	17.75	16.71	15.68	
41	34:36.92	−5:05:11.9	16.92	16.46	16.09	
42	34:36.93	−5:08:14.5	17.01	16.77	16.65	
43	34:37.03	−5:12:39.1	16.80	15.41	14.98	
44	34:37.63	−5:04:49.3	16.82	15.79	14.94	
45	34:37.80	−5:14:29.4	15.17	14.08	13.72	0534377−051429
46	34:37.91	−5:08:48.0	16.25	15.28	15.00	0534379−050847
47	34:37.91	−5:13:32.3	14.46	13.93	13.63	0534378−051332
48	34:38.00	−5:14:06.4	16.77	16.17	15.59	
49	34:38.08	−5:13:09.7	16.75	15.73	15.36	
50	34:38.13	−5:03:05.5	16.26	15.26	14.55	0534381−050305
51	34:38.13	−5:04:57.9	12.30	11.70	11.40	0534381−050458
52	34:38.16	−5:05:17.2	14.35	13.84	13.33	0534381−050517
53	34:38.34	−5:01:26.0	15.26	14.24	13.90	0534383−050125
54	34:38.39	−5:02:06.5	19.60	19.02	16.59	
55	34:38.49	−5:10:42.2	16.42	15.49	15.09	0534384−051042
56	34:38.57	−5:04:19.5	15.73	14.78	14.54	0534385−050419
57	34:38.58	−5:00:44.9	16.12	15.30	15.06	0534385−050044
58	34:38.74	−5:04:05.7	15.72	15.18	14.72	0534387−050405
59	34:38.86	−5:04:40.7	16.05	15.26	14.97	0534388−050440
60	34:38.90	−5:02:36.6	13.65	12.91	12.69	0534388−050236
61	34:38.90	−5:14:28.5	14.45	13.34	13.02	0534388−051428
62	34:39.20	−5:04:52.9	16.44	15.01	14.03	
63	34:39.24	−5:01:48.9	13.81	13.24	12.83	
64	34:39.29	−5:00:47.1	14.62	13.94	13.73	0534393−050046
65	34:39.35	−5:11:35.1	18.24	16.94	16.23	
66	34:39.38	−5:01:46.7	12.86	12.26	11.89	0534393−050146
67	34:39.43	−5:03:43.7	18.09	16.93	16.05	
68	34:39.75	−5:03:06.2	12.95	12.16	11.95	0534397−050306
69	34:39.78	−5:00:34.0	16.19	15.03	14.39	
70	34:39.86	−5:02:13.2	16.73	15.51	15.12	
71	34:39.88	−5:07:59.7	16.48	16.06	15.84	0534398−050800
72	34:39.97	−5:04:08.8	15.91	15.03	14.80	0534399−050408

(cont.)

ID	R.A.[a]	decl.[a]	J^{bcd}	H^{bcd}	K^{cd}	2MASS
73	34:39.98	−5:10:07.0	9.11†	8.96†	8.83†	0534399−051007
74	34:40.41	−5:07:53.2	15.85	15.43	15.21	0534403−050753
75	34:40.64	−5:06:58.6	15.67	13.21	11.87	0534406−050658
76	34:40.66	−5:12:24.5	15.10	14.49	14.20	0534406−051224
77	34:40.73	−5:09:24.1	15.62	14.98	14.66	0534407−050924
78	34:41.32	−5:01:19.6	15.36	14.34	14.03	0534413−050119
79	34:41.53	−5:07:02.4	15.01	14.69	14.29	0534415−050702
80	34:41.55	−5:02:25.2	16.34	15.58	15.37	0534415−050225
81	34:41.84	−5:08:39.7	16.80	16.29	16.04	
82	34:42.01	−5:06:41.9	16.84	16.15	15.81	
83	34:42.03	−5:02:25.0	14.01	13.33	13.09	0534420−050224
84	34:42.05	−5:04:31.8	13.95	13.23	12.52	0534420−050431
85	34:42.28	−5:07:14.6	9.58†	9.64†	9.58†	0534422−050714
86	34:42.39	−5:04:00.0	17.78	16.83	16.44	
87	34:42.41	−5:12:38.2	13.22	12.40	12.13	0534424−051238
88	34:42.42	−5:12:18.8	10.00†	9.18†	8.86†	0534424−051218
89	34:42.60	−5:03:16.2	15.96	15.19	15.00	0534425−050316
90	34:42.69	−4:58:55.9	17.29	16.42	15.67	
91	34:43.09	−5:00:11.6	16.22	15.72	15.43	0534430−050011
92	34:43.16	−5:13:33.8	17.65	16.72	16.31	
93	34:43.18	−5:01:57.9	16.91	15.87	15.66	
94	34:43.24	−5:07:28.3	16.44	15.86	15.77	
95	34:43.48	−5:14:42.5	15.40	14.09	13.84	0534434−051442
96	34:43.50	−5:14:21.2	15.34	14.38	14.12	0534434−051421
97	34:43.52	−5:01:26.3	17.11	16.33	16.06	
98	34:43.58	−5:06:24.5	16.37	15.81	15.53	0534435−050624
99	34:43.65	−5:07:13.1	16.32	15.65	15.27	
100	34:43.88	−5:14:22.2	16.93	15.88	15.80	
101	34:43.89	−5:12:55.8	14.21	13.16	12.89	0534438−051255
102	34:43.96	−4:59:33.4	13.48	12.80	12.58	0534439−045933
103	34:44.08	−5:06:26.5	16.72	16.19	15.55	0534440−050626
104	34:44.14	−4:59:10.1	17.36	16.20	15.70	
105	34:44.18	−5:06:43.2	13.93	13.30	13.12	0534441−050643
106	34:44.48	−5:02:55.3	17.51	16.48	16.23	
107	34:44.57	−5:02:13.5	21.20	18.76	16.66	
108	34:44.90	−5:12:31.8	15.28	14.33	13.75	0534448−051231
109	34:44.98	−5:06:49.5	11.74	11.25	10.75†	0534449−050649
110	34:45.07	−5:06:20.1	13.79	12.58	12.14	0534450−050620
111	34:45.12	−5:14:06.1	14.57	13.46	12.91	0534451−051406

(*cont.*)

ID	R.A.[a]	decl.[a]	J^{bcd}	H^{bcd}	K^{cd}	2MASS
112	34:45.14	−5:00:05.9	16.37	15.48	15.12	0534451−050005
113	34:45.19	−5:10:47.6	12.21	11.43	11.04	0534451−051047
114	34:45.33	−5:00:50.0	14.35	13.53	13.27	0534453−050049
115	34:45.45	−5:02:07.9	12.27	11.58	11.33	0534454−050207
116	34:45.53	−5:10:16.0	14.65	13.99	13.89	0534455−051016
117	34:45.95	−5:04:14.3	17.53	16.73	16.56	
118	34:46.02	−5:08:58.1	16.53	15.63	15.31	0534460−050858
119	34:46.08	−5:13:11.6	15.77	14.96	14.41	
120	34:46.35	−5:01:22.8	>21.62	>20.67	16.70	
121	34:46.60	−5:04:46.6	15.75	15.05	14.92	0534465−050446
122	34:46.63	−5:10:40.4	12.05	11.47	11.32	0534466−051040
123	34:46.79	−5:04:26.8	16.85	15.81	15.37	
124	34:46.82	−4:59:10.9	12.31	11.51	11.16	
125	34:46.91	−5:06:46.3	17.30	16.38	16.23	
126	34:46.94	−4:59:13.1	12.13	11.38	11.03	0534469−045912
127	34:47.03	−5:08:12.4	17.18	16.51	16.19	
128	34:47.08	−5:00:22.3	16.68	15.78	15.51	0534470−050022
129	34:47.09	−5:08:45.6	16.62	15.58	15.45	
130	34:47.68	−4:59:16.3	>21.51	19.94	16.71	
131	34:47.68	−4:59:33.3	13.65	13.00	12.75	0534476−045933
132	34:47.85	−5:00:18.7	16.89	15.83	15.38	
133	34:47.96	−5:04:54.9	11.89	11.16	10.67†	0534479−050455
134	34:48.25	−5:10:39.0	16.00	15.32	15.27	0534482−051038
135	34:48.26	−5:03:30.3	17.05	15.99	15.89	
136	34:48.28	−5:09:00.8	15.25	14.22	13.93	0534482−050900
137	34:48.29	−5:07:13.5	12.83	12.16	12.02	0534482−050713
138	34:48.42	−5:05:01.4	12.47	11.79	11.64	0534484−050501
139	34:48.49	−5:01:52.3	17.57	15.92	15.67	
140	34:48.51	−5:07:11.0	13.30	12.70	12.51	0534485−050710
141	34:48.77	−5:02:46.8	17.02	16.06	15.57	
142	34:48.98	−5:06:56.9	17.71	16.47	15.94	
143	34:49.17	−4:56:01.0	19.36	16.27	14.88	
144	34:49.18	−5:01:38.1	17.53	16.45	15.85	
145	34:49.22	−5:04:38.1	12.48	11.77	11.58	0534492−050438
146	34:49.35	−5:11:53.5	17.21	16.09	15.70	
147	34:49.37	−4:59:38.2	16.61	15.70	15.35	0534493−045937
148	34:49.49	−5:11:34.8	17.78	16.19	15.51	
149	34:49.58	−5:04:59.6	13.31	12.59	12.29	0534495−050459
150	34:49.59	−5:11:18.0	17.13	16.19	16.01	

(*cont.*)

ID	R.A.[a]	decl.[a]	J^{bcd}	H^{bcd}	K^{cd}	2MASS
151	34:49.69	−5:07:07.4	17.62	16.23	15.83	
152	34:49.73	−5:00:03.0	15.27	14.71	14.55	0534497−050003
153	34:50.04	−5:00:37.4	17.73	16.54	16.05	
154	34:50.04	−5:02:57.3	17.62	16.15	15.56	
155	34:50.11	−4:58:18.6	17.16	16.21	15.95	
156	34:50.15	−4:56:42.5	>21.50	>20.55	16.55	
157	34:50.15	−5:16:06.3	18.08	16.58	15.66	
158	34:50.27	−4:59:24.3	16.13	15.20	14.97	0534502−045924
159	34:50.32	−5:15:04.7	17.00	15.84	15.23	
160	34:50.51	−5:11:10.5	11.80	11.21	11.09†	0534505−051110
161	34:50.56	−5:06:38.3	13.97	13.46	13.18	0534505−050638
162	34:50.66	−5:04:07.8	12.66	11.88	11.60	0534506−050407
163	34:50.73	−4:58:37.0	12.25	11.21	10.81†	0534507−045836
164	34:51.15	−4:55:48.4	17.44	15.28	14.48	
165	34:51.15	−5:12:32.2	13.73	12.80	12.41	0534511−051232
166	34:51.31	−4:56:13.5	>21.24	19.75	15.07	
167	34:51.41	−5:00:11.4	13.20	12.46	12.20	0534514−050011
168	34:51.43	−5:13:29.5	12.86	12.19	11.94	0534514−051329
169	34:51.43	−5:14:40.3	19.83	17.94	16.40	
170	34:51.73	−5:08:18.8	16.44	15.34	15.01	0534517−050818
171	34:51.77	−4:55:55.4	15.88	13.67	12.68	0534517−045555
172	34:51.82	−4:58:28.9	>21.76	>20.81	17.72	
173	34:52.08	−4:58:24.1	17.44	16.24	15.82	
174	34:52.26	−5:12:03.2	13.10	12.36	12.17	0534522−051203
175	34:52.40	−5:14:19.6	15.96	14.96	14.67	
176	34:52.61	−5:15:36.5	13.04	11.74	11.18	0534526−051536
177	34:52.69	−5:10:32.7	16.37	15.19	14.92	0534526−051032
178	34:52.76	−5:00:50.8	13.10	12.39	12.16	0534527−050050
179	34:52.93	−5:15:19.5	16.47	15.34	14.64	
180	34:53.05	−5:03:27.0	12.01	11.40	10.73†	0534530−050327
181	34:53.13	−4:56:49.4	>21.49	>20.54	16.66	
182	34:53.43	−5:10:27.8	12.14	11.51	11.29	0534534−051027
183	34:53.44	−4:57:32.8	13.38	12.52	12.29	0534534−045732
184	34:53.53	−5:11:52.2	15.73	14.47	13.91	
185	34:53.56	−5:02:55.1	18.17	16.64	16.21	
186	34:53.56	−5:03:47.5	17.97	16.68	16.06	
187	34:53.63	−5:01:29.0	14.32	13.71	13.38	0534536−050129
188	34:53.64	−5:02:02.6	14.91	13.84	13.55	0534536−050202
189	34:53.65	−5:04:47.3	17.23	15.87	15.15	

(*cont.*)

ID	R.A.[a]	decl.[a]	J^{bcd}	H^{bcd}	K^{cd}	2MASS
190	34:53.71	−4:59:13.8	17.92	16.74	16.23	
191	34:53.73	−5:05:48.8	13.25	12.61	12.39	0534537−050548
192	34:53.80	−5:12:24.2	16.58	15.98	15.45	
193	34:53.87	−4:56:22.4	12.43	11.42	10.99†	0534538−045622
194	34:53.89	−5:07:55.4	16.98	15.72	15.13	
195	34:53.96	−4:56:15.7	17.30	16.14	15.52	
196	34:54.16	−5:11:21.9	18.74	16.88	16.42	
197	34:54.58	−4:56:05.4	12.17	11.15	10.54†	0534545−045605
198	34:54.82	−5:10:12.6	16.77	16.42	15.40	
199	34:54.90	−5:02:43.6	17.07	16.02	15.16	
200	34:54.92	−5:01:28.4	18.00	16.94	16.22	
201	34:54.93	−5:02:21.2	12.13	11.52	11.25†	0534549−050220
202	34:55.00	−4:55:37.6	>21.40	17.64	16.16	
203	34:55.11	−5:17:45.5	>21.00	15.65	14.57	
204	34:55.23	−5:15:58.4	18.36	15.43	14.73	
205	34:55.32	−4:56:59.6	14.37	13.70	13.49	0534553−045659
206	34:55.35	−5:09:25.5	16.80	15.62	15.08	
207	34:55.38	−5:13:51.9	17.57	16.27	16.04	
208	34:55.40	−5:16:00.9	17.33	14.97	14.29	
209	34:55.42	−5:01:39.5	12.51	11.71	11.55	0534554−050139
210	34:55.49	−5:11:15.1	17.97	16.48	15.95	
211	34:55.49	−5:12:59.3	18.24	16.79	16.14	
212	34:55.61	−5:12:34.9	19.06	16.89	16.45	
213	34:55.64	−5:06:30.3	18.14	16.82	15.94	
214	34:55.67	−4:56:12.0	12.39	11.10	10.26†	0534556−045611
215	34:55.81	−4:58:17.1	18.87	17.40	16.07	
216	34:55.86	−5:07:30.0	16.89	15.81	15.31	0534558−050730
217	34:55.87	−5:08:21.2	12.84	11.96	11.74	0534558−050821
218	34:55.91	−5:15:47.8	17.79	16.53	16.54	
219	34:56.10	−5:00:55.2	12.73	11.94	11.74	0534561−050055
220	34:56.13	−5:06:01.8	12.53	11.93	11.64	0534561−050601
221	34:56.17	−5:09:25.4	15.00	14.08	13.76	0534561−050925
222	34:56.18	−5:05:06.6	18.32	17.51	16.73	
223	34:56.36	−5:08:04.8	17.90	16.21	15.85	
224	34:56.38	−5:15:28.3	18.30	15.76	15.13	
225	34:56.54	−5:01:07.2	14.35	13.54	13.10	0534565−050107
226	34:56.63	−5:11:12.4	13.56	12.53	12.04	0534566−051112
227	34:56.66	−5:11:03.3	17.17	15.50	15.00	
228	34:56.67	−5:14:52.3	14.94	13.83	13.25	0534566−051452

(*cont.*)

ID	R.A.[a]	decl.[a]	J^{bcd}	H^{bcd}	K^{cd}	2MASS
229	34:56.72	−5:13:56.9	15.95	14.89	14.35	
230	34:56.77	−5:09:58.5	14.10	12.99	12.58	0534567−050958
231	34:56.78	−5:10:43.8	13.46	12.57	12.31	0534567−051043
232	34:56.82	−5:11:33.0	10.92†	10.19†	9.69†	0534568−051132
233	34:56.95	−4:54:36.0	18.24	15.22	13.64	
234	34:57.07	−5:00:55.9	18.23	16.87	16.20	
235	34:57.16	−5:01:35.5	17.63	16.47	16.15	
236	34:57.18	−5:07:16.8	16.06	14.95	14.55	0534571−050716
237	34:57.20	−5:08:23.9	12.08	11.36	11.01	0534571−050823
238	34:57.31	−5:02:50.8	18.42	16.50	15.68	
239	34:57.32	−5:04:00.2	18.59	16.98	16.58	
240	34:57.38	−5:14:33.5	13.44	12.30	11.81	0534573−051433
241	34:57.46	−4:56:45.5	14.40	12.51	11.63	0534574−045645
242	34:57.61	−4:54:33.3	>21.31	17.55	16.12	
243	34:57.63	−5:06:25.4	18.46	16.91	15.81	
244	34:57.68	−5:12:29.7	14.35	12.89	12.34	0534576−051229
245	34:57.71	−4:57:20.0	17.98	14.82	12.83	
246	34:57.89	−5:08:42.7	16.78	15.85	15.67	0534579−050842
247	34:58.03	−5:06:25.9	17.84	16.33	15.51	
248	34:58.03	−5:16:01.9	14.60	13.09	12.29	0534580−051601
249	34:58.05	−5:17:37.5	12.61	11.48	11.08	0534580−051737
250	34:58.15	−5:07:12.4	16.88	15.66	15.18	
251	34:58.19	−5:09:27.5	12.26	11.51	11.16	0534581−050927
252	34:58.19	−5:11:53.7	12.61	11.98	11.71	0534581−051153
253	34:58.21	−5:00:13.2	17.50	16.38	16.37	
254	34:58.24	−4:54:04.3	17.28	14.71	13.71	0534582−045404
255	34:58.49	−5:07:48.2	15.88	14.84	14.39	0534584−050748
256	34:58.69	−5:12:21.6	16.74	16.10	15.76	
257	34:58.69	−5:12:26.1	14.04	13.44	13.14	0534586−051226
258	34:58.83	−5:15:35.7	17.70	16.75	15.95	
259	34:58.84	−5:17:01.1	15.45	14.57	13.53	
260	34:58.93	−5:13:45.4	13.72	12.43	11.86	0534589−051345
261	34:58.93	−5:16:13.3	16.56	15.57	14.35	
262	34:59.08	−4:57:34.0	18.61	16.82	16.43	
263	34:59.21	−5:17:39.9	16.81	15.72	15.21	
264	34:59.31	−5:05:30.0	12.37	11.63	11.24	0534593−050530
265	34:59.32	−5:06:06.9	16.43	15.37	14.94	0534593−050606
266	34:59.43	−4:57:37.1	18.58	16.59	15.83	
267	34:59.67	−5:02:54.4	18.54	16.60	15.66	

(*cont.*)

ID	R.A.[a]	decl.[a]	J^{bcd}	H^{bcd}	K^{cd}	2MASS
268	34:59.88	−4:55:27.1	17.03	13.89	12.46	0534598−045527
269	34:59.99	−5:10:41.0	17.01	15.79	15.18	
270	35:00.19	−5:00:09.4	17.99	16.78	16.57	
271	35:00.24	−5:15:58.9	14.06	12.81	12.19	0535002−051558
272	35:00.25	−4:56:08.8	14.34	13.12	12.57	0535002−045608
273	35:00.34	−5:00:59.6	19.74	17.68	16.81	
274	35:00.40	−5:09:44.0	13.72	13.01	12.72	0535003−050944
275	35:00.41	−5:09:54.5	16.98	16.25	15.78	
276	35:00.44	−5:15:21.7	18.69	16.69	16.01	
277	35:00.64	−4:58:53.5	15.90	15.10	14.75	0535006−045853
278	35:00.67	−5:05:08.6	10.07†	9.95†	9.89†	0535006−050508
279	35:00.73	−5:11:27.6	16.60	14.95	14.48	
280	35:00.86	−5:09:39.1	12.99	12.28	12.07	0535008−050938
281	35:00.87	−5:09:16.0	17.51	16.59	16.63	
282	35:00.91	−5:05:40.1	15.84	15.25	14.94	0535009−050540
283	35:01.05	−5:05:33.5	15.50	14.90	14.55	0535010−050533
284	35:01.30	−5:06:32.9	17.06	16.08	15.75	
285	35:01.32	−4:55:57.0	14.34	13.25	12.82	0535013−045557
286	35:01.39	−4:56:32.6	16.47	15.62	15.11	0535013−045632
287	35:01.43	−5:09:32.7	12.17	11.48	11.24	0535014−050932
288	35:01.50	−5:11:04.1	13.77	12.69	12.18	0535014−051104
289	35:01.63	−4:56:46.2	18.78	16.41	15.23	
290	35:01.93	−5:00:43.3	14.50	13.72	13.52	0535019−050043
291	35:02.10	−5:15:37.5	12.69	11.93	11.77	0535020−051537
292	35:02.18	−5:04:38.4	13.77	12.99	12.65	0535021−050438
293	35:02.38	−5:15:47.9	9.32†	8.32†	7.91†	0535023−051547
294	35:02.64	−4:54:17.3	16.84	14.89	13.89	0535026−045417
295	35:02.64	−5:09:59.1	15.53	14.44	14.07	0535026−050959
296	35:02.66	−4:56:42.8	18.07	16.37	15.06	
297	35:02.67	−4:58:18.3	19.53	18.16	16.27	
298	35:02.73	−5:08:53.9	18.38	16.93	16.09	
299	35:02.75	−5:00:03.0	12.24	11.18	10.64†	0535027−050002
300	35:02.89	−5:03:03.7	15.82	15.06	14.76	0535028−050303
301	35:02.90	−4:54:30.2	20.83	16.27	13.64	
302	35:02.99	−4:57:04.0	14.05	13.33	13.16	0535029−045703
303	35:03.04	−4:59:59.9	12.57	11.85	11.66	0535030−045959
304	35:03.07	−5:13:54.4	13.32	11.84	11.07	0535030−051354
305	35:03.12	−5:09:17.1	13.14	12.28	12.05	0535031−050917
306	35:03.24	−5:17:53.1	12.89	12.44	11.51	0535032−051753

(*cont.*)

ID	R.A.[a]	decl.[a]	J^{bcd}	H^{bcd}	K^{cd}	2MASS
307	35:03.25	−5:17:26.2	15.11	13.82	12.80	
308	35:03.34	−5:16:22.7	14.02	12.38	11.56	0535033−051622
309	35:03.35	−4:56:43.0	12.36	11.59	11.34	0535033−045643
310	35:03.37	−5:08:53.8	13.41	12.09	11.66	0535033−050853
311	35:03.40	−5:05:40.3	9.14†	8.95†	8.90†	0535034−050540
312	35:03.40	−5:16:45.1	19.92	18.68	16.13	
313	35:03.59	−5:16:00.4	14.38	13.40	12.87	0535035−051600
314	35:03.65	−5:13:10.4	15.17	13.39	12.60	0535036−051310
315	35:03.85	−5:07:49.0	16.96	15.83	15.52	
316	35:04.08	−5:17:42.8	17.18	15.47	14.68	
317	35:04.20	−5:15:21.4	13.70	12.64	12.35	0535042−051521
318	35:04.29	−5:13:13.9	17.78	17.15	15.34	
319	35:04.30	−5:08:12.6	8.19†	7.52†	7.34†	0535042−050812
320	35:04.38	−5:01:15.4	16.24	14.48	13.68	0535043−050115
321	35:04.38	−5:08:16.7	<10.05	<10.01	<9.56	
322	35:04.40	−4:57:15.4	14.73	13.80	13.43	0535044−045715
323	35:04.40	−5:07:35.6	14.81	13.67	13.09	
324	35:04.60	−4:56:45.4	>21.87	>20.93	16.99	
325	35:04.62	−4:58:29.0	13.43	12.10	11.01†	0535046−045829
326	35:04.63	−5:09:55.7	11.32	10.74†	10.45†	0535046−050955
327	35:04.64	−4:54:02.5	15.85	14.30	13.56	0535046−045402
328	35:04.70	−5:01:50.9	17.50	16.04	16.11	
329	35:04.78	−5:17:41.9	11.49	10.11†	9.34†	0535047−051742
330	35:05.10	−4:53:51.4	18.10	16.05	14.91	
331	35:05.20	−5:14:50.2	8.23†	7.71†	7.38†	0535051−051450
332	35:05.21	−5:03:39.8	19.48	17.22	16.26	
333	35:05.45	−4:57:05.5	20.56	16.77	14.57	
334	35:05.48	−4:54:34.6	19.85	18.29	16.83	
335	35:05.61	−5:11:50.5	12.90	11.46	10.86†	0535056−051150
336	35:05.68	−4:58:53.9	11.99	11.30	11.04	0535056−045853
337	35:05.68	−5:02:40.1	17.54	15.96	15.27	
338	35:05.70	−4:58:33.4	15.95	13.51	12.46	0535056−045833
339	35:05.74	−5:11:35.0	12.94	11.46	10.83†	0535057−051135
340	35:05.78	−4:55:36.2	19.76	17.02	15.38	
341	35:05.88	−5:08:38.2	15.06	14.00	13.46	0535058−050838
342	35:06.06	−5:03:19.0	17.19	15.81	15.23	
343	35:06.10	−5:14:24.9	13.28	12.14	11.50	0535061−051425
344	35:06.19	−5:12:15.8	8.09†	8.22†	8.21†	0535061−051216
345	35:06.34	−4:58:41.6	14.08	13.46	13.27	0535063−045841

(*cont.*)

ID	R.A.[a]	decl.[a]	J^{bcd}	H^{bcd}	K^{cd}	2MASS
346	35:06.37	−5:03:02.5	16.23	14.80	14.28	0535063−050302
347	35:06.50	−5:17:22.2	14.39	12.66	11.67	
348	35:06.72	−5:11:45.3	13.91	12.43	11.75	0535067−051145
349	35:06.73	−5:16:46.4	16.89	15.31	14.13	
350	35:06.83	−5:10:38.4	12.56	11.81	11.61	0535068−051038
351	35:06.86	−5:11:50.0	15.84	14.91	14.23	
352	35:06.88	−5:11:33.0	17.78	15.82	14.77	
353	35:06.97	−4:57:04.2	18.48	15.25	13.59	
354	35:07.02	−5:01:34.8	17.68	16.86	16.16	
355	35:07.04	−4:54:56.7	12.80	11.95	11.84	0535070−045456
356	35:07.13	−4:55:50.7	14.66	13.18	12.61	0535071−045550
357	35:07.16	−4:55:43.2	19.38	17.23	15.15	
358	35:07.40	−5:07:38.0	17.37	15.56	14.62	
359	35:07.54	−5:11:14.5	12.31	11.32	10.70†	0535075−051114
360	35:07.75	−5:04:55.4	17.43	15.87	15.64	
361	35:07.79	−5:01:19.3	16.23	15.27	14.80	0535077−050119
362	35:07.83	−5:02:43.7	17.37	15.91	15.64	
363	35:07.85	−5:10:06.7	16.79	15.72	15.21	
364	35:07.88	−5:17:05.4	13.99	12.98	12.14	
365	35:08.23	−4:54:10.2	14.57	11.91	10.28†	0535082−045410
366	35:08.37	−5:16:20.5	14.19	13.21	12.49	
367	35:08.38	−5:16:25.0	>21.73	>20.78	17.45	
368	35:08.45	−5:07:13.5	14.37	13.13	12.71	0535084−050713
369	35:08.48	−5:03:48.5	18.74	16.64	15.92	
370	35:08.60	−5:09:30.3	17.51	15.57	14.78	
371	35:08.64	−5:16:47.6	14.83	13.64	12.52	
372	35:08.67	−5:16:13.3	16.30	13.73	11.79	
373	35:08.68	−4:59:33.2	19.66	17.98	16.78	
374	35:08.69	−4:54:41.2	17.32	15.82	14.83	
375	35:08.69	−5:17:30.6	>20.98	16.59	14.18	
376	35:08.75	−5:04:40.8	12.81	11.61	11.44	0535087−050440
377	35:08.83	−5:17:05.9	15.54	13.94	12.59	
378	35:08.95	−5:05:52.5	9.05†	7.91†	7.41†	0535089−050552
379	35:08.96	−5:10:25.8	14.54	13.83	13.74	0535089−051025
380	35:09.09	−5:10:09.7	17.39	15.81	15.31	
381	35:09.14	−5:16:44.2	16.43	14.36	12.63	
382	35:09.15	−5:06:47.1	11.24	10.57†	10.33†	0535091−050647
383	35:09.19	−5:16:51.7	14.62	13.31	12.14	
384	35:09.24	−5:10:22.6	18.20	16.59	16.13	

(*cont.*)

ID	R.A.^{a}	decl.^{a}	J^{bcd}	H^{bcd}	K^{cd}	2MASS
385	35:09.28	−5:16:55.7	13.82	12.76	11.96	
386	35:09.38	−4:55:13.6	>21.53	>20.58	16.21	
387	35:09.45	−5:16:53.3	14.19	13.14	12.15	
388	35:09.46	−4:57:11.7	13.92	13.11	12.51	0535094−045711
389	35:09.49	−4:59:41.4	12.94	12.24	11.85	0535094−045941
390	35:09.54	−5:14:45.8	>21.13	15.94	14.78	
391	35:09.82	−4:53:52.6	17.07	14.95	13.66	
392	35:09.83	−4:56:19.4	19.33	17.22	15.26	
393	35:09.88	−4:58:49.5	16.45	13.22	11.51	
394	35:09.96	−5:14:50.1	14.42	13.50	12.99	0535099−051450
395	35:10.01	−5:14:57.3	16.57	15.74	15.00	
396	35:10.05	−5:16:22.9	17.50	15.35	14.58	
397	35:10.09	−5:17:06.8	14.68	12.98	11.84	
398	35:10.11	−5:16:16.9	17.08	15.06	14.06	
399	35:10.28	−5:03:05.0	16.85	14.99	14.16	
400	35:10.30	−5:17:19.9	21.61	16.60	14.72	
401	35:10.38	−4:54:41.2	15.93	15.26	15.31	0535103−045441
402	35:10.53	−4:58:46.0	>21.08	13.84	<10.45	
403	35:10.73	−5:08:16.8	16.11	14.32	13.14	
404	35:10.79	−5:10:34.4	12.46	11.11	10.41†	0535107−051034
405	35:10.83	−5:17:33.2	15.09	13.40	11.99	
406	35:10.85	−4:55:42.6	17.33	15.75	14.36	
407	35:10.97	−4:56:39.7	16.85	14.71	13.33	0535109−045639
408	35:11.00	−5:15:21.8	15.99	13.39	11.13	
409	35:11.11	−5:16:01.8	16.30	14.38	13.16	
410	35:11.12	−4:54:14.9	20.13	17.13	16.00	
411	35:11.16	−4:55:28.0	18.16	15.46	14.00	
412	35:11.18	−4:55:38.1	17.86	15.89	14.60	
413	35:11.23	−5:17:20.8	12.30	<10.55	<9.22	
414	35:11.24	−5:17:42.3	16.28	14.89	13.96	
415	35:11.34	−5:14:01.6	13.80	12.80	12.37	0535113−051401
416	35:11.41	−5:17:46.5	15.12	13.23	11.72	
417	35:11.45	−5:05:16.2	16.20	14.94	14.28	0535114−050516
418	35:11.50	−5:17:57.3	15.42	12.56	<10.52	
419	35:11.63	−5:16:57.6	9.56†	8.94†	8.70†	0535116−051657
420	35:11.76	−5:16:52.0	14.16	13.20	12.99	
421	35:11.79	−4:54:21.5	16.06	14.37	13.56	0535118−045421
422	35:11.81	−5:03:30.2	14.54	13.07	12.51	0535118−050330
423	35:11.87	−5:01:56.0	17.30	15.91	14.93	

(*cont.*)

ID	R.A.^{a}	decl.^{a}	J^{bcd}	H^{bcd}	K^{cd}	2MASS
424	35:11.87	−5:17:25.9	13.55	11.89	11.81†	0535118−051725
425	35:11.89	−5:00:00.9	>21.66	>20.71	16.69	
426	35:12.05	−5:14:14.7	13.93	13.47	13.41	0535120−051414
427	35:12.13	−5:03:48.4	18.26	16.44	15.59	
428	35:12.17	−5:14:30.1	17.18	15.93	15.10	
429	35:12.30	−5:04:26.4	18.51	16.27	14.99	
430	35:12.31	−5:17:50.0	15.74	13.99	12.49	
431	35:12.51	−5:16:26.3	>21.04	15.79	13.45	
432	35:12.52	−4:57:13.7	15.07	13.13	12.35	0535125−045713
433	35:12.55	−4:56:01.9	17.49	14.20	12.28	
434	35:12.57	−5:16:33.2	13.87	11.40	9.90†	0535125−051633
435	35:12.69	−4:54:02.6	16.76	14.12	12.24	
436	35:12.69	−4:56:00.7	17.45	14.20	12.24	
437	35:12.69	−5:02:26.5	18.55	16.65	15.69	
438	35:12.70	−5:12:00.7	14.17	13.53	13.20	0535127−051200
439	35:12.70	−5:12:28.8	16.43	14.95	13.78	
440	35:12.72	−5:16:13.4	12.60	11.91	11.67	0535127−051613
441	35:12.73	−5:15:43.2	14.67	14.05	13.33	
442	35:12.74	−5:16:52.7	11.37	10.18†	9.62†	0535127−051652
443	35:12.81	−5:01:46.3	13.92	13.27	13.01	0535128−050146
444	35:12.82	−5:15:23.9	11.65	10.72†	10.39†	0535128−051524
445	35:12.87	−5:15:48.7	14.82	14.35	13.71	
446	35:12.95	−5:02:08.7	14.42	13.76	13.42	0535129−050208
447	35:12.97	−5:07:47.5	>21.67	19.58	17.14	
448	35:13.03	−5:00:26.1	>20.99	16.97	14.43	
449	35:13.10	−4:55:52.3	12.98	11.00	9.29†	0535131−045552
450	35:13.11	−5:15:26.3	13.27	12.60	12.27	
451	35:13.19	−5:17:30.6	15.22	12.86	<10.88	
452	35:13.32	−5:09:19.5	12.06	11.35	11.19	0535133−050919
453	35:13.41	−4:54:21.2	18.16	17.42	17.41	
454	35:13.47	−5:17:10.5	11.49	10.39†	9.93†	0535134−051710
455	35:13.48	−5:05:51.3	19.09	17.51	16.51	
456	35:13.52	−5:17:31.4	13.14	11.45	11.16†	0535135−051731
457	35:13.53	−5:17:17.6	13.24	12.09	11.61	0535135−051717
458	35:13.59	−5:17:45.9	17.34	14.23	11.40	
459	35:13.64	−5:14:22.2	15.74	15.02	15.15	0535136−051421
460	35:13.86	−4:58:03.5	15.58	13.12	11.93	0535138−045803
461	35:13.92	−4:53:58.4	>22.11	>21.16	16.97	
462	35:13.92	−5:18:53.2	13.57	12.89	11.55†	0535139−051853

(*cont.*)

ID	R.A.[a]	decl.[a]	J^{bcd}	H^{bcd}	K^{cd}	2MASS
463	35:13.96	−5:18:44.9	17.30	14.14	12.54	
464	35:14.04	−5:18:48.9	14.74	12.62	11.41	
465	35:14.06	−5:04:06.1	20.01	17.14	16.13	
466	35:14.19	−5:02:01.7	18.49	16.75	16.20	
467	35:14.22	−5:14:27.8	14.86	12.95	11.82	0535142−051427
468	35:14.29	−5:13:40.7	16.51	14.52	13.22	0535143−051340
469	35:14.29	−5:18:21.2	15.92	13.89	12.65	
470	35:14.38	−4:55:22.0	12.86	11.75	11.28	0535143−045522
471	35:14.46	−4:55:24.7	13.08	11.87	11.28	
472	35:14.47	−5:17:25.6	12.36	11.19†	10.73†	0535144−051725
473	35:14.47	−5:17:40.5	17.03	16.74	17.11	
474	35:14.47	−5:18:24.8	15.56	12.67	11.40	
475	35:14.64	−5:02:25.0	12.82	11.76	11.26	
476	35:14.64	−5:16:46.3	13.04	12.18	11.70	0535146−051646
477	35:14.65	−5:06:25.3	14.38	13.25	12.72	0535146−050625
478	35:14.68	−5:03:12.7	12.59	11.74	11.59	0535146−050312
479	35:14.68	−5:08:52.1	12.80	12.21	12.06	0535146−050852
480	35:14.69	−5:06:22.3	16.54	15.36	14.75	
481	35:14.70	−5:16:52.4	>21.78	>20.83	15.54	
482	35:14.70	−5:18:43.3	16.54	12.58	11.31	
483	35:14.73	−5:02:27.9	13.14	11.86	11.26	
484	35:14.74	−5:14:00.5	15.62	13.84	12.81	0535147−051400
485	35:14.77	−5:14:15.1	>21.52	17.51	16.32	
486	35:14.82	−5:18:30.6	16.91	13.79	12.84	
487	35:14.86	−5:06:48.9	16.90	14.54	13.39	
488	35:14.88	−5:17:53.7	>21.93	17.99	17.32	
489	35:14.89	−5:07:47.7	13.91	13.20	12.77	0535148−050747
490	35:14.91	−5:03:21.5	18.22	17.33	16.67	
491	35:14.97	−4:56:08.3	>20.86	15.86	13.61	
492	35:14.99	−4:56:04.6	>21.05	15.12	13.22	
493	35:15.00	−5:16:49.6	17.22	17.26	15.28	
494	35:15.05	−5:00:08.2	17.21	15.86	14.10	
495	35:15.05	−5:16:39.5	18.01	14.75	13.19	
496	35:15.07	−5:06:53.8	15.83	13.59	12.47	0535150−050653
497	35:15.08	−4:58:08.6	17.57	15.96	14.84	
498	35:15.26	−5:00:33.4	13.45	12.73	12.52	0535152−050033
499	35:15.29	−5:15:48.2	14.39	13.62	13.06	
500	35:15.33	−5:13:15.5	>21.70	>20.75	16.21	
501	35:15.33	−5:13:38.2	13.49	12.61	12.46	

(*cont.*)

ID	R.A.[a]	decl.[a]	J^{bcd}	H^{bcd}	K^{cd}	2MASS
502	35:15.34	−4:57:46.1	16.76	14.82	13.82	
503	35:15.36	−5:02:50.6	17.91	15.53	14.16	
504	35:15.37	−5:19:02.1	12.68	11.36	11.07	
505	35:15.44	−5:16:00.3	17.74	16.47	14.78	
506	35:15.46	−5:17:39.1	12.36	<10.73	<9.51	
507	35:15.47	−5:00:11.4	18.15	17.32	15.60	
508	35:15.47	−5:01:12.5	14.47	13.70	13.41	0535154−050112
509	35:15.53	−5:17:36.5	12.37	<10.83	<9.66	
510	35:15.54	−5:08:58.9	15.96	15.32	15.19	0535155−050858
511	35:15.55	−5:01:43.8	13.73	12.56	12.05	0535155−050143
512	35:15.57	−5:07:28.6	19.08	18.06	16.34	
513	35:15.59	−5:09:51.2	13.41	12.23	11.60	0535155−050951
514	35:15.60	−4:59:27.9	14.60	12.70	11.47	0535156−045927
515	35:15.61	−5:05:50.6	17.46	14.93	13.11	
516	35:15.62	−5:09:31.9	13.93	13.35	13.11	0535156−050931
517	35:15.64	−4:57:13.9	13.30	12.40	12.15	0535156−045713
518	35:15.65	−5:06:12.5	16.51	15.62	13.29	
519	35:15.69	−5:17:47.3	12.63	<10.65	<9.50	
520	35:15.78	−5:15:27.6	20.48	18.42	17.29	
521	35:15.79	−5:16:42.3	18.60	17.62	15.42	
522	35:15.80	−5:12:26.4	18.38	15.20	13.88	
523	35:15.82	−5:03:26.1	13.84	13.11	12.72	0535158−050326
524	35:15.84	−5:00:36.3	12.27	11.80	11.75	0535158−050036
525	35:15.86	−4:57:11.3	13.99	13.13	12.77	0535158−045711
526	35:15.89	−5:06:54.3	18.88	17.06	16.37	
527	35:15.93	−5:06:13.6	16.80	15.91	13.75	
528	35:15.96	−5:14:59.1	11.58	11.15	<10.72	
529	35:15.97	−5:16:57.6	12.57	11.55	11.07	
530	35:16.14	−5:15:48.4	14.76	14.40	14.12	
531	35:16.17	−5:00:02.5	16.16	13.57	11.77	
532	35:16.17	−5:09:19.1	12.94	12.27	12.11	0535161−050919
533	35:16.19	−5:14:13.3	16.81	14.88	13.80	
534	35:16.20	−5:19:02.9	15.75	13.18	11.98	
535	35:16.21	−5:19:06.8	15.89	13.35	12.37	
536	35:16.29	−4:57:36.0	19.85	16.59	15.31	
537	35:16.31	−5:17:44.3	13.96	12.21	11.03	
538	35:16.34	−5:15:38.1	12.03	11.27	<10.74	
539	35:16.35	−5:04:36.4	14.85	14.13	13.82	0535163−050436
540	35:16.37	−5:18:09.0	15.50	12.70	11.49	

(*cont.*)

ID	R.A.[a]	decl.[a]	J^{bcd}	H^{bcd}	K^{cd}	2MASS
541	35:16.41	−4:58:02.0	13.25	12.54	12.35	0535164−045802
542	35:16.42	−4:56:53.1	19.40	17.61	16.45	
543	35:16.44	−4:58:07.0	13.41	12.59	12.34	0535164−045807
544	35:16.48	−5:06:00.6	15.32	14.60	12.86	
545	35:16.49	−4:57:17.4	18.77	16.53	15.41	
546	35:16.51	−5:03:30.4	12.42	11.72	11.35	0535164−050330
547	35:16.51	−5:17:47.3	14.55	12.73	11.65	
548	35:16.61	−5:16:20.6	15.72	14.19	13.36	
549	35:16.63	−5:17:23.5	13.12	11.10	<9.77	
550	35:16.74	−5:18:44.8	15.29	13.38	11.64	
551	35:16.78	−5:17:16.9	13.43	12.47	11.91	
552	35:16.79	−5:07:27.3	14.89	14.41	14.01	0535167−050727
553	35:16.80	−5:14:47.4	15.61	14.58	13.92	
554	35:16.80	−5:16:53.3	13.60	12.79	12.42	
555	35:16.80	−5:19:01.1	13.95	11.40	<10.03	
556	35:16.86	−5:07:47.8	13.77	12.81	12.32	0535168−050747
557	35:16.89	−5:15:09.2	16.05	14.15	12.96	
558	35:16.91	−5:17:03.0	13.40	11.73	<10.80	
559	35:16.94	−5:18:41.0	15.50	12.81	11.38	
560	35:17.01	−5:17:32.1	14.92	13.67	13.34	
561	35:17.02	−5:15:44.3	12.50	11.64	11.01	
562	35:17.07	−5:17:29.0	15.03	13.79	13.87	
563	35:17.12	−5:19:00.7	14.91	12.38	<10.87	
564	35:17.13	−5:18:13.7	>21.68	>20.73	14.59	
565	35:17.14	−4:58:06.1	16.48	14.91	14.06	0535171−045806
566	35:17.15	−4:57:47.1	15.14	14.47	14.17	0535171−045747
567	35:17.15	−5:12:39.4	13.51	12.89	12.78	
568	35:17.15	−5:18:06.5	>21.67	18.80	13.47	
569	35:17.38	−5:12:29.6	13.70	13.08	12.93	
570	35:17.42	−4:59:57.0	17.47	14.64	12.89	
571	35:17.42	−5:05:00.9	>21.57	19.73	17.19	
572	35:17.43	−5:17:12.5	16.56	15.63	15.02	
573	35:17.45	−5:16:57.0	14.42	13.61	13.12	
574	35:17.50	−5:09:49.1	13.34	12.70	12.47	0535174−050949
575	35:17.50	−5:16:50.0	19.02	18.79	16.76	
576	35:17.51	−5:18:22.5	16.02	12.73	<10.90	
577	35:17.54	−5:17:40.1	<10.79	<9.46	<8.53	
578	35:17.56	−5:16:13.0	12.56	11.22	<10.60	
579	35:17.57	−5:19:28.9	13.53	11.07	<10.16	

(*cont.*)

ID	R.A.[a]	decl.[a]	J^{bcd}	H^{bcd}	K^{cd}	2MASS
580	35:17.61	−5:18:32.6	15.01	12.13	<10.92	
581	35:17.62	−5:16:41.1	16.84	16.39	16.31	
582	35:17.67	−5:16:55.2	17.10	16.17	15.64	
583	35:17.71	−5:03:08.9	20.58	17.33	15.93	
584	35:17.73	−5:00:30.9	15.07	12.86	11.55	0535177−050031
585	35:17.79	−5:16:15.7	11.94	<10.65	<9.97	
586	35:17.82	−5:19:28.1	14.11	11.75	<10.82	
587	35:17.83	−4:55:17.2	16.98	14.51	13.14	
588	35:17.91	−5:17:56.0	14.61	12.86	11.84	
589	35:17.91	−5:18:35.2	13.54	<10.59	<9.54	
590	35:17.92	−5:15:32.9	11.50	<10.76	<10.18	
591	35:17.98	−4:57:59.7	18.33	16.27	15.40	
592	35:17.98	−5:16:45.1	13.07	11.70	11.04	
593	35:17.99	−5:15:38.7	13.06	12.39	11.86	
594	35:18.01	−5:16:13.6	11.18	<10.00	<9.37	
595	35:18.11	−5:15:46.1	13.37	12.69	12.39	
596	35:18.19	−5:17:53.9	15.01	13.54	13.04	
597	35:18.21	−5:03:54.5	9.26†	9.26†	9.21†	0535182−050354
598	35:18.21	−5:16:33.9	12.76	11.97	11.70	
599	35:18.21	−5:17:22.0	15.69	13.41	11.57	
600	35:18.22	−5:15:06.1	17.15	16.69	16.08	
601	35:18.24	−5:13:06.9	11.36	<10.34	<9.78	
602	35:18.24	−5:17:44.9	13.68	11.61	<10.12	
603	35:18.30	−5:08:04.8	17.02	14.58	13.24	
604	35:18.33	−5:00:32.9	18.02	14.57	<10.72	
605	35:18.35	−5:13:16.7	14.53	13.30	12.67	
606	35:18.37	−5:19:17.5	15.32	13.72	12.26	
607	35:18.38	−5:15:01.5	>21.86	>20.91	16.27	
608	35:18.44	−5:07:14.5	>21.28	17.44	15.94	
609	35:18.45	−5:06:45.9	19.93	18.01	15.45	
610	35:18.46	−5:16:37.7	<10.36	<9.89	<9.57	
611	35:18.47	−5:08:30.7	13.79	13.13	12.97	0535184−050830
612	35:18.52	−5:13:38.4	11.34	<10.40	<9.61	
613	35:18.55	−5:18:20.5	12.80	11.36	11.10	
614	35:18.60	−4:55:11.0	17.52	15.58	14.09	
615	35:18.60	−4:59:42.3	13.86	13.22	12.96	0535186−045942
616	35:18.60	−5:16:35.1	<10.54	<9.97	<9.57	
617	35:18.62	−5:06:47.3	20.70	16.47	14.15	
618	35:18.62	−5:13:27.5	16.73	14.53	13.25	

(*cont.*)

ID	R.A.[a]	decl.[a]	J^{bcd}	H^{bcd}	K^{cd}	2MASS
619	35:18.64	−5:18:35.1	14.10	12.05	11.26	
620	35:18.69	−5:17:43.1	16.17	14.42	13.27	
621	35:18.71	−5:15:18.1	14.30	13.67	12.97	
622	35:18.75	−5:18:02.2	12.86	<10.95	<10.00	
623	35:18.82	−5:17:29.0	11.37	<9.68	<8.35	
624	35:18.83	−5:17:20.9	16.09	14.46	13.11	
625	35:18.86	−5:14:45.6	12.48	11.65	11.27	
626	35:18.90	−5:19:02.9	15.06	12.91	11.58	
627	35:18.92	−5:16:13.9	11.66	11.36	11.27	
628	35:18.95	−5:12:31.3	17.89	16.33	15.64	
629	35:18.97	−5:17:51.6	16.50	14.60	13.01	
630	35:19.10	−4:54:07.8	13.70	13.03	12.97	0535191−045407
631	35:19.20	−5:18:39.1	15.89	14.38	13.63	
632	35:19.25	−5:16:58.6	15.89	14.59	13.71	
633	35:19.25	−5:17:31.6	15.58	13.24	11.43	
634	35:19.28	−5:19:11.9	16.94	13.47	12.02	
635	35:19.31	−5:16:44.6	12.87	12.03	11.46	0535193−051644
636	35:19.32	−4:55:44.9	15.67	12.87	10.25†	0535193−045545
637	35:19.32	−5:16:10.0	14.48	13.97	13.66	
638	35:19.33	−5:19:03.7	15.49	13.24	12.12	
639	35:19.38	−5:17:19.2	14.59	13.53	12.51	
640	35:19.48	−5:16:08.2	13.93	13.26	12.77	
641	35:19.52	−5:08:58.1	19.28	18.18	15.59	
642	35:19.58	−5:17:03.1	13.55	11.31	10.06†	0535195−051703
643	35:19.66	−5:02:28.8	14.54	13.82	13.50	0535196−050228
644	35:19.66	−5:13:26.5	13.61	11.55	10.24†	0535196−051326
645	35:19.67	−5:03:34.5	18.72	16.57	15.30	
646	35:19.74	−5:19:30.1	>21.03	13.23	11.27	
647	35:19.75	−5:04:54.3	18.13	>20.17	13.94	
648	35:19.80	−4:57:06.7	>21.06	16.83	14.86	
649	35:19.80	−5:14:05.4	13.52	13.10	12.68	0535197−051405
650	35:19.82	−5:15:35.4	17.65	14.95	12.35	
651	35:19.84	−5:15:09.0	13.00	11.43	10.26†	0535198−051508
652	35:19.98	−5:01:02.4	>21.05	18.28	14.20	
653	35:20.00	−5:18:46.6	13.91	12.04	12.09	
654	35:20.03	−5:12:50.5	13.87	11.96	11.17	0535200−051250
655	35:20.12	−4:58:40.2	19.20	16.74	16.25	
656	35:20.20	−5:15:59.5	14.55	12.94	12.26	
657	35:20.24	−5:13:16.2	9.89†	7.86†	6.49†	0535201−051315

(*cont.*)

ID	R.A.[a]	decl.[a]	J^{bcd}	H^{bcd}	K^{cd}	2MASS
658	35:20.25	−5:12:11.2	14.74	14.04	14.12	
659	35:20.25	−5:13:59.6	12.24	10.75†	9.88†	0535202−051359
660	35:20.36	−5:02:26.4	12.39	11.57	11.28	0535203−050226
661	35:20.42	−5:17:14.6	11.99	11.21†	10.43†	0535204−051714
662	35:20.44	−5:01:15.1	18.43	17.32	15.98	
663	35:20.45	−5:06:38.0	>21.18	17.87	15.62	
664	35:20.62	−5:03:00.7	18.44	14.47	11.94	
665	35:20.67	−5:09:02.9	14.19	13.17	12.70	0535206−050902
666	35:20.68	−5:01:54.2	13.16	12.35	12.20	0535206−050154
667	35:20.74	−5:19:26.6	15.93	14.66	11.78	
668	35:20.76	−5:15:49.4	9.90†	8.68†	8.08†	0535207−051549
669	35:20.77	−4:58:34.0	12.39	11.34	10.72†	0535207−045834
670	35:20.78	−5:13:23.4	15.11	14.03	13.40	
671	35:20.82	−4:57:16.9	16.28	13.91	12.74	0535208−045717
672	35:20.83	−5:02:57.7	19.49	16.28	13.60	
673	35:20.86	−4:56:53.0	>21.30	18.48	15.16	
674	35:20.99	−5:16:37.7	12.02	11.36	11.19	0535210−051637
675	35:21.07	−5:01:16.6	16.87	16.34	14.12	
676	35:21.09	−5:19:16.3	>21.78	>20.83	15.56	
677	35:21.15	−5:06:32.4	>21.23	18.56	15.31	
678	35:21.17	−5:18:21.4	12.85	11.97	11.57	0535211−051821
679	35:21.22	−5:06:47.8	>21.15	>20.20	15.39	
680	35:21.28	−5:09:16.2	8.38†	7.84†	7.31†	0535212−050916
681	35:21.32	−4:58:35.0	15.41	14.13	13.62	
682	35:21.32	−5:12:12.8	8.71†	8.11†	7.90†	0535213−051212
683	35:21.38	−5:12:44.4	15.91	15.20	14.63	
684	35:21.41	−5:09:42.3	13.74	12.89	12.57	0535213−050942
685	35:21.45	−4:57:18.8	20.51	17.02	15.77	
686	35:21.47	−5:09:03.7	12.27	11.32	<10.91	
687	35:21.47	−5:17:11.0	>21.46	>20.51	14.56	
688	35:21.50	−5:01:15.9	15.48	14.82	12.37	
689	35:21.52	−5:01:53.9	18.49	15.06	13.06	
690	35:21.59	−5:09:38.9	13.75	12.85	12.56	0535215−050938
691	35:21.60	−5:09:49.7	13.01	12.03	11.75	0535215−050949
692	35:21.63	−5:15:17.2	17.42	14.81	13.55	
693	35:21.65	−5:17:19.4	13.26	11.82	10.78†	0535216−051718
694	35:21.67	−5:17:17.0	13.22	11.83	<10.83	
695	35:21.69	−4:56:48.7	13.65	12.94	12.80	
696	35:21.74	−5:17:40.2	>21.20	>20.25	14.16	

(*cont.*)

ID	R.A.^{a}	decl.^{a}	J^{bcd}	H^{bcd}	K^{cd}	2MASS
697	35:21.76	−5:06:32.8	>21.18	17.30	14.77	
698	35:21.76	−5:17:27.8	>21.32	>20.37	15.58	
699	35:21.77	−5:09:45.3	19.37	17.65	16.23	
700	35:21.85	−4:55:21.8	20.53	19.61	17.31	
701	35:21.88	−4:54:07.6	11.91	11.31	<10.68	
702	35:21.88	−4:59:19.7	15.35	13.86	13.01	
703	35:21.88	−5:07:01.9	11.27	10.00†	9.40†	0535218−050701
704	35:21.90	−5:18:12.0	17.22	16.66	14.62	
705	35:21.91	−5:15:01.2	12.78	11.85	11.49	0535219−051501
706	35:21.94	−4:57:09.1	18.35	15.70	14.35	
707	35:21.94	−5:08:03.7	13.05	12.03	11.58	0535219−050803
708	35:21.94	−5:17:04.4	12.19	11.81	11.65	0535219−051704
709	35:21.95	−5:14:27.7	11.79	11.36	11.31	0535219−051427
710	35:22.08	−5:00:14.1	20.75	17.25	16.13	
711	35:22.08	−5:12:22.6	17.36	16.69	15.15	
712	35:22.13	−5:18:57.6	>21.67	>20.72	14.17	
713	35:22.23	−4:53:57.3	16.59	15.50	15.44	
714	35:22.25	−5:00:39.0	18.22	17.03	14.99	
715	35:22.26	−5:18:08.8	13.20	12.49	11.98	0535222−051808
716	35:22.37	−5:07:39.2	14.47	12.05	10.46†	0535223−050739
717	35:22.37	−5:08:17.9	14.67	13.72	13.52	
718	35:22.39	−5:17:32.9	14.55	12.98	12.01	
719	35:22.42	−5:08:05.2	10.97†	9.51†	8.59†	0535224−050805
720	35:22.46	−5:09:11.1	11.44	11.12	10.62†	0535224−050911
721	35:22.57	−5:08:00.7	11.26†	9.78†	9.21†	0535225−050800
722	35:22.58	−4:59:05.3	17.46	16.02	15.39	
723	35:22.60	−5:13:28.4	14.33	13.64	13.37	0535226−051328
724	35:22.63	−5:14:11.3	11.69	11.19	10.24†	0535226−051411
725	35:22.67	−5:15:08.6	13.38	12.63	12.35	0535226−051508
726	35:22.71	−5:16:14.0	13.21	12.19	11.82	0535226−051613
727	35:22.75	−5:18:38.1	>21.25	16.67	12.80	
728	35:22.89	−4:59:09.4	14.61	13.69	13.12	
729	35:22.93	−5:13:40.0	14.10	12.52	11.66	0535229−051339
730	35:22.99	−5:15:21.8	16.43	16.44	13.73	
731	35:23.01	−5:17:45.2	>21.07	14.04	11.63	
732	35:23.03	−5:14:39.5	>21.75	20.73	16.14	
733	35:23.07	−5:11:50.3	15.88	14.27	13.20	
734	35:23.08	−5:00:36.4	16.48	13.89	12.45	
735	35:23.13	−5:13:43.6	13.02	11.33	10.49†	0535231−051343

(*cont.*)

ID	R.A.^{a}	decl.^{a}	J^{bcd}	H^{bcd}	K^{cd}	2MASS
736	35:23.22	−5:08:43.6	>21.11	14.94	13.10	
737	35:23.24	−5:00:38.2	16.56	13.90	12.44	
738	35:23.27	−5:13:48.3	15.35	14.29	13.53	
739	35:23.27	−5:15:43.4	15.68	14.41	13.39	
740	35:23.30	−4:56:11.9	16.82	14.30	12.95	
741	35:23.30	−4:57:20.7	12.90	12.05	11.92	
742	35:23.30	−5:03:23.4	17.01	14.35	12.85	
743	35:23.33	−5:07:09.6	17.79	14.78	12.51	
744	35:23.35	−5:08:21.6	16.80	14.26	12.41	
745	35:23.40	−5:12:03.1	15.96	15.03	12.47	
746	35:23.41	−5:18:50.6	11.32	10.54†	10.23†	0535234−051850
747	35:23.42	−5:01:28.7	>21.42	>20.47	16.17	
748	35:23.47	−5:10:51.7	11.17	11.06	10.40†	0535234−051051
749	35:23.48	−5:15:23.4	14.05	13.38	12.93	
750	35:23.54	−5:18:57.0	11.89	11.14	11.18†	0535235−051857
751	35:23.59	−5:19:11.9	14.82	13.71	12.60	
752	35:23.60	−5:19:30.2	14.97	13.83	12.28	
753	35:23.65	−5:01:40.3	>21.02	18.18	14.94	
754	35:23.77	−5:18:39.8	12.53	11.62	10.90†	0535237−051839
755	35:23.82	−4:59:24.9	18.49	16.70	15.75	
756	35:23.89	−5:18:20.9	17.02	13.95	12.20	
757	35:23.92	−5:04:11.6	16.88	15.33	14.48	
758	35:23.92	−5:06:01.7	>20.88	17.96	14.38	
759	35:23.96	−5:07:53.4	18.29	16.64	15.26	
760	35:23.99	−5:19:07.4	12.92	12.14	11.80	0535239−051907
761	35:24.00	−5:18:38.6	12.59	11.86	11.20	
762	35:24.06	−4:56:30.4	15.69	15.06	14.82	
763	35:24.06	−5:19:23.3	14.72	13.35	11.96	
764	35:24.09	−5:09:06.7	14.50	13.90	13.51	0535240−050906
765	35:24.14	−5:18:54.3	17.72	15.13	13.41	
766	35:24.14	−5:19:27.6	15.23	13.55	12.01	
767	35:24.34	−5:01:20.5	>21.10	19.52	14.68	
768	35:24.38	−5:14:57.9	16.57	14.66	13.58	0535243−051457
769	35:24.49	−5:16:59.8	11.65	10.75†	10.18†	0535244−051659
770	35:24.55	−5:00:21.1	18.00	15.88	14.07	
771	35:24.61	−5:11:58.3	11.29	10.96†	10.41†	0535246−051158
772	35:24.63	−5:11:29.5	12.20	11.30	10.18†	0535246−051129
773	35:24.63	−5:19:33.2	16.14	13.21	11.47	
774	35:24.64	−5:06:25.9	16.95	15.16	13.89	

(*cont.*)

ID	R.A.^{a}	decl.^{a}	J^{bcd}	H^{bcd}	K^{cd}	2MASS
775	35:24.66	−5:19:09.5	11.86	10.99†	10.29†	0535246−051909
776	35:24.69	−5:09:26.4	17.01	15.11	13.55	
777	35:24.72	−5:16:41.1	13.85	11.68	<10.38	
778	35:24.74	−5:06:56.2	>21.27	>20.32	14.23	
779	35:24.79	−5:10:29.5	<10.34	<8.79	<7.16	
780	35:24.88	−5:06:21.5	16.08	12.93	11.07	0535248−050621
781	35:24.94	−5:12:55.9	16.92	16.23	13.87	
782	35:25.03	−5:09:09.4	14.60	13.96	13.49	0535250−050909
783	35:25.03	−5:19:25.2	15.21	13.95	12.76	
784	35:25.15	−5:06:30.7	>21.86	18.45	16.58	
785	35:25.24	−5:15:35.7	11.72	10.81†	10.09†	0535252−051535
786	35:25.25	−5:09:27.5	11.82	11.25	11.03	0535252−050927
787	35:25.34	−5:12:05.7	14.09	13.19	12.93	0535253−051205
788	35:25.38	−5:16:36.1	16.39	16.17	15.55	
789	35:25.42	−5:10:48.0	11.62	11.30	11.15	0535254−051048
790	35:25.44	−4:54:04.1	13.72	13.01	12.89	0535254−045403
791	35:25.56	−4:55:27.0	14.90	12.98	12.33	0535255−045526
792	35:25.64	−5:07:57.1	11.82†	11.15†	10.22†	0535256−050757
793	35:25.66	−5:09:41.9	18.17	15.94	13.91	
794	35:25.66	−5:18:04.2	13.96	12.94	12.31	
795	35:25.67	−4:57:18.3	13.33	12.59	12.40	0535256−045718
796	35:25.68	−5:00:28.4	18.50	16.77	15.98	
797	35:25.70	−5:07:03.0	13.74	13.14	12.95	0535257−050703
798	35:25.73	−5:07:46.2	14.74	12.36	11.22	0535257−050746
799	35:25.75	−5:09:49.3	11.39	10.48†	10.11†	0535257−050949
800	35:25.76	−5:05:57.9	16.37	14.53	12.66	
801	35:25.87	−5:07:56.4	11.14	11.00†	10.09†	0535258−050756
802	35:26.02	−5:19:12.9	>22.41	>21.46	18.22	
803	35:26.06	−5:08:37.7	12.61†	12.02†	10.03†	0535260−050837
804	35:26.16	−5:08:33.4	14.55	13.78	13.39	
805	35:26.29	−5:08:39.9	10.10†	9.60†	9.35†	0535262−050840
806	35:26.31	−5:15:11.4	10.25†	9.85†	9.80†	0535263−051511
807	35:26.34	−5:17:44.7	15.16	14.58	13.75	
808	35:26.36	−5:10:50.4	14.76	14.04	13.74	0535263−051050
809	35:26.38	−5:18:57.6	16.24	14.02	12.77	
810	35:26.41	−5:16:12.5	13.31	11.91	11.15	0535264−051612
811	35:26.41	−5:16:37.9	14.79	13.87	13.00	
812	35:26.44	−5:00:56.9	14.85	14.12	13.88	0535264−050057
813	35:26.45	−5:15:05.4	16.74	15.63	15.14	

(*cont.*)

ID	R.A.^{a}	decl.^{a}	J^{bcd}	H^{bcd}	K^{cd}	2MASS
814	35:26.46	−4:59:51.9	16.69	14.21	12.89	
815	35:26.49	−5:17:56.6	12.20	<10.86	<10.21	
816	35:26.51	−5:19:18.7	>21.18	>20.23	12.61	
817	35:26.58	−4:58:25.3	20.52	16.46	14.51	
818	35:26.58	−5:17:53.1	12.11	12.38†	11.36†	0535265−051753
819	35:26.58	−5:19:08.3	18.17	>19.94	13.97	
820	35:26.59	−4:56:06.7	11.85	11.32	10.93†	0535265−045606
821	35:26.62	−5:03:54.6	14.89	12.12	<9.96	
822	35:26.63	−5:18:18.6	16.06	16.08	14.53	
823	35:26.64	−4:55:02.5	16.78	15.56	15.03	0535266−045502
824	35:26.73	−5:16:45.2	13.11	12.28	11.87	0535267−051645
825	35:26.75	−5:14:44.5	15.07	14.39	14.19	0535267−051444
826	35:26.75	−5:19:02.8	16.15	14.06	12.70	
827	35:26.87	−5:09:24.6	15.18	13.21	11.17	0535268−050924
828	35:26.87	−5:17:13.1	16.50	15.52	14.37	
829	35:26.94	−5:04:07.0	15.90	14.22	13.11	
830	35:26.94	−5:18:06.9	13.84	12.63	11.88	
831	35:26.96	−5:11:07.7	8.92†	7.84†	7.03†	0535268−051107
832	35:26.96	−5:15:37.2	>21.73	17.64	16.10	
833	35:26.97	−5:10:17.4	13.23	10.93†	8.73†	0535269−051017
834	35:27.00	−5:13:14.5	10.55†	10.26†	10.17†	0535269−051314
835	35:27.01	−5:09:54.4	14.83	13.41	12.55	0535269−050954
836	35:27.06	−5:15:44.8	12.83	11.91	11.25	0535270−051544
837	35:27.18	−4:58:34.5	18.91	16.61	15.59	
838	35:27.21	−4:55:18.6	14.05	12.74	12.29	0535272−045518
839	35:27.27	−5:05:27.2	17.52	16.54	15.61	
840	35:27.44	−5:02:42.4	16.87	13.34	11.04	
841	35:27.44	−5:09:04.0	15.09	14.35	13.80	0535274−050903
842	35:27.45	−5:17:12.1	11.21	<10.36	<9.78	
843	35:27.47	−5:17:09.8	11.20	10.24†	9.53†	0535274−051709
844	35:27.49	−5:09:44.3	11.66	10.51†	9.83†	0535274−050944
845	35:27.49	−5:11:50.1	16.27	15.16	13.45	
846	35:27.53	−5:13:56.4	13.71	12.46	11.60	0535275−051356
847	35:27.53	−5:16:55.9	14.09	12.79	12.26	
848	35:27.55	−5:18:19.3	14.72	13.06	12.17	
849	35:27.56	−5:10:08.6	16.73	14.27	12.36	
850	35:27.59	−5:05:31.2	19.96	19.55	14.99	
851	35:27.64	−4:54:38.9	19.09	17.33	16.73	
852	35:27.64	−5:09:32.9	13.65	<10.76	<9.14	

(*cont.*)

ID	R.A.[a]	decl.[a]	J^{bcd}	H^{bcd}	K^{cd}	2MASS
853	35:27.65	−5:18:00.7	13.15	12.08	11.64	
854	35:27.66	−5:09:37.3	13.25	10.51†	8.94†	0535276−050937
855	35:27.67	−5:09:49.9	12.74	11.74	<10.60	
856	35:27.75	−5:18:04.7	11.98	10.55†	9.94†	0535277−051804
857	35:27.81	−5:07:54.8	14.20	13.28	12.75	0535277−050754
858	35:27.86	−5:05:36.3	>20.99	16.26	13.37	
859	35:27.93	−5:16:57.3	11.60	10.86†	10.56†	0535279−051657
860	35:27.94	−4:53:52.1	17.69	16.71	17.04	
861	35:27.97	−5:18:59.8	13.90	11.53	<10.25	
862	35:28.01	−5:19:14.0	>21.03	15.78	14.36	
863	35:28.05	−5:05:38.4	>21.13	17.06	14.49	
864	35:28.05	−5:17:20.3	12.45	11.69	11.43	0535280−051720
865	35:28.06	−5:01:35.1	19.33	14.97	12.18	
866	35:28.14	−5:18:57.0	13.36	11.03	9.68†	0535281−051857
867	35:28.15	−5:10:13.9	11.03	10.19†	9.48†	0535281−051013
868	35:28.17	−5:00:49.8	12.53	11.74	11.51	0535281−050049
869	35:28.17	−5:15:50.9	14.06	12.54	11.78	
870	35:28.19	−5:03:41.3	20.45	16.61	11.84	
871	35:28.19	−5:11:37.6	12.96	11.87	11.55	0535281−051137
872	35:28.19	−5:16:01.4	14.70	14.16	13.15	
873	35:28.26	−4:58:38.4	>21.06	14.14	11.39	
874	35:28.30	−4:55:42.1	17.33	15.79	14.98	
875	35:28.36	−5:17:54.4	13.52	12.36	11.93	0535283−051754
876	35:28.36	−5:18:23.0	12.34	11.64	11.46	0535283−051823
877	35:28.40	−4:58:07.2	14.65	13.88	13.67	0535283−045807
878	35:28.41	−5:07:44.1	16.22	13.07	11.11	
879	35:28.43	−5:19:02.0	13.23	12.59	12.22	0535284−051902
880	35:28.44	−4:53:57.0	17.46	16.42	15.95	
881	35:28.46	−4:57:16.9	>21.66	18.34	16.67	
882	35:28.52	−5:07:46.9	16.23	13.05	11.09	
883	35:28.60	−4:55:03.8	11.05	10.54†	10.38†	0535286−045503
884	35:28.60	−5:05:44.6	15.04	12.50	10.81†	0535285−050544
885	35:28.66	−5:02:44.8	15.49	13.83	12.99	0535286−050244
886	35:28.71	−5:19:25.7	18.58	15.86	15.62	
887	35:28.87	−4:57:39.3	18.00	14.98	13.11	
888	35:28.88	−4:54:36.1	18.68	16.39	15.28	
889	35:28.93	−4:54:20.6	14.93	13.91	13.34	0535289−045420
890	35:28.94	−5:16:18.4	11.22	10.37†	9.63†	0535289−051618
891	35:29.04	−5:06:04.0	11.93	11.41	11.12	0535290−050604

(*cont.*)

ID	R.A.[a]	decl.[a]	J^{bcd}	H^{bcd}	K^{cd}	2MASS
892	35:29.14	−5:00:18.2	19.47	16.58	15.32	
893	35:29.17	−5:18:18.3	14.67	13.71	13.09	
894	35:29.37	−5:02:14.8	17.10	15.44	14.63	
895	35:29.37	−5:11:46.9	>22.20	>21.25	17.51	
896	35:29.44	−4:55:45.6	19.09	17.69	16.73	
897	35:29.45	−5:16:33.4	11.52	10.73†	10.44†	0535294−051633
898	35:29.45	−5:17:55.4	14.22	13.20	12.26	
899	35:29.46	−5:18:45.6	13.94	12.26	10.09†	0535294−051845
900	35:29.51	−5:00:00.5	>21.44	20.19	16.25	
901	35:29.54	−5:17:47.1	>22.05	17.75	18.06	
902	35:29.55	−5:18:40.0	14.99	13.46	12.71	
903	35:29.56	−4:59:56.7	18.00	17.21	16.07	
904	35:29.79	−4:56:11.7	>21.50	>20.55	16.53	
905	35:29.82	−5:16:06.3	11.98	11.26	10.94†	0535298−051606
906	35:29.90	−5:18:53.0	16.14	13.70	12.19	
907	35:29.91	−5:12:10.3	12.64	12.08	11.85	0535299−051210
908	35:29.93	−4:57:08.2	14.25	12.35	11.60	0535299−045708
909	35:29.93	−5:08:20.2	>20.63	15.57	13.12	
910	35:30.00	−5:12:27.4	12.06	11.43	10.86†	0535299−051227
911	35:30.07	−5:19:06.4	17.64	16.37	15.14	
912	35:30.11	−5:09:09.5	13.75	13.24	13.20	0535301−050909
913	35:30.15	−5:14:18.5	13.03	12.00	11.35	0535301−051418
914	35:30.22	−5:16:57.5	16.83	15.90	15.04	
915	35:30.23	−5:08:19.1	>20.74	15.79	13.52	
916	35:30.26	−5:09:32.3	17.07	16.03	14.21	
917	35:30.28	−4:55:53.9	18.00	16.27	15.42	
918	35:30.29	−5:13:52.6	13.29	11.66	10.70†	0535302−051352
919	35:30.36	−5:18:05.6	14.26	13.36	13.09	
920	35:30.46	−5:19:00.7	15.10	13.39	12.33	
921	35:30.48	−5:19:33.9	13.94	12.03	<10.40	
922	35:30.53	−5:17:15.2	14.01	12.68	12.15	
923	35:30.54	−5:03:34.5	15.42	13.67	12.75	0535305−050334
924	35:30.59	−4:57:21.5	14.77	14.01	13.74	0535305−045721
925	35:30.61	−4:59:36.0	14.36	11.65	10.48†	0535306−045936
926	35:30.63	−5:15:16.2	13.65	13.11	12.88	0535306−051516
927	35:30.65	−5:04:11.1	18.32	16.47	15.02	
928	35:30.70	−5:18:07.0	12.38	11.40	10.98†	0535307−051807
929	35:30.72	−5:03:35.5	15.35	13.61	12.75	0535307−050335
930	35:30.73	−4:55:49.8	18.66	16.92	16.51	

(*cont.*)

ID	R.A.^{a}	decl.^{a}	J^{bcd}	H^{bcd}	K^{cd}	2MASS
931	35:30.80	−4:54:43.7	16.89	15.23	14.54	
932	35:30.80	−4:58:13.6	17.83	14.50	12.78	
933	35:30.89	−4:55:17.9	11.87	11.36	10.68†	0535308−045517
934	35:30.91	−5:18:17.8	11.54	10.80†	10.46†	0535309−051817
935	35:30.97	−5:04:44.1	18.65	16.77	15.46	
936	35:31.03	−4:57:20.0	12.27	11.50	11.28	0535310−045720
937	35:31.04	−5:18:07.6	16.56	15.84	16.24	
938	35:31.04	−5:18:44.9	13.26	12.37	11.87	0535310−051845
939	35:31.07	−5:04:15.0	11.51	11.01	10.41†	0535310−050415
940	35:31.13	−5:15:13.9	15.79	15.33	15.62	
941	35:31.14	−4:54:15.2	9.15†	9.20†	9.22†	0535311−045415
942	35:31.14	−5:13:43.8	16.40	14.14	13.03	0535311−051343
943	35:31.20	−4:57:27.1	14.92	13.00	12.03	0535312−045727
944	35:31.21	−5:12:27.9	12.94	12.33	12.16	0535312−051228
945	35:31.22	−5:19:31.8	14.28	13.21	12.63	
946	35:31.26	−5:15:10.7	15.14	14.12	13.83	
947	35:31.27	−5:18:55.7	10.80†	10.07†	9.58†	0535312−051855
948	35:31.29	−5:15:32.9	10.07†	9.28†	8.93†	0535312−051533
949	35:31.33	−5:12:01.5	14.53	13.95	13.61	0535313−051201
950	35:31.44	−5:16:03.4	5.84†	5.64†	5.56†	0535313−051602
951	35:31.48	−4:57:47.7	16.88	15.07	14.14	0535314−045747
952	35:31.51	−5:05:01.6	12.16	11.57	10.77†	0535315−050501
953	35:31.54	−5:05:47.2	11.83	9.77†	8.28†	0535315−050547
954	35:31.54	−5:14:44.7	>21.33	15.97	15.55	
955	35:31.55	−5:16:36.8	12.59	11.58	10.85†	0535315−051636
956	35:31.58	−5:15:23.4	13.00	12.07	11.74	0535315−051523
957	35:31.59	−4:54:37.4	14.88	13.40	12.74	0535316−045437
958	35:31.59	−5:06:25.0	17.22	16.52	14.95	
959	35:31.61	−5:19:27.5	>21.29	14.48	13.09	
960	35:31.62	−5:16:58.1	12.98	11.72	11.25	0535316−051658
961	35:31.64	−5:00:14.0	14.11	11.53	9.96†	0535316−050014
962	35:31.64	−5:03:46.1	18.52	16.24	15.38	
963	35:31.75	−5:16:39.8	12.80	11.86	11.32	0535317−051639
964	35:31.76	−5:14:52.2	11.38	10.03†	9.47†	0535317−051452
965	35:31.97	−5:09:27.9	9.37†	8.54†	8.09†	0535319−050927
966	35:31.98	−5:15:59.5	11.64	11.03	<10.68	
967	35:32.00	−5:16:20.0	12.56	12.58	11.56	0535319−051620
968	35:32.01	−4:55:35.7	13.94	13.24	13.02	0535320−045535
969	35:32.02	−5:08:05.8	14.14	12.90	12.76	0535320−050805

(*cont.*)

ID	R.A.^{a}	decl.^{a}	J^{bcd}	H^{bcd}	K^{cd}	2MASS
970	35:32.03	−5:08:34.9	17.13	15.86	15.32	
971	35:32.04	−4:56:42.4	16.08	14.79	14.24	0535320−045642
972	35:32.10	−5:00:23.7	15.57	13.91	13.06	0535321−050023
973	35:32.19	−5:11:57.8	12.63	11.92	11.54	0535321−051157
974	35:32.20	−5:16:03.6	12.52	12.04	11.70	
975	35:32.31	−5:16:26.8	14.00	12.92	12.51	
976	35:32.34	−5:11:43.3	14.92	12.67	11.57	0535323−051143
977	35:32.34	−5:18:07.7	11.53	10.67†	10.30†	0535323−051807
978	35:32.36	−4:58:29.9	17.18	14.54	13.10	
979	35:32.38	−5:12:53.0	16.61	16.22	16.39	
980	35:32.39	−5:12:10.8	19.48	18.40	17.13	
981	35:32.43	−5:14:24.6	12.68	11.66	11.43	0535324−051424
982	35:32.44	−5:15:06.7	13.61	12.25	11.45	0535324−051506
983	35:32.53	−5:01:57.2	15.39	13.60	12.65	0535325−050157
984	35:32.54	−5:02:09.8	14.86	14.12	13.76	0535325−050209
985	35:32.62	−5:05:37.9	15.74	14.84	14.32	0535326−050538
986	35:32.64	−4:56:52.7	18.45	17.21	16.33	
987	35:32.64	−5:15:51.2	13.49	12.42	11.99	0535326−051551
988	35:32.72	−5:12:00.7	15.52	14.12	12.80	
989	35:32.81	−5:17:38.5	13.36	12.28	11.78	0535328−051738
990	35:32.84	−5:18:19.8	13.90	12.44	11.62	
991	35:32.90	−5:16:30.2	15.67	17.92	13.98	
992	35:32.92	−5:16:05.3	11.31	10.50†	10.16†	0535329−051605
993	35:32.93	−5:02:46.8	17.58	15.60	14.51	
994	35:32.95	−4:57:54.2	18.03	16.25	15.59	
995	35:32.96	−5:12:04.8	13.27	11.52	10.55†	0535329−051204
996	35:32.96	−5:16:40.1	14.52	13.98	13.07	
997	35:33.02	−5:17:39.2	14.51	13.26	12.54	
998	35:33.10	−5:13:38.7	>20.85	>19.90	13.92	
999	35:33.12	−5:17:34.0	12.84	12.17	11.85	0535331−051733
1000	35:33.17	−5:14:10.6	12.20	11.47	11.06	0535331−051410
1001	35:33.21	−5:16:05.4	11.45	10.74†	9.82†	0535331−051605
1002	35:33.34	−5:11:45.6	17.54	16.58	15.32	
1003	35:33.43	−4:56:54.5	17.67	16.44	16.02	
1004	35:33.44	−5:08:45.1	>21.31	20.12	15.20	
1005	35:33.45	−4:56:01.7	13.03	12.33	12.09	0535334−045601
1006	35:33.49	−5:05:00.9	19.23	15.91	14.19	
1007	35:33.58	−4:55:17.9	18.29	17.84	16.25	
1008	35:33.59	−5:15:23.2	13.13	11.84	11.42	0535335−051523

(*cont.*)

ID	R.A.^{a}	decl.^{a}	J^{bcd}	H^{bcd}	K^{cd}	2MASS
1009	35:33.63	−5:00:42.0	12.01	11.30	11.03	0535336−050042
1010	35:33.65	−5:03:08.0	13.29	11.77	10.89†	0535336−050308
1011	35:33.68	−4:57:20.9	18.10	16.72	16.10	
1012	35:33.80	−4:55:00.2	13.11	12.09	11.69	0535338−045500
1013	35:33.83	−5:04:27.5	13.82	11.82	10.53†	0535338−050427
1014	35:33.84	−5:17:10.8	14.40	13.27	12.82	0535338−051710
1015	35:33.88	−5:09:06.2	13.96	13.40	13.22	0535338−050905
1016	35:33.88	−5:15:54.4	14.27	12.96	12.32	0535338−051554
1017	35:33.92	−5:03:33.5	12.54	11.97	11.74	0535339−050333
1018	35:34.00	−4:53:57.6	13.48	12.08	11.61	0535340−045357
1019	35:34.03	−4:54:11.0	12.73	12.04	11.88	0535340−045411
1020	35:34.08	−5:17:09.5	15.36	14.47	14.26	
1021	35:34.29	−5:06:21.2	9.31†	9.33†	9.31†	0535342−050621
1022	35:34.42	−5:18:38.4	15.23	13.66	12.42	
1023	35:34.46	−5:03:30.7	17.73	16.63	16.21	
1024	35:34.50	−5:00:52.4	16.28	13.69	11.28	0535345−050052
1025	35:34.52	−5:13:07.7	14.32	13.74	13.58	0535345−051307
1026	35:34.58	−5:19:05.5	15.89	14.33	13.08	
1027	35:34.59	−4:57:52.8	18.56	16.93	15.64	
1028	35:34.62	−5:15:52.8	12.98	12.33	11.71	0535346−051552
1029	35:34.80	−5:15:31.2	14.13	12.94	12.29	0535348−051531
1030	35:34.98	−5:12:30.2	14.50	13.83	12.85	
1031	35:35.04	−5:11:09.2	14.16	13.40	12.91	0535350−051109
1032	35:35.20	−5:00:14.7	18.10	15.92	15.55	
1033	35:35.28	−4:58:02.2	18.40	16.61	15.68	
1034	35:35.34	−5:05:24.9	17.22	15.08	13.89	
1035	35:35.35	−5:11:11.5	12.61	12.05	11.88	0535353−051111
1036	35:35.40	−5:08:46.9	13.84	13.09	12.87	0535353−050846
1037	35:35.45	−5:07:53.1	16.31	14.28	13.07	0535354−050753
1038	35:35.55	−5:06:58.6	13.28	12.68	12.53	0535355−050658
1039	35:35.60	−5:15:43.2	11.66	10.96†	10.64†	0535356−051543
1040	35:35.68	−5:10:51.0	14.13	13.64	13.36	0535356−051050
1041	35:35.77	−4:56:27.6	16.93	15.51	15.05	
1042	35:35.79	−5:12:20.5	9.79†	9.84†	9.76†	0535357−051220
1043	35:35.93	−5:12:06.8	18.41	17.06	16.92	
1044	35:35.97	−5:11:01.0	19.93	17.39	16.26	
1045	35:35.98	−5:02:12.6	18.14	16.66	16.04	
1046	35:36.01	−5:12:25.2	11.20†	9.57†	8.47†	0535360−051225
1047	35:36.02	−4:57:41.6	14.47	13.73	13.47	0535360−045741

(*cont.*)

ID	R.A.^{a}	decl.^{a}	J^{bcd}	H^{bcd}	K^{cd}	2MASS
1048	35:36.05	−5:13:17.1	15.97	14.77	13.81	
1049	35:36.21	−5:04:55.8	13.26	11.76	11.12	0535362−050455
1050	35:36.26	−5:14:30.2	15.10	14.49	13.77	
1051	35:36.42	−5:01:15.5	10.93†	10.07†	9.56†	0535364−050115
1052	35:36.44	−5:09:11.5	16.63	15.89	15.41	
1053	35:36.48	−5:15:32.3	>21.51	16.71	15.55	
1054	35:36.57	−5:04:39.2	12.05	11.30	11.15	0535365−050439
1055	35:36.57	−5:12:31.0	15.77	15.32	14.56	
1056	35:36.58	−4:58:40.6	17.64	16.53	16.11	
1057	35:36.67	−5:18:50.7	14.14	13.27	12.80	0535366−051850
1058	35:36.68	−5:04:14.4	11.98	11.31	10.87†	0535366−050414
1059	35:36.69	−4:57:29.9	13.61	12.38	11.92	0535366−045730
1060	35:36.72	−4:56:44.7	18.74	16.79	15.43	
1061	35:36.74	−5:10:00.4	13.13	12.43	11.97	0535367−051000
1062	35:36.87	−4:58:05.0	13.73	12.21	11.64	0535368−045805
1063	35:36.89	−5:18:21.6	>21.55	18.84	14.64	
1064	35:36.95	−5:03:26.2	21.10	>20.73	16.43	
1065	35:36.95	−5:04:05.3	17.02	15.84	15.19	
1066	35:36.97	−5:05:26.2	14.58	14.06	13.70	0535369−050526
1067	35:37.06	−5:18:58.6	20.98	18.74	16.60	
1068	35:37.16	−5:10:29.5	12.47	11.81	11.64	0535371−051029
1069	35:37.18	−4:56:53.7	>21.64	>20.69	17.39	
1070	35:37.33	−5:02:36.3	15.27	13.16	12.10	0535373−050236
1071	35:37.45	−4:58:26.8	11.69	10.01†	9.19†	0535374−045826
1072	35:37.59	−5:04:47.1	17.66	14.88	13.22	
1073	35:37.70	−5:06:31.9	12.96	12.25	12.10	0535376−050631
1074	35:37.89	−5:03:46.2	17.42	15.77	15.20	0535378−050346
1075	35:38.01	−4:57:13.5	18.79	16.98	16.46	
1076	35:38.01	−5:03:29.8	17.46	16.05	15.64	
1077	35:38.03	−5:19:26.7	17.55	16.59	15.96	
1078	35:38.10	−5:18:11.1	11.53	9.78†	8.86†	0535381−051811
1079	35:38.13	−4:55:41.0	14.26	13.11	12.65	
1080	35:38.19	−4:56:37.0	>21.54	>20.59	16.73	
1081	35:38.19	−5:03:33.7	15.76	14.92	14.51	0535382−050333
1082	35:38.22	−5:14:18.8	11.58	11.01	10.56†	0535382−051418
1083	35:38.32	−4:58:01.2	16.18	14.81	14.31	
1084	35:38.45	−5:10:08.9	14.82	14.36	14.13	0535384−051009
1085	35:38.51	−5:04:51.4	16.33	13.69	11.95	0535385−050451
1086	35:38.53	−4:59:41.0	12.05	11.10	<10.44	

(*cont.*)

ID	R.A.[a]	decl.[a]	J^{bcd}	H^{bcd}	K^{cd}	2MASS
1087	35:38.57	−5:08:03.2	15.91	13.46	12.01	0535385−050803
1088	35:38.58	−5:07:31.1	17.86	16.89	16.76	
1089	35:38.65	−5:09:56.6	12.98	12.47	12.22	0535386−050956
1090	35:38.72	−5:16:58.9	13.44	12.21	11.53	0535387−051659
1091	35:38.75	−5:04:55.4	14.95	12.91	11.94	0535387−050455
1092	35:38.79	−5:11:19.9	17.78	16.48	15.82	
1093	35:38.82	−4:57:10.2	16.12	14.98	14.62	
1094	35:38.83	−5:12:15.3	15.40	14.86	14.84	0535388−051215
1095	35:38.86	−5:12:41.8	10.87†	9.72†	9.02†	0535388−051241
1096	35:38.98	−5:19:14.6	>21.93	17.56	16.32	
1097	35:39.03	−5:07:04.2	12.52	11.97	11.79	0535390−050704
1098	35:39.08	−5:08:56.3	10.87†	10.20†	10.00†	0535390−050856
1099	35:39.16	−5:12:20.2	13.72	12.74	12.27	0535391−051220
1100	35:39.20	−5:16:35.5	14.12	12.31	11.36	0535392−051635
1101	35:39.27	−5:13:50.9	18.14	17.41	15.06	
1102	35:39.29	−5:18:31.7	>21.45	>20.50	16.55	
1103	35:39.68	−5:12:08.2	>21.07	>20.12	15.03	
1104	35:39.73	−5:01:54.0	>21.09	15.62	14.18	
1105	35:39.84	−5:15:49.4	14.01	12.89	12.42	0535398−051549
1106	35:39.92	−4:57:31.3	<10.95	<10.32	<9.93	
1107	35:39.93	−4:58:39.1	12.82	11.99	11.82	
1108	35:39.97	−5:06:36.7	12.48	11.47	10.91†	0535399−050636
1109	35:40.01	−5:02:36.9	14.12	13.45	13.16	0535400−050236
1110	35:40.03	−4:57:28.9	11.06	<10.38	<9.97	
1111	35:40.04	−5:11:38.2	16.71	14.97	14.49	
1112	35:40.17	−5:09:56.0	18.34	15.51	14.88	
1113	35:40.17	−5:16:25.8	>20.88	14.62	14.13	
1114	35:40.19	−5:16:32.0	15.58	13.67	13.17	0535402−051631
1115	35:40.20	−5:17:29.1	10.35†	9.53†	9.01†	0535401−051729
1116	35:40.25	−5:15:47.5	>20.89	14.35	14.04	
1117	35:40.29	−5:13:36.8	15.97	14.15	13.65	
1118	35:40.32	−5:03:05.5	16.79	15.31	14.78	
1119	35:40.34	−5:12:31.9	15.97	14.79	14.37	0535403−051232
1120	35:40.38	−4:55:44.0	13.28	12.39	12.03	
1121	35:40.53	−4:59:12.8	15.76	15.08	14.84	
1122	35:40.56	−5:03:04.6	16.41	14.99	14.52	0535405−050304
1123	35:40.62	−5:12:19.3	12.95	12.21	11.99	0535406−051219
1124	35:40.63	−5:13:20.6	16.92	16.11	16.06	
1125	35:40.63	−5:19:02.6	>21.13	16.06	14.31	

(*cont.*)

ID	R.A.[a]	decl.[a]	J^{bcd}	H^{bcd}	K^{cd}	2MASS
1126	35:40.65	−5:19:33.0	15.35	15.56	14.81	
1127	35:40.68	−5:01:53.3	19.32	17.28	16.29	
1128	35:40.68	−5:03:56.1	14.79	13.36	12.71	0535406−050356
1129	35:40.77	−5:12:47.9	13.29	12.17	11.64	0535407−051247
1130	35:40.78	−5:09:01.6	10.36†	9.43†	9.09†	0535407−050901
1131	35:40.78	−5:11:11.9	13.67	13.09	12.74	0535407−051111
1132	35:40.80	−5:18:39.1	16.80	15.84	14.95	
1133	35:40.86	−5:18:32.6	16.50	15.81	15.29	
1134	35:40.96	−5:16:57.5	15.97	15.31	15.12	
1135	35:41.05	−5:06:25.3	12.55	11.89	11.67	0535410−050625
1136	35:41.10	−5:09:55.8	17.10	16.59	16.00	
1137	35:41.29	−5:03:52.2	14.17	12.92	12.43	0535413−050352
1138	35:41.35	−5:16:41.3	16.81	16.07	15.21	
1139	35:41.38	−5:04:38.8	17.04	15.70	14.93	
1140	35:41.38	−5:12:58.9	11.31	10.61†	10.38†	0535413−051258
1141	35:41.42	−5:16:44.9	16.88	16.09	15.40	
1142	35:41.57	−5:16:05.5	15.10	14.21	14.06	0535415−051605
1143	35:41.63	−4:56:53.9	12.64	11.87	11.85	0535416−045653
1144	35:41.74	−5:03:29.0	13.41	12.66	12.55	0535417−050329
1145	35:41.74	−5:05:19.8	15.37	13.77	12.82	0535417−050519
1146	35:41.77	−4:58:13.9	17.95	16.52	15.92	
1147	35:41.85	−5:01:25.6	15.61	14.45	14.08	0535418−050125
1148	35:41.99	−5:17:36.4	18.17	17.01	16.63	
1149	35:42.01	−5:10:11.5	11.46	10.16†	9.50†	0535420−051011
1150	35:42.04	−5:16:11.2	17.89	16.91	16.75	
1151	35:42.07	−5:12:59.4	13.37	12.81	12.51	0535420−051259
1152	35:42.16	−5:19:09.3	17.89	16.96	15.86	
1153	35:42.28	−5:15:59.3	16.86	15.44	14.91	
1154	35:42.29	−5:15:08.0	13.55	12.08	11.60	0535422−051507
1155	35:42.35	−5:11:57.4	15.19	14.73	14.54	
1156	35:42.50	−5:12:38.2	15.35	14.70	13.89	
1157	35:42.52	−4:59:40.1	12.58	11.83	11.79	0535425−045940
1158	35:42.61	−5:01:03.4	15.77	14.38	13.89	0535426−050103
1159	35:42.69	−5:10:16.3	14.64	14.13	13.73	0535426−051016
1160	35:42.69	−5:10:49.1	19.76	17.17	16.24	
1161	35:42.77	−5:11:54.7	11.20	10.47†	10.28†	0535427−051154
1162	35:42.97	−4:57:43.5	13.62	12.85	12.44	0535429−045743
1163	35:42.97	−5:19:05.8	18.82	16.97	16.16	
1164	35:43.06	−5:03:07.3	13.32	12.62	12.39	0535430−050307

(*cont.*)

ID	R.A.[a]	decl.[a]	J^{bcd}	H^{bcd}	K^{cd}	2MASS
1165	35:43.11	−5:13:46.8	10.19†	9.89†	9.78†	0535431−051346
1166	35:43.20	−4:58:23.3	13.76	13.03	12.78	0535431−045823
1167	35:43.25	−5:09:17.0	10.78†	10.08†	9.88†	0535432−050917
1168	35:43.55	−5:05:41.3	11.96	11.48	10.38†	0535435−050541
1169	35:43.55	−5:08:49.4	15.95	14.90	14.26	0535435−050849
1170	35:43.58	−5:09:33.9	16.28	15.26	14.92	0535435−050933
1171	35:43.63	−5:01:52.8	>21.29	16.89	15.99	
1172	35:43.64	−4:56:54.0	20.18	>20.76	16.23	
1173	35:43.65	−5:17:28.7	15.20	14.08	13.64	
1174	35:43.66	−5:07:07.5	17.82	16.87	16.09	
1175	35:43.66	−5:17:25.5	14.27	13.20	12.79	0535436−051725
1176	35:43.80	−5:14:39.1	14.37	13.47	13.16	0535437−051439
1177	35:43.81	−5:09:58.7	13.61	13.07	12.71	0535438−050958
1178	35:43.93	−5:14:04.8	18.45	17.19	16.53	
1179	35:43.96	−5:03:42.9	14.86	13.63	13.09	0535439−050343
1180	35:44.03	−4:56:18.5	15.56	14.54	13.90	0535440−045618
1181	35:44.09	−5:08:37.5	14.63	13.19	11.86	0535440−050837
1182	35:44.09	−5:12:56.5	17.36	16.33	16.22	
1183	35:44.35	−4:57:16.8	13.23	12.49	12.25	0535443−045716
1184	35:44.50	−5:07:31.6	12.07	11.60	11.30	0535445−050731
1185	35:44.53	−5:08:56.3	16.12	14.93	14.18	0535445−050856
1186	35:44.59	−4:56:24.8	16.81	15.58	15.37	0535446−045625
1187	35:44.61	−4:59:57.6	14.92	14.26	14.01	0535446−045957
1188	35:44.70	−5:00:39.6	14.00	13.13	12.89	0535447−050039
1189	35:44.71	−4:58:35.3	16.18	15.01	14.65	0535447−045835
1190	35:44.79	−4:58:12.4	18.62	16.83	15.90	
1191	35:44.86	−5:07:16.8	10.09†	9.45†	9.17†	0535448−050716
1192	35:44.93	−4:57:01.4	16.12	15.40	14.80	0535449−045701
1193	35:44.95	−5:15:20.1	14.43	13.39	13.10	0535449−051520
1194	35:44.99	−4:56:02.8	13.57	12.63	12.39	0535450−045602
1195	35:45.07	−5:13:55.4	16.70	15.34	15.36	
1196	35:45.12	−5:17:27.3	17.34	16.38	16.36	
1197	35:45.19	−5:00:47.6	17.01	15.85	15.53	
1198	35:45.24	−5:07:09.1	15.94	15.02	14.43	0535452−050709
1199	35:45.32	−5:19:08.3	15.69	13.82	12.92	0535453−051908
1200	35:45.38	−5:10:11.9	16.30	14.67	13.90	0535453−051012
1201	35:45.48	−5:09:45.9	>21.90	>20.95	16.95	
1202	35:45.58	−5:18:02.9	13.07	13.80	13.04	
1203	35:45.61	−5:18:13.3	12.55	11.52	11.37	0535456−051813

(*cont.*)

ID	R.A.[a]	decl.[a]	J^{bcd}	H^{bcd}	K^{cd}	2MASS
1204	35:45.62	−5:18:42.6	15.22	13.83	14.24	0535456−051842
1205	35:45.63	−5:13:36.0	13.30	12.25	11.67	0535456−051336
1206	35:45.70	−5:14:05.4	18.67	16.52	15.35	
1207	35:45.71	−5:10:55.3	13.54	13.03	12.68	0535457−051055
1208	35:45.97	−5:16:09.9	17.08	15.40	14.04	
1209	35:46.02	−4:55:31.9	15.16	13.98	13.78	0535460−045532
1210	35:46.07	−5:17:49.4	12.34	11.56	11.06	0535460−051749
1211	35:46.10	−5:12:24.3	18.08	16.30	14.77	
1212	35:46.11	−4:56:41.0	17.39	16.15	15.69	
1213	35:46.13	−5:10:51.8	12.02	11.23	10.62†	0535461−051051
1214	35:46.14	−5:00:43.4	16.68	15.45	15.16	0535461−050043
1215	35:46.20	−4:58:49.5	15.20	14.07	13.70	0535462−045849
1216	35:46.24	−5:15:39.8	12.64	11.53	10.98†	0535462−051539
1217	35:46.26	−5:16:25.4	>21.37	19.05	15.92	
1218	35:46.31	−5:13:24.3	16.52	15.40	15.04	
1219	35:46.48	−4:58:17.8	15.33	14.24	13.94	0535464−045817
1220	35:46.49	−4:57:54.8	16.59	15.51	15.29	0535465−045755
1221	35:46.56	−5:03:41.3	16.37	15.17	14.65	0535465−050341
1222	35:46.57	−5:05:33.4	18.25	16.39	15.24	
1223	35:46.62	−5:15:55.5	20.34	15.90	15.59	
1224	35:46.67	−5:17:29.0	16.52	14.43	14.33	
1225	35:46.82	−5:16:46.8	12.77	11.88	11.55	0535468−051646
1226	35:46.90	−5:10:43.3	14.12	13.22	12.89	0535469−051043
1227	35:47.01	−5:16:14.1	15.22	14.52	13.88	
1228	35:47.02	−5:06:13.2	16.22	15.31	14.96	0535470−050612
1229	35:47.09	−5:14:42.5	15.26	14.10	13.59	
1230	35:47.17	−5:00:44.9	16.16	15.39	15.13	0535471−050044
1231	35:47.20	−5:16:43.2	16.27	15.94	15.85	
1232	35:47.25	−5:17:43.1	12.15	11.64	11.33	
1233	35:47.39	−5:13:18.5	11.58	11.03†	10.84†	0535473−051318
1234	35:47.42	−5:10:28.5	11.17	10.88†	10.47†	0535474−051028
1235	35:47.47	−5:00:46.9	17.69	16.58	16.01	
1236	35:47.48	−5:14:30.4	17.51	16.25	15.35	
1237	35:47.50	−5:16:57.8	10.42†	10.06†	9.93†	0535475−051657
1238	35:47.54	−5:07:10.0	16.16	14.95	14.57	0535475−050709
1239	35:47.54	−5:12:18.1	12.88	12.23	12.21	0535475−051218
1240	35:47.79	−5:10:30.8	10.18†	9.64†	9.40†	0535477−051030
1241	35:47.83	−4:58:54.6	17.95	16.38	15.91	
1242	35:47.83	−5:13:17.7	13.29	12.08	11.70	0535478−051317

(*cont.*)

ID	R.A.^{a}	decl.^{a}	J^{bcd}	H^{bcd}	K^{cd}	2MASS
1243	35:48.01	−5:08:11.6	15.09	13.61	12.78	0535480−050811
1244	35:48.11	−5:15:24.7	17.39	16.42	15.92	
1245	35:48.26	−5:11:10.3	12.79	11.95	11.66	0535482−051110
1246	35:48.32	−4:57:42.7	15.50	14.25	13.83	0535483−045742
1247	35:48.39	−5:01:28.7	12.24	11.39	10.21†	0535483−050128
1248	35:48.39	−5:08:52.6	17.59	16.40	16.05	
1249	35:48.50	−5:15:35.7	17.07	16.18	16.14	
1250	35:48.55	−5:15:21.1	15.18	14.39	14.15	0535485−051521
1251	35:48.59	−5:17:42.6	14.56	13.49	13.09	0535485−051742
1252	35:48.62	−5:11:53.2	16.72	15.78	15.56	
1253	35:48.76	−4:58:38.5	18.01	16.70	16.50	
1254	35:48.84	−5:00:28.5	12.66	11.67	11.35	0535488−050028
1255	35:48.96	−5:09:53.0	16.15	15.11	14.60	0535489−050953
1256	35:48.99	−5:01:39.3	12.61	11.81	11.53	0535489−050139
1257	35:49.02	−5:15:52.7	15.95	15.15	14.97	
1258	35:49.03	−5:15:37.6	13.13	12.33	12.07	0535490−051537
1259	35:49.66	−5:00:34.7	17.47	16.16	15.67	
1260	35:49.66	−5:06:02.7	13.26	12.57	12.46	0535496−050602
1261	35:49.80	−5:06:39.7	15.48	14.69	14.41	0535497−050639
1262	35:49.81	−4:58:22.8	18.37	17.06	16.58	
1263	35:49.84	−5:15:12.8	15.44	14.65	14.24	
1264	35:50.03	−5:17:17.8	14.98	14.01	13.54	0535500−051718
1265	35:50.06	−4:58:16.1	14.46	13.26	12.93	0535500−045816
1266	35:50.08	−5:09:46.2	13.39	12.48	12.15	0535500−050946
1267	35:50.12	−5:10:29.4	13.22	12.64	12.37	0535501−051029
1268	35:50.12	−5:17:00.2	17.76	16.51	16.60	
1269	35:50.41	−5:17:29.9	17.44	16.38	15.82	
1270	35:50.56	−5:16:11.7	17.67	16.13	15.93	
1271	35:50.58	−5:09:20.9	14.29	13.18	12.55	
1272	35:50.66	−5:14:58.9	17.39	16.58	14.31	
1273	35:50.76	−5:16:28.9	11.36	10.65†	10.21†	0535507−051629
1274	35:50.77	−5:16:53.3	16.69	15.46	15.15	
1275	35:50.82	−5:09:29.9	8.49†	7.20†	6.67†	0535508−050930
1276	35:50.84	−5:05:49.3	17.74	16.41	15.65	
1277	35:51.02	−5:17:33.0	17.36	16.41	15.31	
1278	35:51.05	−5:15:08.7	11.81	11.30†	10.97†	0535510−051508
1279	35:51.08	−5:07:08.7	11.01	10.28†	9.87†	0535510−050708
1280	35:51.31	−5:11:59.9	15.92	14.98	14.80	
1281	35:51.45	−5:08:01.8	15.59	14.30	13.52	

(*cont.*)

ID	R.A.^{a}	decl.^{a}	J^{bcd}	H^{bcd}	K^{cd}	2MASS
1282	35:51.49	−5:17:22.6	18.49	17.31	16.57	
1283	35:51.55	−4:57:11.7	19.45	17.98	16.07	
1284	35:51.56	−5:09:01.9	14.97	14.14	13.83	0535515−050901
1285	35:51.56	−5:15:09.5	16.79	16.19	16.02	
1286	35:51.64	−5:08:09.0	10.05†	9.42†	9.24†	0535516−050809
1287	35:51.78	−5:17:39.2	13.44	12.68	12.33	0535517−051739
1288	35:51.82	−4:58:45.2	18.65	17.21	16.33	
1289	35:51.84	−5:15:30.5	14.80	14.14	13.93	0535518−051530
1290	35:51.93	−5:14:16.5	14.72	13.87	13.65	0535519−051416
1291	35:51.94	−5:11:56.6	17.41	16.52	16.25	
1292	35:51.96	−5:03:57.7	16.41	15.48	15.23	0535519−050357
1293	35:52.26	−5:13:18.6	13.91	13.24	12.87	0535522−051318
1294	35:52.29	−4:57:41.8	19.09	17.02	16.54	
1295	35:52.33	−5:12:57.0	12.84	12.18	11.87	0535523−051256
1296	35:52.48	−4:57:04.2	16.64	15.07	14.52	0535524−045704
1297	35:52.63	−5:05:05.3	9.84†	9.26†	8.85†	0535526−050505
1298	35:52.77	−5:12:59.0	11.73	11.07†	10.45†	0535527−051258
1299	35:52.88	−4:59:29.0	17.84	16.52	15.81	
1300	35:52.89	−4:58:34.4	16.04	15.09	14.68	0535528−045834
1301	35:52.95	−5:01:18.3	15.21	14.19	13.93	0535529−050118
1302	35:53.09	−5:10:17.9	17.42	16.02	15.65	
1303	35:53.31	−4:57:26.2	18.56	17.09	16.09	
1304	35:53.35	−5:01:16.0	16.83	15.78	15.36	
1305	35:53.44	−5:09:39.2	17.62	16.07	15.47	
1306	35:53.54	−5:02:34.7	13.02	11.99	11.62	0535535−050234
1307	35:53.64	−5:10:09.6	15.67	14.75	14.35	0535536−051009
1308	35:53.71	−5:02:00.7	15.88	15.00	14.64	0535537−050200
1309	35:53.99	−5:10:40.0	15.59	15.00	14.85	0535539−051039
1310	35:54.00	−5:00:30.7	17.72	16.51	16.37	
1311	35:54.03	−5:04:14.9	11.32†	10.43†	9.90†	0535540−050414
1312	35:54.09	−5:13:29.8	16.79	16.30	16.01	
1313	35:54.19	−5:14:43.8	17.22	15.95	15.93	
1314	35:54.20	−5:05:45.4	13.43	12.79	12.61	0535542−050545
1315	35:54.21	−5:13:43.0	>21.91	18.38	17.27	
1316	35:54.32	−5:11:52.3	16.98	16.07	15.82	
1317	35:54.68	−5:05:21.3	16.93	15.79	15.08	
1318	35:54.69	−5:06:27.8	13.34	12.85	12.57	0535546−050627
1319	35:54.77	−5:10:55.6	13.21	12.61	12.47	0535547−051055
1320	35:54.80	−4:59:20.2	16.40	15.17	14.75	0535548−045920

(*cont.*)

ID	R.A.[a]	decl.[a]	J^{bcd}	H^{bcd}	K^{cd}	2MASS
1321	35:54.82	−5:00:48.9	16.87	16.00	15.89	
1322	35:54.95	−5:10:27.9	14.48	13.78	13.59	0535549−051027
1323	35:54.97	−5:13:15.5	12.35	11.64	11.43	0535549−051315
1324	35:55.03	−5:03:16.1	17.03	16.00	15.66	
1325	35:55.06	−5:02:37.3	13.87	12.72	12.33	0535550−050237
1326	35:55.18	−5:14:42.6	17.25	15.87	15.60	
1327	35:55.35	−5:14:13.9	16.14	15.32	14.92	
1328	35:55.41	−5:09:50.2	16.39	15.60	15.19	0535554−050950
1329	35:55.45	−5:13:55.3	13.15	12.52	12.29	0535554−051355
1330	35:55.60	−5:10:00.4	17.22	16.58	16.52	
1331	35:55.65	−5:08:11.9	16.71	15.97	15.40	
1332	35:55.74	−5:04:37.7	16.99	15.68	14.18	
1333	35:55.94	−5:00:07.5	14.69	14.09	13.78	0535559−050007
1334	35:56.00	−5:05:37.3	16.50	15.40	15.15	0535560−050537
1335	35:56.02	−5:00:51.5	14.08	13.05	12.79	0535560−050051
1336	35:56.02	−5:12:09.0	15.23	14.29	14.01	0535560−051209
1337	35:56.04	−5:09:03.1	11.17	10.39†	10.12†	0535560−050903
1338	35:56.49	−5:10:30.7	14.95	14.35	14.22	0535564−051030
1339	35:56.55	−4:58:58.7	15.45	14.44	14.09	0535565−045858
1340	35:56.61	−5:02:16.5	17.40	16.29	15.75	
1341	35:56.69	−4:59:02.5	17.56	16.43	16.07	
1342	35:56.81	−5:13:28.9	17.65	16.58	16.44	
1343	35:56.84	−4:59:14.6	14.68	14.11	13.77	0535568−045914
1344	35:57.04	−5:01:40.7	17.66	16.52	15.94	
1345	35:57.14	−5:01:49.5	17.64	16.79	16.23	
1346	35:57.15	−5:02:31.9	16.93	15.94	15.72	
1347	35:57.24	−5:07:22.4	16.84	16.11	15.70	
1348	35:57.47	−5:05:32.2	17.36	16.29	16.05	
1349	35:57.47	−5:08:29.1	17.85	16.80	16.48	
1350	35:57.54	−5:08:44.2	17.44	16.60	16.34	
1351	35:57.61	−5:10:02.6	17.82	16.83	16.22	
1352	35:57.66	−5:10:29.6	16.73	16.19	15.97	0535576−051029
1353	35:57.89	−5:12:18.0	15.97	15.33	15.05	
1354	35:58.07	−5:12:54.3	9.15†	8.30†	8.05†	0535580−051254
1355	35:58.10	−5:11:42.7	15.49	14.54	13.85	
1356	35:58.17	−5:13:07.1	17.04	16.27	15.98	
1357	35:58.19	−5:11:22.9	17.81	16.58	16.00	
1358	35:58.20	−5:11:53.8	16.51	15.61	15.41	
1359	35:58.22	−5:09:32.1	15.16	14.45	14.22	0535582−050932

(*cont.*)

ID	R.A.[a]	decl.[a]	J^{bcd}	H^{bcd}	K^{cd}	2MASS
1360	35:58.27	−5:00:48.1	17.23	16.09	15.93	
1361	35:58.37	−5:06:27.4	16.95	16.19	15.75	
1362	35:58.74	−4:59:32.2	16.07	14.85	14.47	0535587−045932
1363	35:58.75	−4:59:19.8	17.66	16.18	15.81	
1364	35:58.81	−5:10:14.6	15.23	14.66	14.12	0535588−051014
1365	35:58.82	−5:00:26.1	15.03	14.38	14.07	0535588−050026
1366	35:59.00	−5:12:40.6	16.12	15.38	14.96	
1367	35:59.22	−4:58:46.2	12.55	11.83	11.61	0535592−045846
1368	35:59.22	−5:07:33.7	14.81	14.24	13.82	0535592−050733
1369	35:59.48	−5:10:21.4	16.22	15.19	15.27	0535595−051021
1370	35:59.61	−5:01:28.7	13.52	12.62	12.34	0535596−050128
1371	35:59.67	−5:01:38.6	16.59	15.31	15.03	0535596−050138
1372	35:59.71	−5:06:43.5	15.10	14.28	14.26	0535597−050643
1373	35:59.91	−5:00:37.7	17.34	15.98	15.49	
1374	35:59.93	−5:04:31.0	13.41	12.82	12.52	0535599−050430
1375	36:00.33	−5:05:00.0	12.23	11.66	11.44	0536003−050459
1376	36:00.45	−5:05:54.0	13.38	12.83	12.61	0536004−050553
1377	36:00.77	−4:59:14.7	15.46	14.53	14.24	0536007−045914
1378	36:00.92	−5:08:48.7	15.88	15.17	15.04	0536009−050848
1379	36:01.07	−5:03:09.3	17.11	15.99	15.79	
1380	36:01.27	−5:01:59.2	16.18	15.23	14.93	0536012−050158
1381	36:01.28	−5:07:56.4	15.44	14.71	14.54	0536012−050756
1382	36:01.31	−5:07:39.1	17.35	16.34	15.62	
1383	36:01.51	−5:00:21.0	16.58	15.18	14.86	0536015−050020
1384	36:01.55	−5:11:55.3	16.95	16.08	15.86	
1385	36:01.64	−5:10:40.4	15.41	14.83	14.69	0536016−051040
1386	36:01.74	−4:59:24.3	16.37	15.21	14.67	0536017−045924
1387	36:01.87	−5:08:35.1	16.09	15.18	14.88	0536018−050835
1388	36:02.17	−5:00:13.1	17.98	16.97	16.42	
1389	36:02.52	−5:04:42.1	16.79	16.05	15.89	
1390	36:02.63	−5:07:36.5	11.31	10.69†	10.49†	0536026−050736
1391	36:02.70	−5:04:18.9	17.08	16.35	16.23	
1392	36:02.73	−5:04:44.7	17.67	16.57	16.25	
1393	36:02.81	−5:04:23.7	14.89	14.39	14.06	0536028−050423
1394	36:02.87	−5:09:00.1	17.25	16.43	16.27	
1395	36:02.91	−5:05:43.3	17.61	16.30	15.65	
1396	36:02.93	−5:07:53.2	17.16	16.13	15.45	
1397	36:03.03	−5:01:11.0	17.58	16.40	15.97	
1398	36:03.04	−5:07:18.7	15.16	14.36	14.16	0536030−050718

(*cont.*)

ID	R.A.[a]	decl.[a]	J^{bcd}	H^{bcd}	K^{cd}	2MASS
1399	36:03.14	−4:58:58.9	15.71	14.73	14.49	0536031−045859
1400	36:03.21	−5:07:09.6	17.39	16.55	16.44	
1401	36:03.25	−5:09:30.9	15.66	14.86	14.70	0536032−050931
1402	36:03.29	−5:04:21.4	16.59	15.69	15.32	0536032−050421
1403	36:03.39	−5:04:00.8	14.81	14.03	13.85	0536034−050401
1404	36:03.53	−5:07:31.0	17.73	16.81	16.54	
1405	36:03.54	−5:04:45.2	16.74	15.73	15.51	0536035−050445
1406	36:04.15	−5:04:08.8	12.77	12.02	11.67	0536041−050409
1407	36:04.32	−5:07:15.6	10.90†	10.57†	10.51†	0536043−050715
1408	36:04.48	−5:07:00.9	16.04	15.36	15.16	
1409	36:04.51	−5:09:28.6	11.23	10.62†	10.44†	0536045−050929
1410	36:04.70	−5:13:03.7	>21.66	>20.71	16.61	
1411	36:04.86	−4:58:45.1	15.67	14.73	14.39	0536048−045845
1412	36:04.87	−4:59:32.6	17.29	16.02	15.72	
1413	36:04.97	−4:59:41.5	11.71	11.20	11.13	
1414	36:05.07	−5:10:57.9	>21.41	>20.46	16.14	
1415	36:05.09	−4:59:42.9	11.75	11.24	11.16	0536050−045943
1416	36:05.09	−5:03:12.2	15.35	14.68	14.55	0536050−050312
1417	36:05.11	−5:11:13.4	13.78	13.32	12.88	0536051−051113
1418	36:05.43	−5:04:43.6	16.45	16.03	14.15	
1419	36:05.48	−5:07:58.6	15.62	15.09	14.69	0536054−050758
1420	36:05.87	−5:02:18.2	14.63	13.97	13.77	0536058−050218
1421	36:05.94	−5:00:41.0	14.92	14.00	13.77	0536059−050041
1422	36:06.02	−5:08:15.0	15.81	15.30	15.19	0536060−050815
1423	36:06.14	−5:07:55.6	>21.72	>20.77	17.05	
1424	36:06.30	−5:09:40.1	17.30	16.32	16.03	
1425	36:06.62	−5:06:07.4	>21.91	>20.96	17.47	
1426	36:06.65	−5:03:06.9	17.24	16.24	16.06	
1427	36:06.79	−5:06:15.5	12.98	12.25	12.04	0536067−050615
1428	36:06.85	−5:08:50.8	16.68	15.93	15.93	0536068−050851
1429	36:07.07	−5:04:07.6	16.98	15.86	15.61	0536070−050407
1430	36:07.61	−5:04:15.4	17.67	16.56	16.43	
1431	36:07.71	−5:09:41.2	16.28	15.28	15.07	0536077−050941
1432	36:07.96	−5:09:37.3	17.23	16.64	16.72	
1433	36:08.26	−5:04:01.8	14.68	14.04	13.97	0536082−050401
1434	36:08.49	−5:06:14.7	17.02	16.07	15.90	
1435	36:08.54	−5:01:41.3	15.28	14.42	14.30	0536085−050141
1436	36:08.68	−5:04:34.3	15.97	14.91	14.70	0536086−050434
1437	36:08.70	−5:02:32.2	18.70	18.08	16.34	
1438	36:08.80	−5:02:35.0	15.94	14.96	14.70	0536088−050235
1439	36:08.92	−5:02:26.3	17.84	16.69	16.44	
1440	36:09.81	−5:00:49.9	13.48	12.63	12.53	0536098−050050
1441	36:09.96	−5:05:35.3	14.23	13.63	13.39	0536099−050535
1442	36:10.17	−5:07:02.7	15.51	14.55	14.36	0536101−050703
1443	36:10.57	−5:06:12.3	17.79	16.88	16.83	
1444	36:10.76	−5:04:34.0	15.85	14.92	14.84	0536107−050434
1445	36:10.81	−5:02:31.3	15.50	14.80	14.71	0536108−050231
1446	36:11.00	−5:03:41.2	8.44†	8.22†	8.14†	0536110−050341
1447	36:11.79	−5:00:32.4	12.12	11.37	11.24	0536118−050032
1448	36:12.41	−5:00:53.8	15.97	14.73	14.55	0536124−050053

(*cont.*)

[a] The equinox in J2000.0 in the format of 05:mm:ss.ss for R.A. and dd:mm:ss.s for decl. The hundredth and the tenth order in the arc seconds of R.A. and decl. are truncated.

[b] The upper limit of the flux (the lower limit of the magnitude) is given for sources with less than a 3 σ detection in the J, H, or both bands.

[c] Saturated QUIRC magnitudes (brighter than 11 mag) are replaced with 2MASS magnitudes and labeled † if they have the 2MASS counterpart.

[d] Saturated QUIRC magnitudes (brighter than 11 mag) are labeled $<$ if they have no 2MASS counterpart, recognizing the value as the lower limit of the flux (the upper limit of the magnitude).

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